

Klystrons

(and other vacuum electronics
power amplification devices)

Gap-Coupling, Beam-Loading
Admittance

Eric Colby,

SLAC

USPAS, January 21st

Outline

- How a linac's properties interact with the power source
- General principles of microwave tubes
 - Beam loading admittance
 - Radial gap-coupling coefficient
 - Bunching
- Specific examples
- General scaling laws
- Other examples

Linac / Power Source Parameter Couplings

Linac Gradient $G' = \sqrt{P' Z'_s T}$

G' = gradient per unit length

P' = rf peak power per unit length

Z'_s = shunt impedance per unit length

T = transit time factor

Linac fill time $\tau = Q_L / 2\omega \ll \tau_{rf} = \text{rf pulse length}$

Q_L = Accelerating loaded cavity $Q_L = Q_o / (1 + \beta)$

Average Beam Power $P_b = \eta P_{rf}$

P_b = (average beam current) x (beam voltage gain)

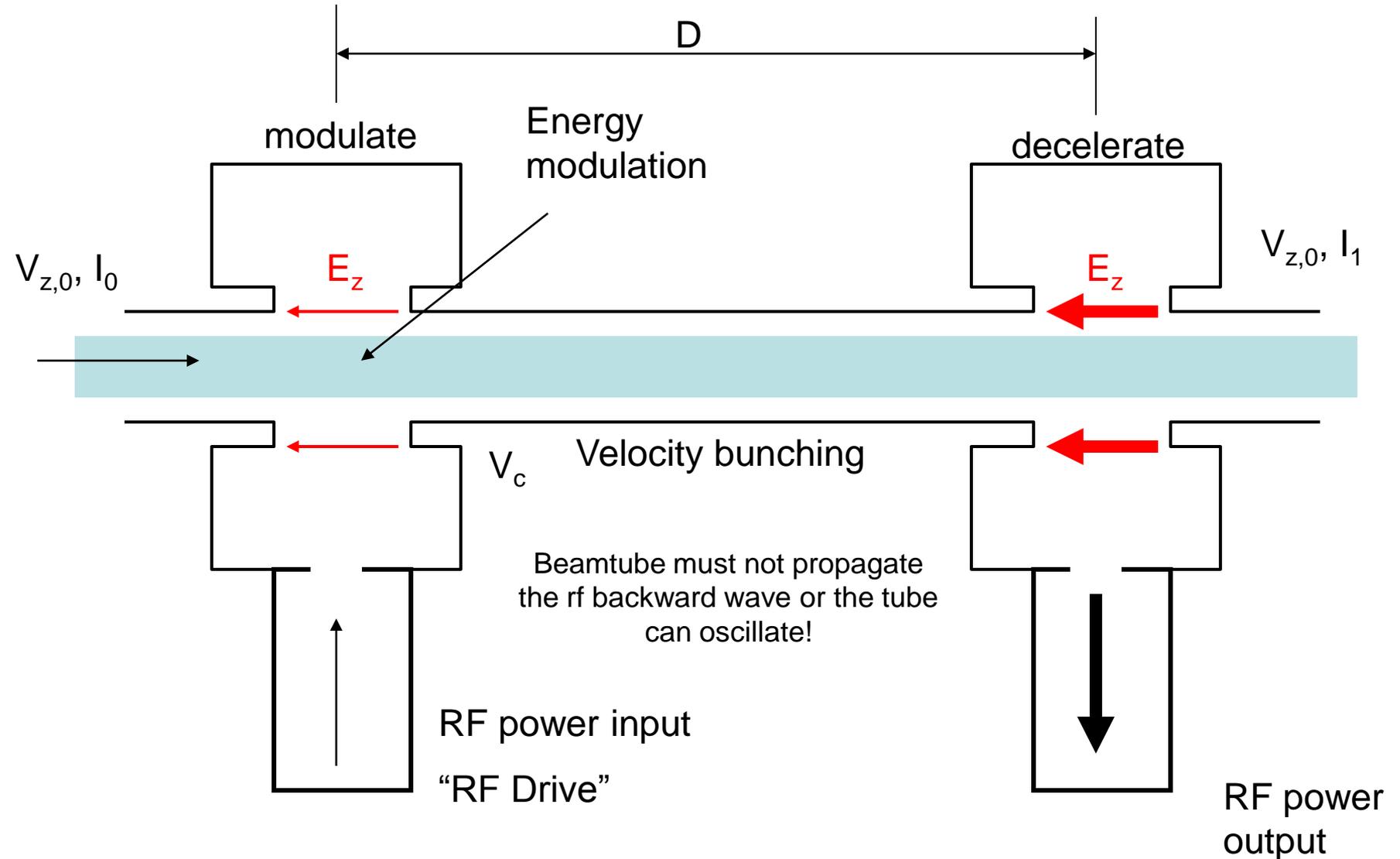
η = net efficiency of coupling rf to beam

P_{rf} = rf average power ($= f * \tau_{rf} * P_{peak}$)

Linac coupling factor $\beta = \text{VSWR (oc) @ tube}$

Phase stability of linac \Leftrightarrow phase stability of source
frequency stability of source

Basic Klystron



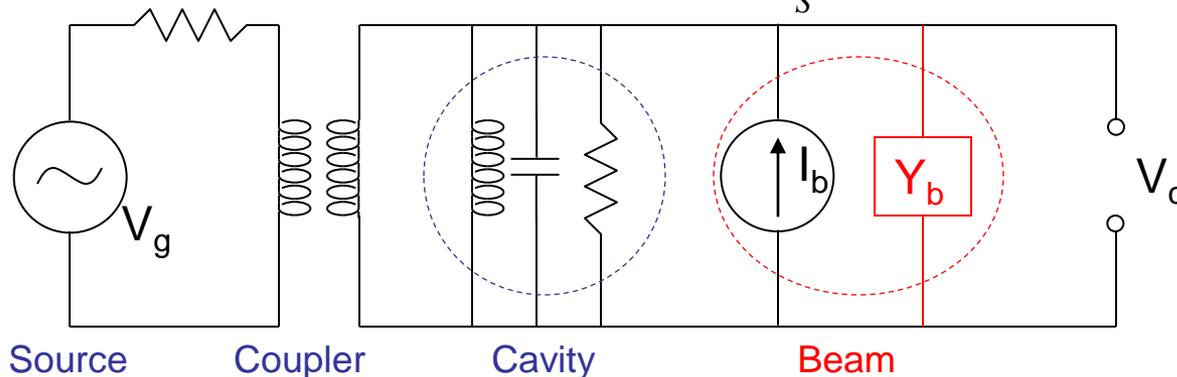
Beam Loading Admittance

So far, we've taken $\beta \rightarrow 1$ in treating beam/cavity interaction, which allows for the approximation of the beam as an ideal current source (i.e. the beam current is **not** modified by the gap voltage in the cavity).

For vacuum power tubes, however, $\beta \sim (0.2-0.8) \Leftrightarrow (V=10\text{kV}-500\text{kV})$, and the beam current changes significantly while *in the gap*.

$$\tilde{J}_1 = \tilde{\rho}_1 V_{z,0} + \rho_0 \tilde{V}_{z,1} \quad \begin{array}{l} 0 \rightarrow \text{DC} \\ 1 \rightarrow \text{rf} \end{array}$$

$$I_b = \int_S \tilde{J}_1 dS$$



Beam Loading Admittance

Proceeding from the Continuity equation and Lorentz force law, the action of the rf fields on the beam velocity and density can be described, yielding:

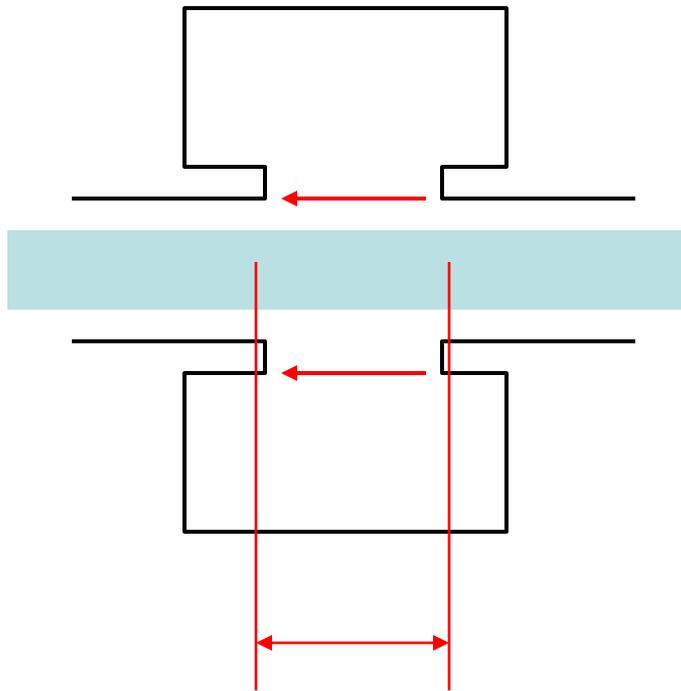
$$\tilde{I}_{rf} = -I_{dc} \frac{e}{m\gamma^3 V_{z,0}^2} \frac{1}{j\theta} \left[1 + e^{-j\theta} + \frac{2j}{\theta} (1 - e^{-j\theta}) \right] \tilde{V}_c = \tilde{Y}_c \tilde{V}_c$$

writing

$$\tilde{Y}_c = G_c + jB_c$$

$$\theta = \frac{\omega L}{V_{z,0}} \quad \text{Gap transit angle}$$

Beam Loading Admittance



Gap length L

Gap transit angle

$$\theta = \omega L / \beta c$$

Normal range

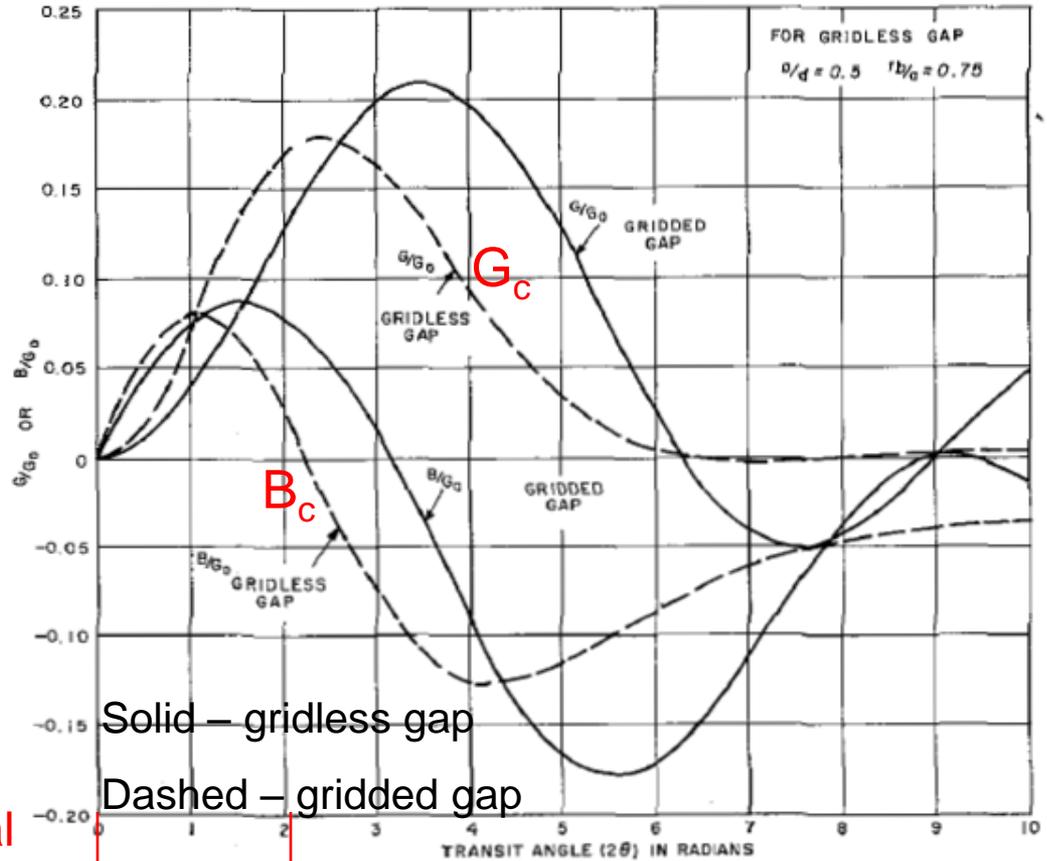


Fig. 1. The normalized conductance and susceptance of a gridless gap as a function of transit angle with $\alpha = a/d = 0.5$ and with the beam radius 75 percent of the tunnel radius. Also shown are the corresponding quantities for an ideal gridded gap which corresponds to the gridless gap with $a/d = 0$.

From E. Craig, "The Beam-Loading Admittance of Gridless Klystron Gaps", IEEE Trans. Elec. Dev. **14** (5), p. 273ff, (1967).

Gap Coupling

The beams used in vacuum tubes generally occupy a significant fraction of the beam pipe (often $r_b = 2/3 R_{\text{pipe}}$), so the variation of the gap fields with r must be accounted for. This is done with the **radial gap coupling coefficient**.

The solution to wave equation in cylindrical coordinates yields the form of the gap voltage with r :

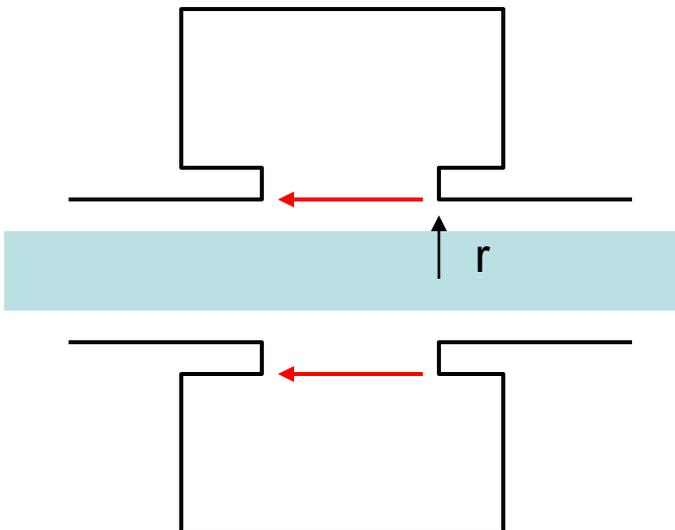
$$\tilde{V}_c(r, \omega) = \tilde{V}_c(R, \omega) \frac{I_0(\Gamma r)}{I_0(\Gamma R)}$$

$$\Gamma = \omega / \gamma \mathcal{W}_z$$

Averaging over the beam's cross section yields the gap coupling coefficient:

$$M = \frac{2I_1(\Gamma r_b)}{\Gamma r_b I_0(\Gamma R)}$$

$$[R/Q] \rightarrow [R/Q] * M^2$$



Ballistic Bunching

- Ballistic (no space charge) Theory

$$I_1 = 2I_{DC} J_1(B) \exp(j(\varphi_0 - \theta_0 + \pi / 2))$$

with

$$B = \frac{\theta_0}{\gamma_0^3 (V_{z,0} / c)^2} \frac{e|V_c|}{mc^2} \quad \text{Bunching parameter}$$

$$\theta_T = \omega D / V_{z,0} \quad \text{Transit angle of drift D}$$

- Matlab demonstration (space charge free)

Klystron Efficiency vs. Perveance

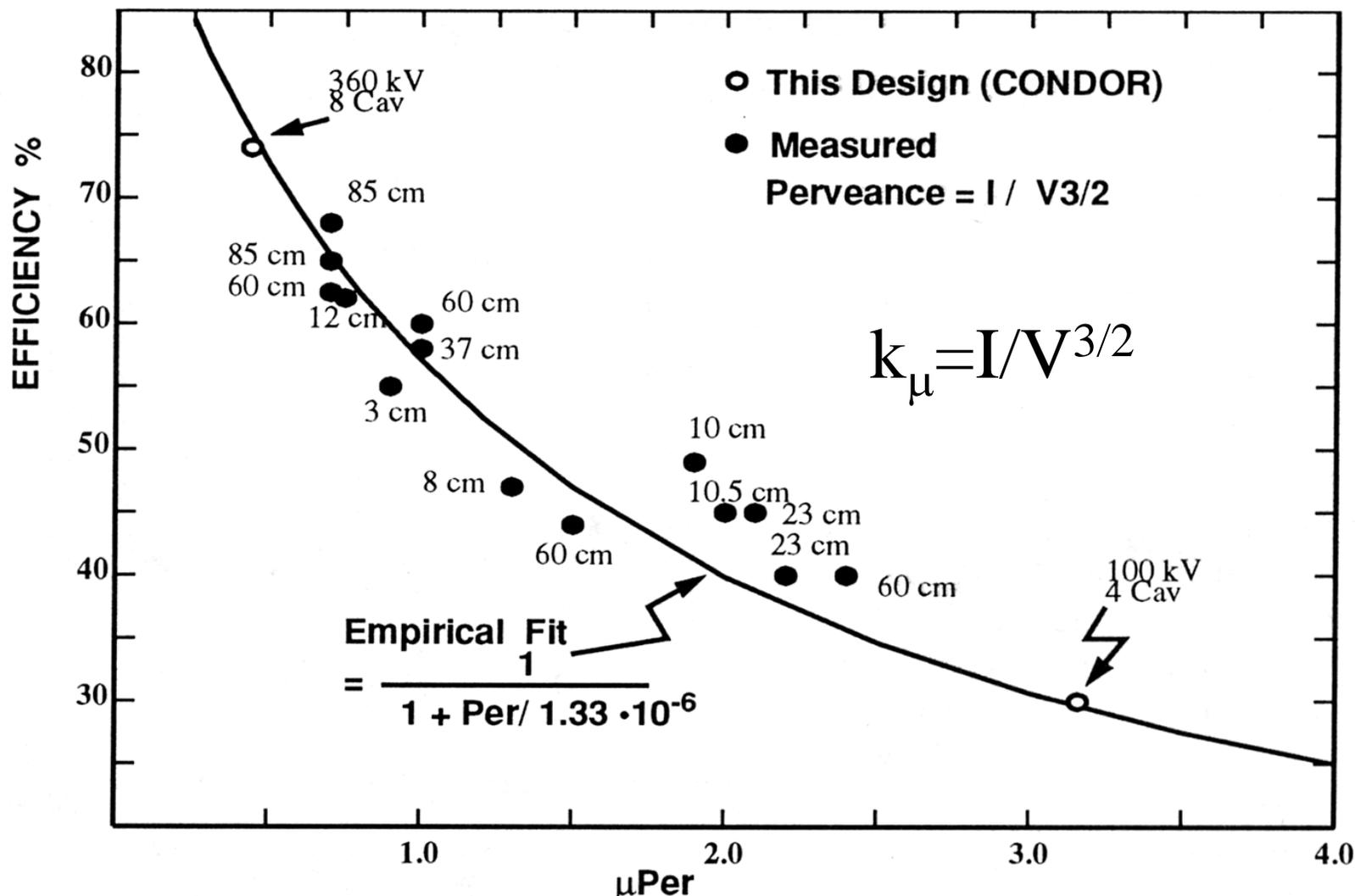
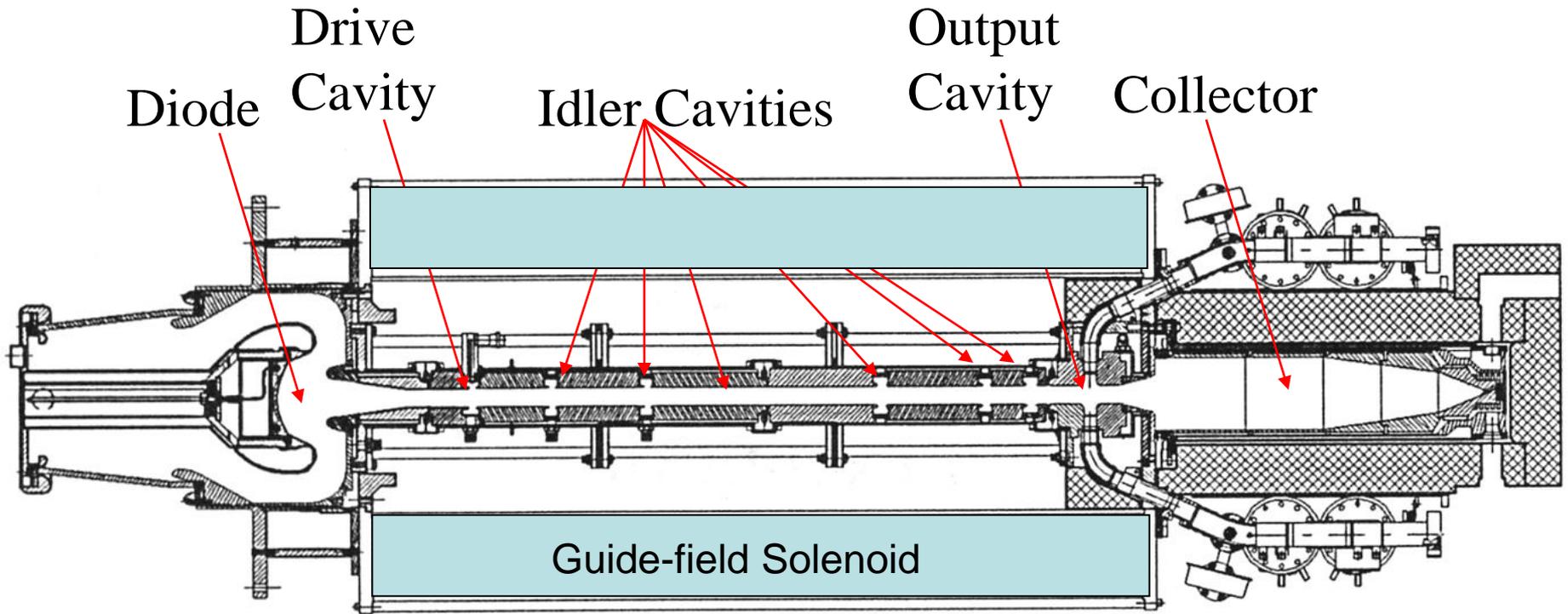


Fig.1 The empirical relation of efficiency to the perveance.

Taken from R. Palmer, *et al*, "Status of the BNL-MIT-SLAC Cluster Klystron Project", AIP Conf. Proc. 337, p. 94ff, (1994).

Real Klystron Schematic



DESY S-Band Tube (short-pulse)

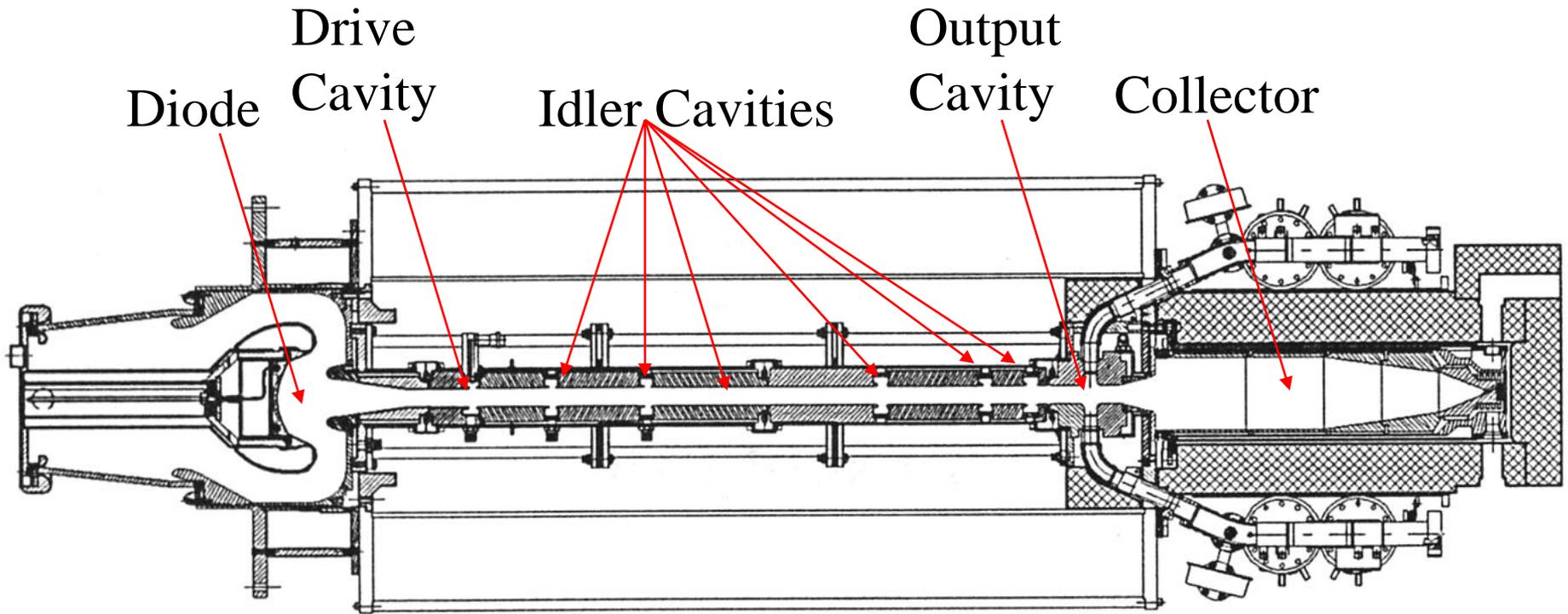


FIGURE 1. The 150-MW klystron assembly shown with magnets and lead.

$f = 2996 \text{ MHz}$

Gain = 55 dB

Efficiency: >40%

$P = 150 \text{ MW}$

$B \sim 2100 \text{ Gauss}$

PRF: 60Hz

$K = 1.8 \mu\text{P}$

Group Delay 150 nsec

Pulse length: 3 μs

$V_b = 535 \text{ kV}$

$J_{\text{cath}} = 6 \text{ A/cm}^2$

$I_b = 700 \text{ Amps}$

B-Factory Tube (CW)

$f = 476 \text{ MHz}$

Gain > 43 dB

Efficiency: >50%

$P = 1.2 \text{ MW CW}$

BW = +/- 3.0 MHz at -1dB points

VSWR tolerance: 1.2

$K = 0.83\text{-}1.3$

Group Delay 150 nsec

$V_b = 83 \text{ kV}$

$J_{\text{cath}} = 0.63 \text{ A/cm}^2$

$I_b = 24 \text{ Amps CW}$

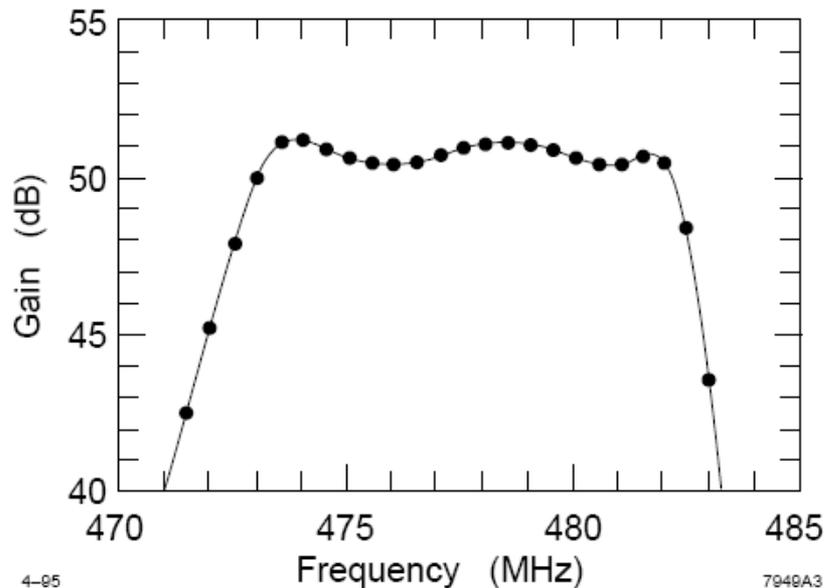


Figure 1. Predicted Small-Signal Gain versus Frequency Response at 90 KV

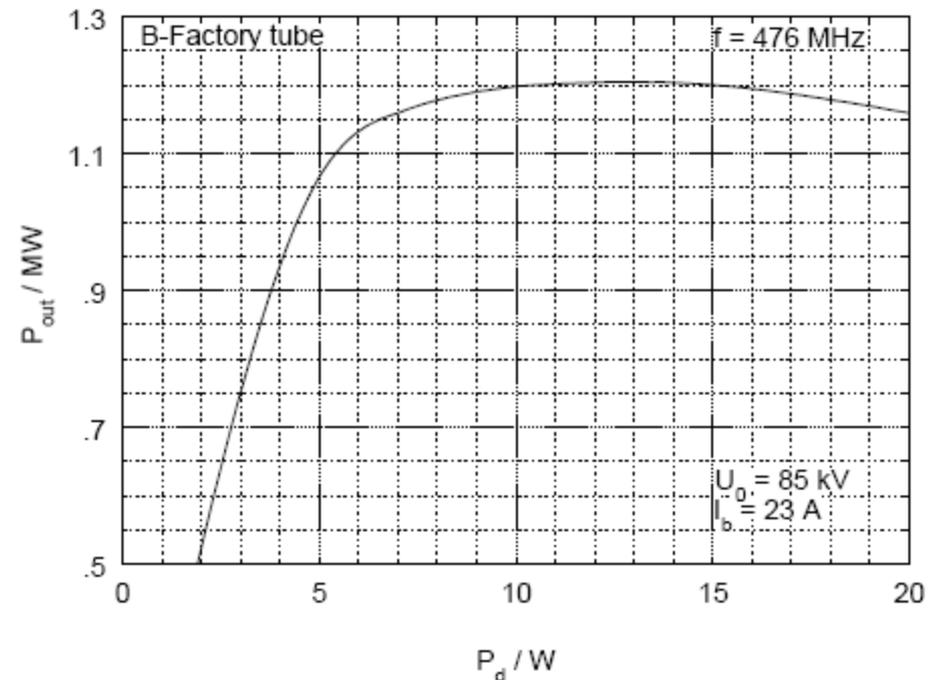
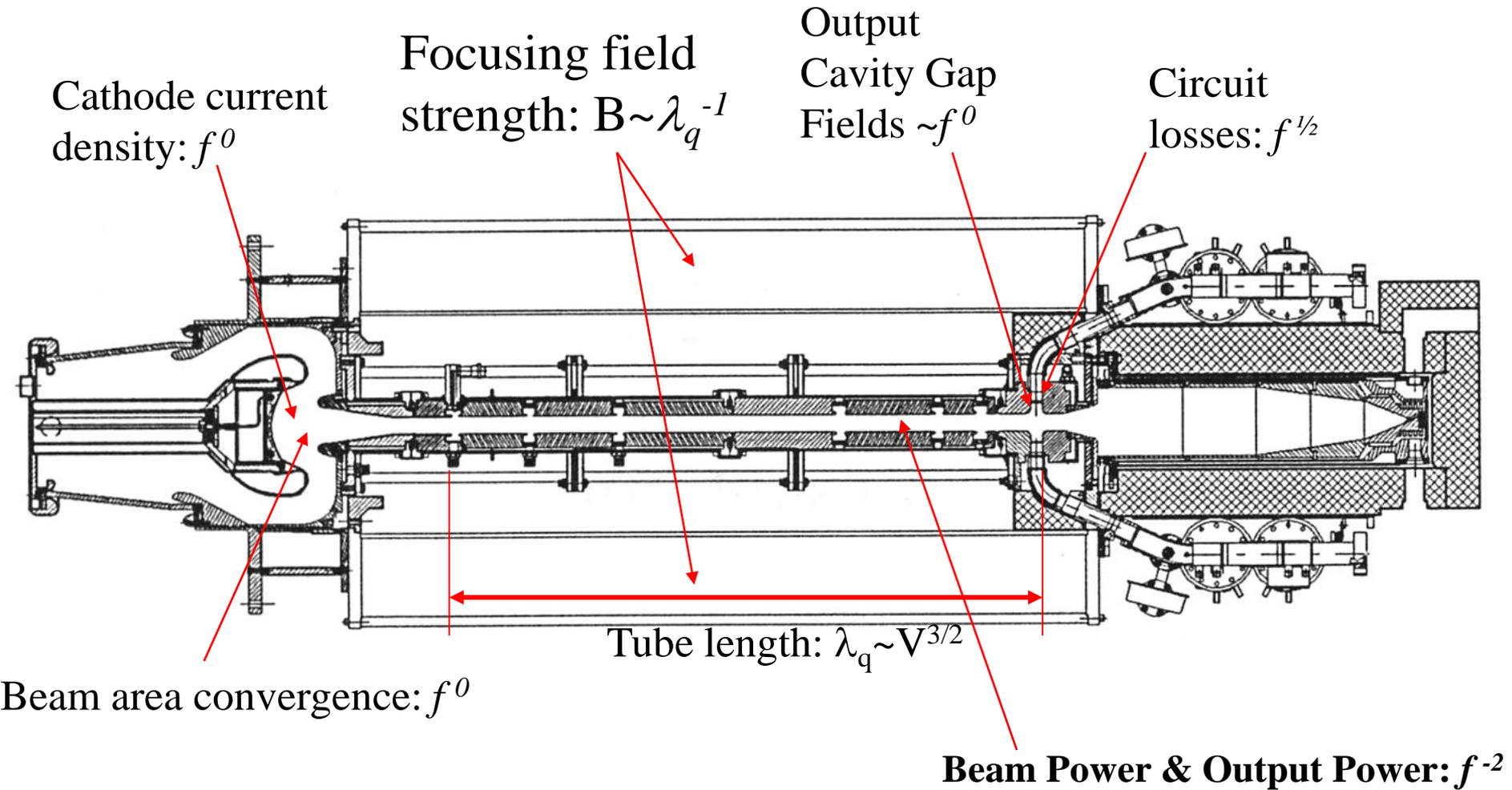


Figure 2: output power versus drive power.

Klystron Amplifier Scalings



RF PULSED HEATING OF CAVITY WALLS

[O.A. Nezhevenko, PAC'97, Vancouver, 1997, p. 3013.]

Structure lifetime is limited by fatigue from pulsed heating in a thermal-diffusion layer.

 Assume OFE copper NLC-like cavities, “safe” temperature excursions of 110 °C, and collider operation at 120 Hz, 6000 hrs/yr. [For 20 year life, $n = 5 \times 10^{10}$ pulses.]

[Experiments (Pritzkau, 2000) achieved $\Delta T = 120$ °C for 5.5×10^7 pulses, observed surface damage.]

CALCULATED EFFECTS OF RF PULSE HEATING OF IDEAL CAVITY WALLS.

| freq ω (GHz) | $G_0 \sim \omega^{5/6}$ (MV/m) | $\tau_p \sim \omega^{-1.5}$ (ns) | $\Delta T \sim \omega^{1.42}$ (°C) | $T_o = 27$ °C | | $T_o = 77$ °C | |
|------------------------|-----------------------------------|-------------------------------------|---------------------------------------|----------------------|----------|-------------------|----------|
| | | | | n | lifetime | n | lifetime |
| 11.424 | 100 | 250 | 23 | | | | |
| 22.848 | 180 | 90 | 61 | | | | |
| 34.272 | 250 | 50 | 112 | | | | |
| 45.696 | 320 | 30 | 165 | 3.6×10^{10} | 14 yrs | 2.5×10^9 | 1 yr |
| 57.120 | 380 | 22 | 225 | 1.5×10^9 | 3500 hrs | 1.8×10^8 | 420 hrs |
| 68.544 | 440 | 17 | 290 | 1×10^8 | 230 hrs | 1.8×10^7 | 42 hrs |
| 79.968 | 500 | 14 | 370 | 9.8×10^6 | 23 hrs | 2.6×10^6 | 6 hrs |
| 91.392 | 560 | 11 | 440 | 1.4×10^6 | 3 hrs | 4.7×10^5 | 1 hr |
| 102.816 | 620 | 9 | 520 | 2.6×10^5 | 36 min | 1×10^5 | 14 min |
| 114.240 | 680 | 8 | 600 | 6×10^4 | 8 min | 3×10^4 | 4 min |

For collider applications, RF technology development >40 GHz appears unjustified/

Omega-P, Inc. – Yale University
Beam Physics Collaboration

Typical operational problems with a klystron

- Input drive power is too low and klystron is not saturated (although this is *required* if the tube is included in a feedback loop) → reduced output power, efficiency, and increased output power amplitude jitter
- Wrong cathode filament current → wrong perveance → poor efficiency (minor) or beam interception (major)
- Modulator voltage jitter causes jitter in
 - The tube output power $P \sim V^{5/2} \rightarrow \delta P/P \sim (5/2)(\delta V/V)$
 - The tube transit time, which in turn causes phase jitter on the output $\partial\phi / \partial V = (Lf/\beta^3\gamma^3c)$
- Mismatched load causes significant output instability

Typical Klystron Failure Modes

- Burnt-out or shorted cathode filament
- Beam interception erodes output cavity
- Rf breakdown erodes the output cavity
- Multipactoring on or near the output window sputters metal onto the ceramic window, resulting in breakdowns
- Slow vacuum leak contaminates (“poisons”) the cathode
- Focusing magnetic field changes, resulting in a current density change that either:
 - If minor: changes the gap-coupling and space charge forces
 - If major: results in beam interception
- Input cavity erodes (rare)
- Output cavity oscillates (not necessarily at an integer multiple of the rf frequency)
- Gun supports an rf resonance

Examples of other beam-based power amplifiers

17.136 GHz Klystron (Haimson Research, NRL)

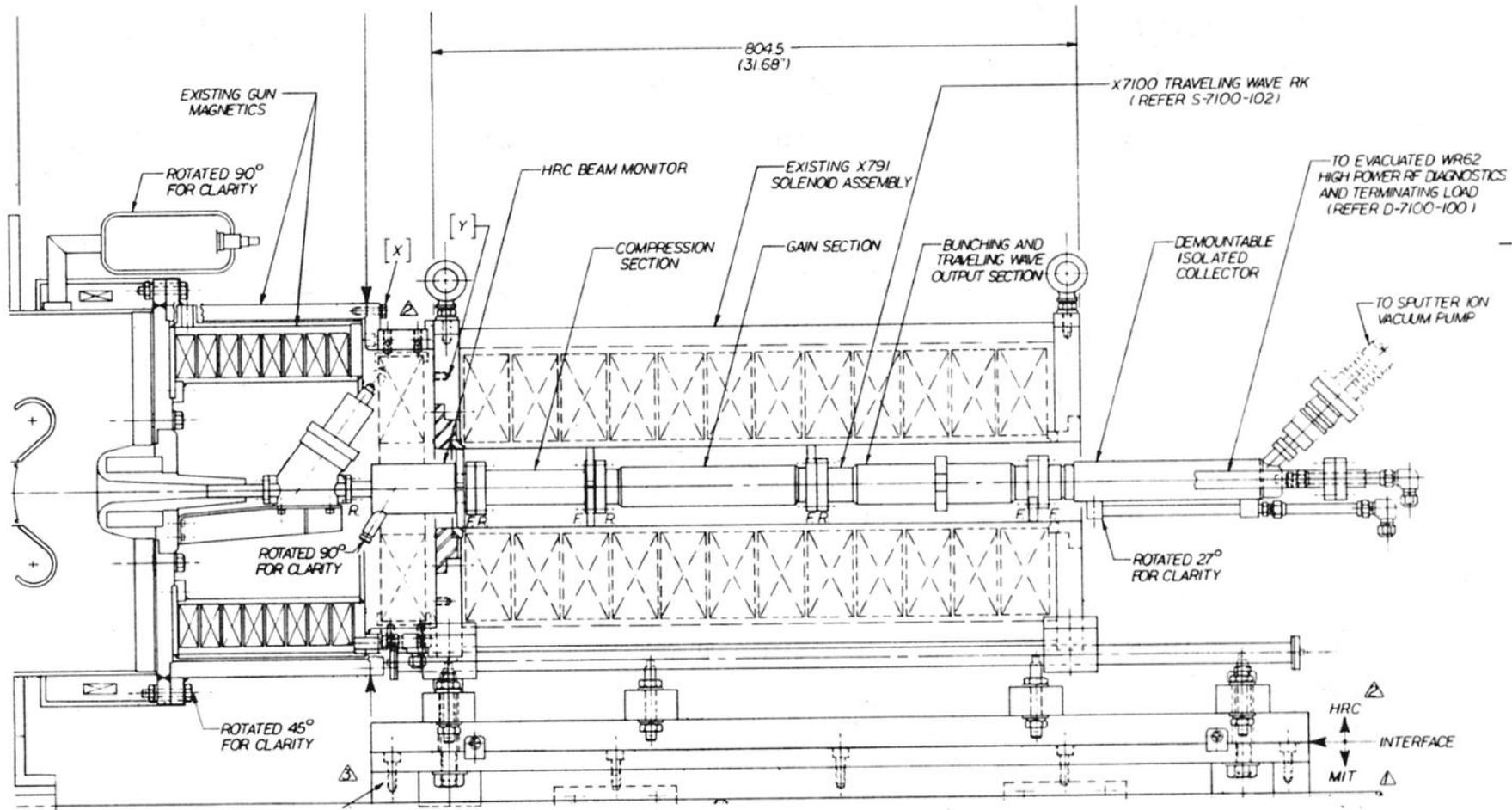


FIGURE 2. Prototype 17 GHz TWRK and solenoid assembly showing the electron gun vacuum isolation valve and the precision alignment supporting mechanism.

Taken from J. Haimson, B. Mecklenberg, B. Danly, "Initial Performance of a High Gain, High Efficiency 17 GHz Traveling Wave Relativistic Klystron for High Gradient Accelerator Research", in AIP Conf. Proc. 337, p.146ff, (1994).

17.136 GHz Klystron

(Haimson Research, NRL)

Measured Properties:

- Output Power: 26 MW @ 150 ns
- Efficiency: 49%
- Saturated Gain: 67 dB
- 560 kV/95 A beam



FIGURE 1. Model X7100 17 GHz TWRK showing the WR62 rectangular waveguide and body water-cooling connections, and the high vacuum port for the independently pumped, electrically isolated beam collector.

Taken from J. Haimson, B. Mecklenberg, B. Danly, "Initial Performance of a High Gain, High Efficiency 17 GHz Traveling Wave Relativistic Klystron for High Gradient Accelerator Research", in AIP Conf. Proc. 337, p.146ff, (1994).

1 MW, 91.4 GHz Sheet Beam Klystron (SLAC/MRC)

SBK Details:

$$I_o = 15 \text{ A}$$

$$V_o = 140 \text{ kV}$$

$$\text{Perv/} \square = 14.4 \text{ nP/} \square$$

$$L \times W = 8 \text{ mm} \times 0.4 \text{ mm}$$

$$P_{in} = 400 \text{ W (at ports)}$$

$$f_{rf} = 91.41 \text{ GHz}$$

PIC Details:

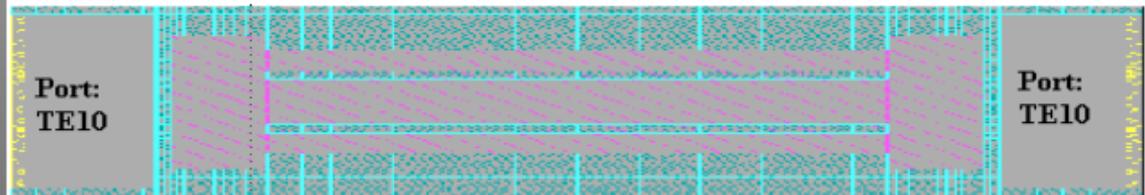
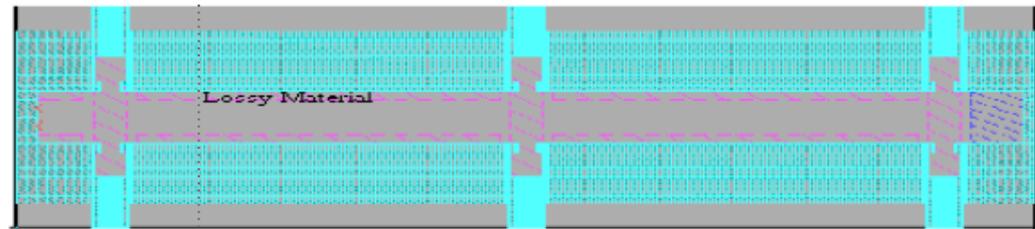
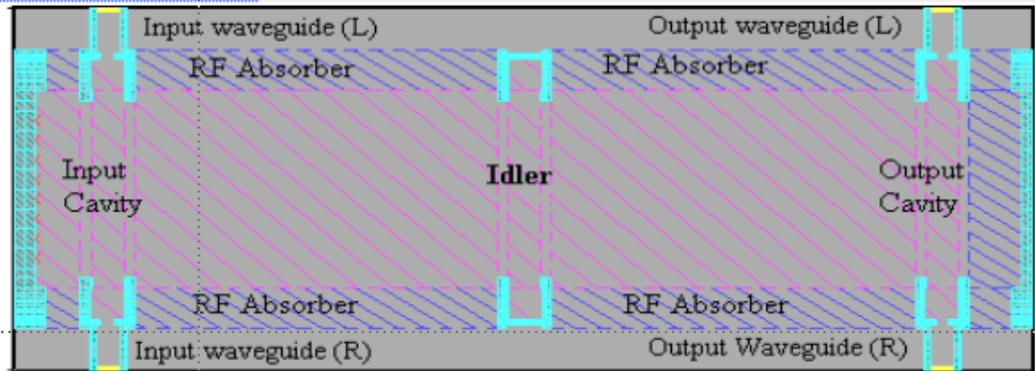
$$\delta t = 0.135 \text{ ps} = \tau_{rf}/84$$

$$\delta z = 200 \text{ } \mu\text{m} = \lambda/16$$

$$N = 198,000 \quad n = 17/\text{cell}$$

$$dt_{sim}/dt_{real} = 3.1 \text{ ns/day}$$

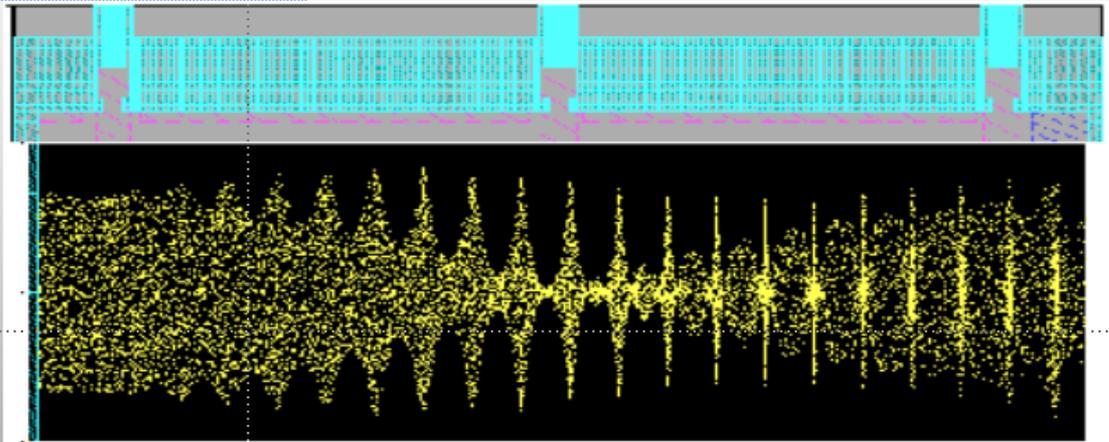
October 6, 1998



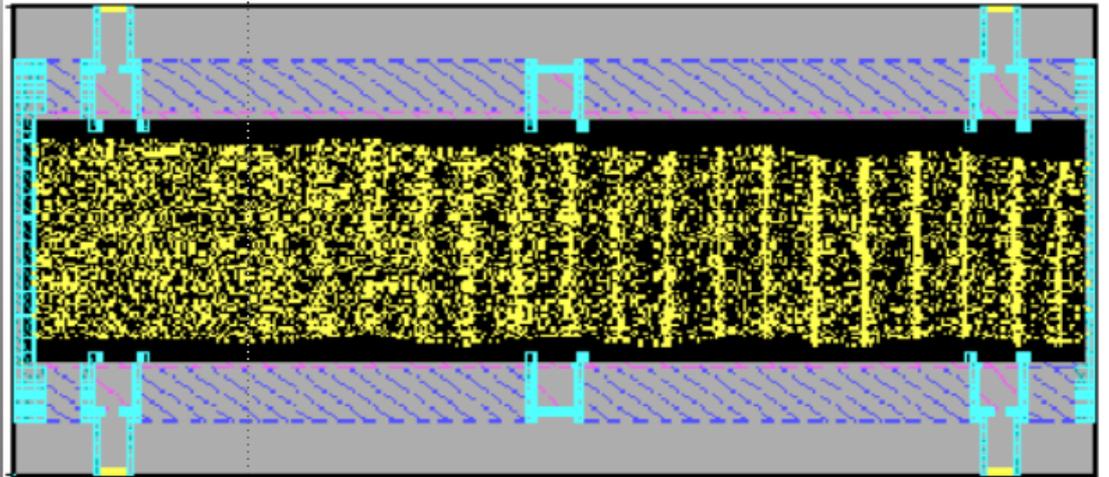
SBK Simulation

Configuration Space
Images of Bunching

Vertical Plane



Horizontal Plane



October 6, 1998

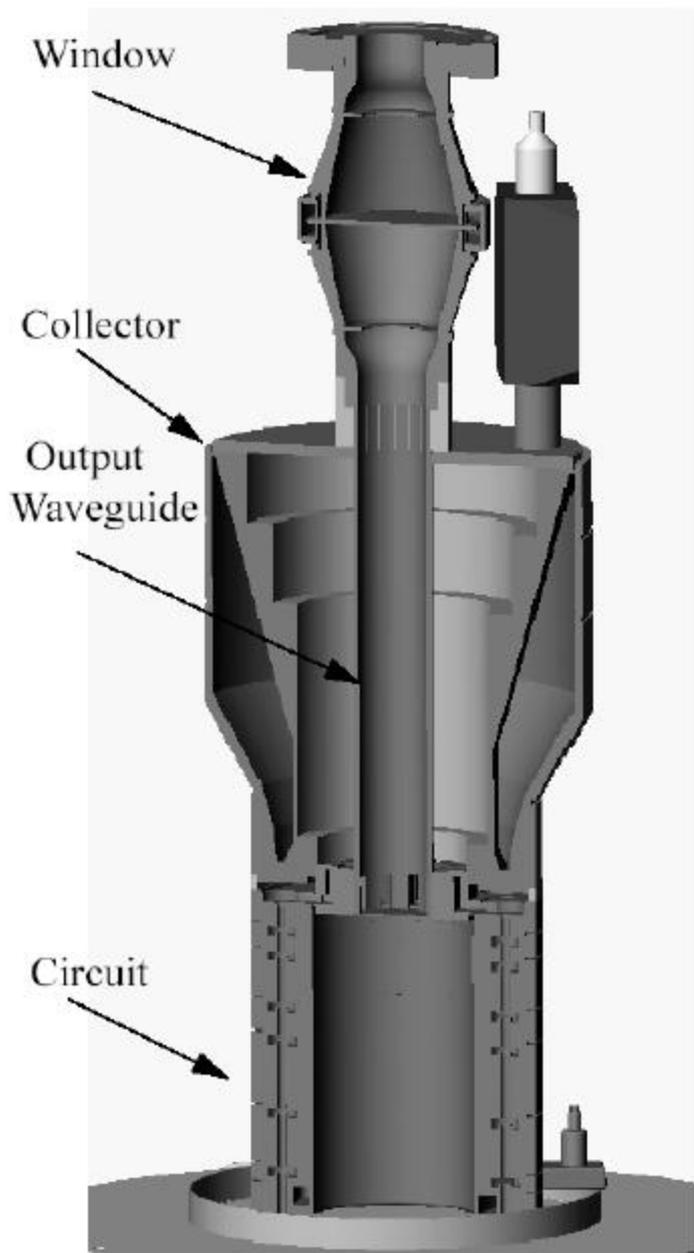


Figure 6: Cross-section of MBK

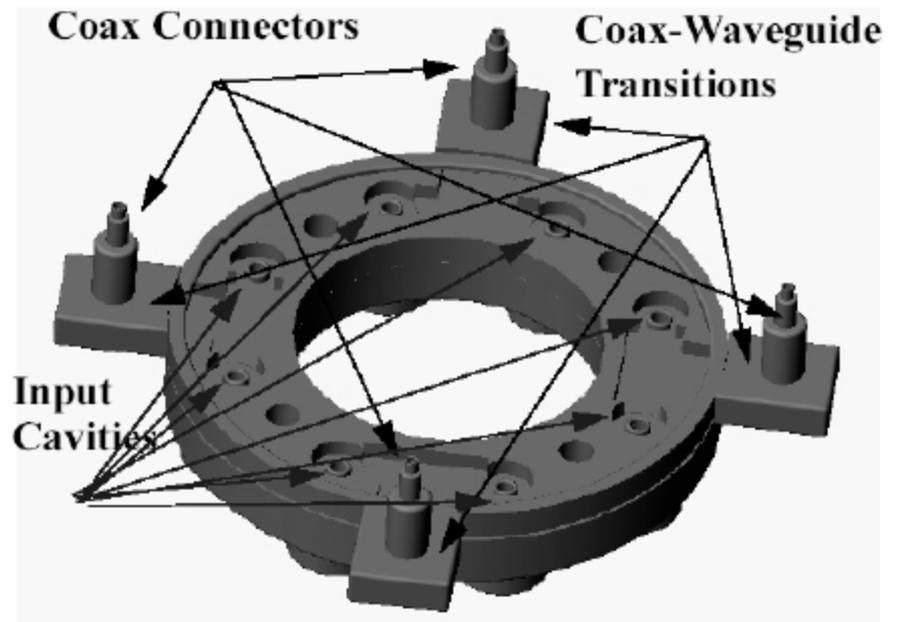


Figure 2. Input cavity assembly showing four coaxial inputs to waveguide transitions.

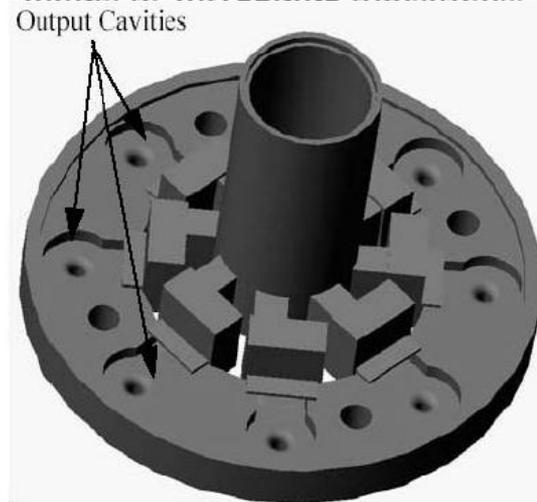
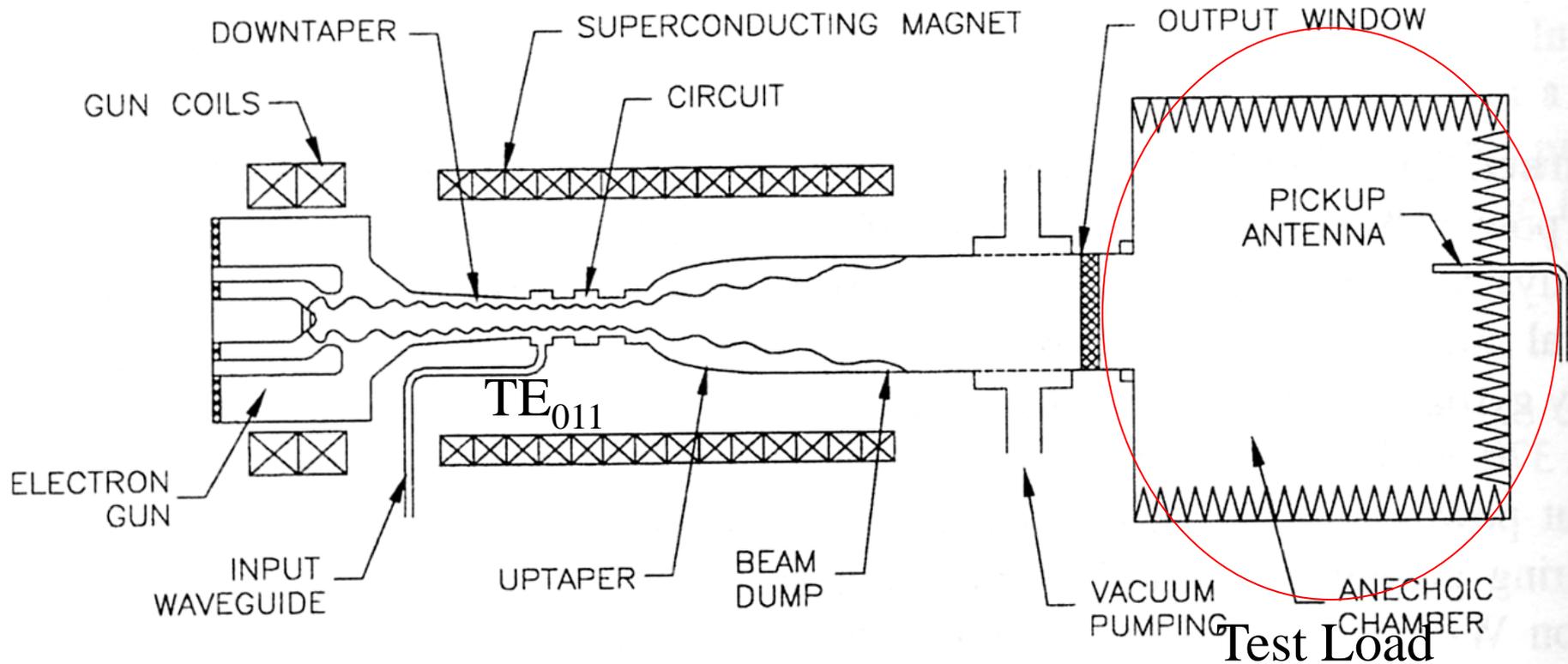


Figure 4. Output waveguide system consisting of eight waveguide twists, a TE_{01} circular waveguide, and the TE_{01} window

The Gyroklystron

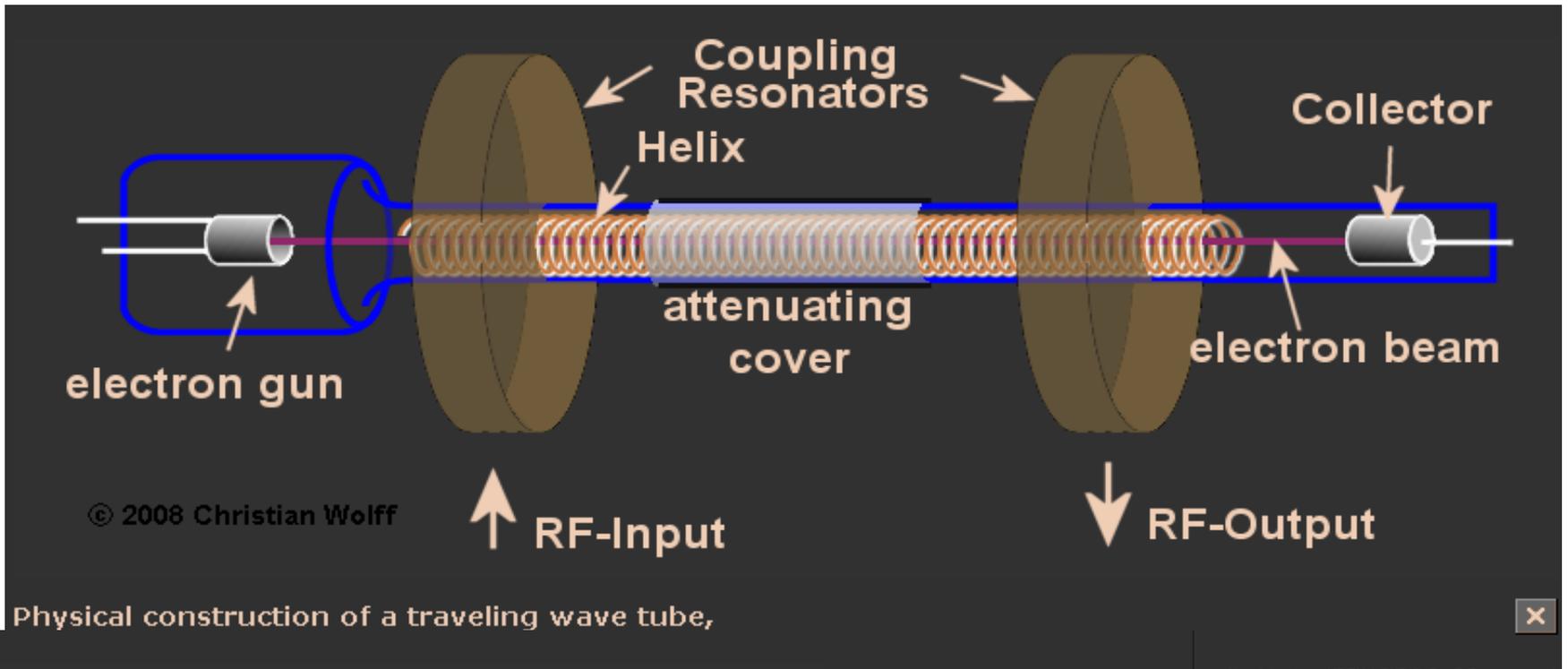


Synchronism condition: $\omega - k_{\parallel} z_{\parallel} = n\Omega_c$

FIGURE 1. Schematic diagram of the Ka-Band gyrokystron experiment.

Taken from M. Blank, *et al*, "Experimental Demonstration of High Power Millimeter Wave Gyro-Amplifiers", AIP Conf. Proc. 474, p. 165ff, (1998).

Helix TWT (Preamplifier)



Physical construction of a traveling wave tube,

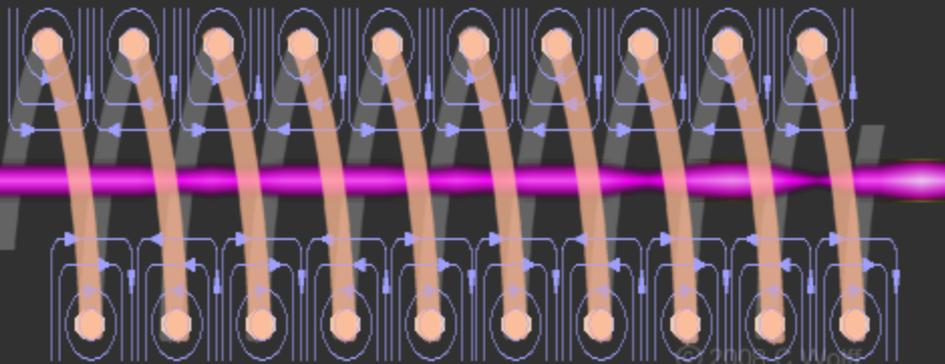


Figure 5. - electron- beam bouncing and a detail-foto of a helix
(Measure detail for 20 windings)

Generally used for driving klystrons.
Gain (typ)~20-30 dB
Output~1 kW or less

34 GHz Magnicon

(Omega-P)

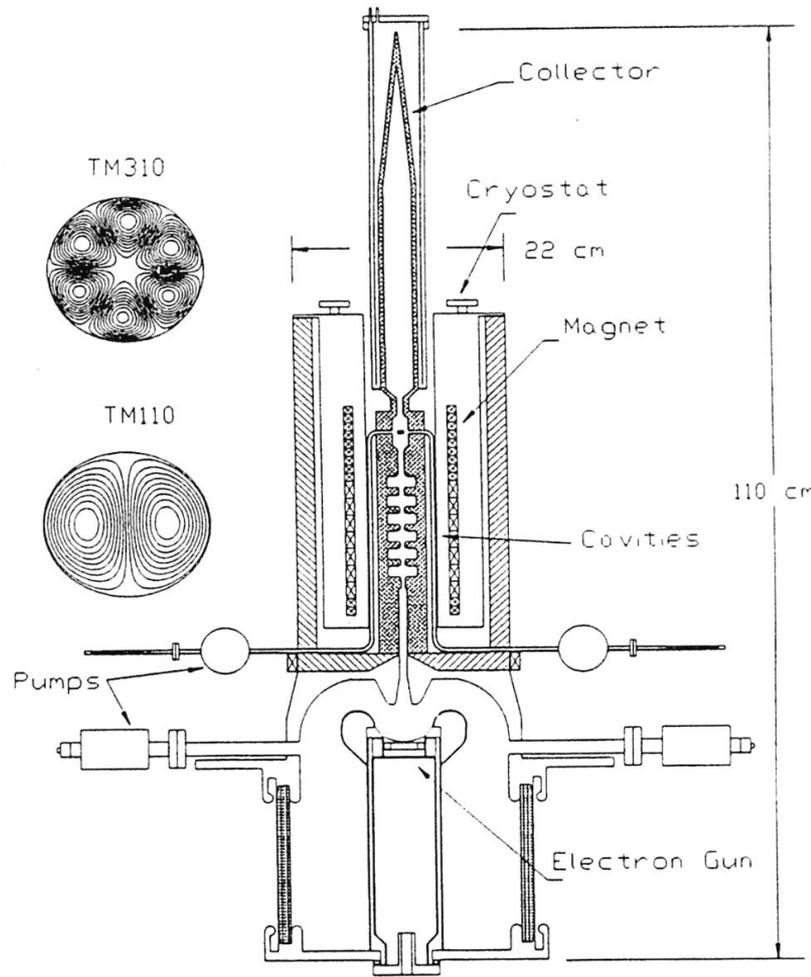


FIGURE 1. Schematic diagram of 34.272 GHz magnicon amplifier tube. Inserts at upper left show rf field patterns for drive and all deflection cavities (TM₁₁₀ mode at 11.424 GHz), and for the output cavity (TM₃₁₀ mode at 34.272 GHz).

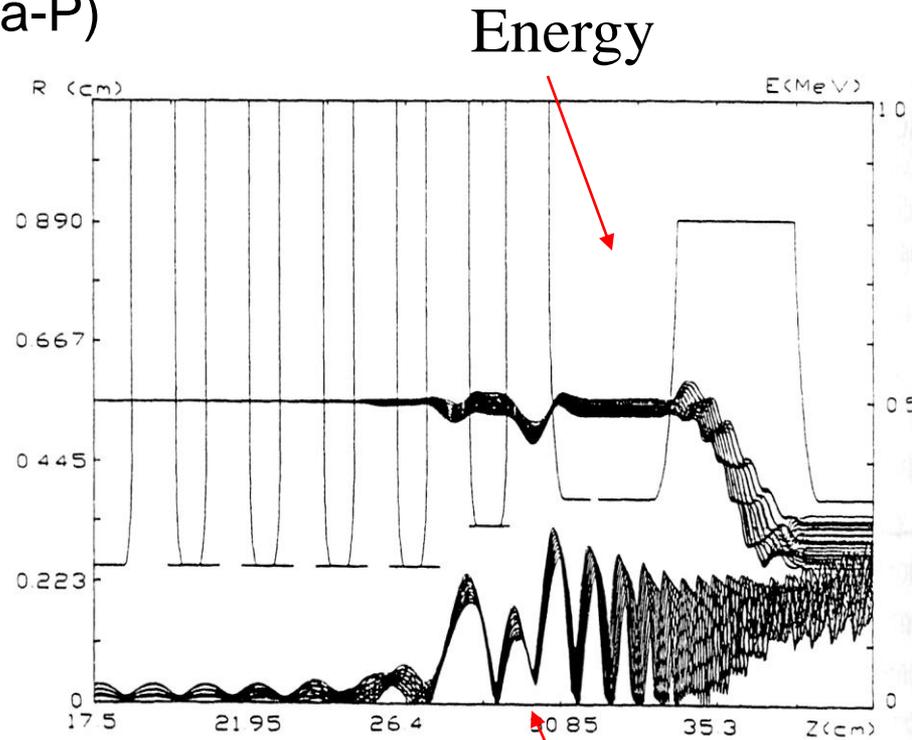
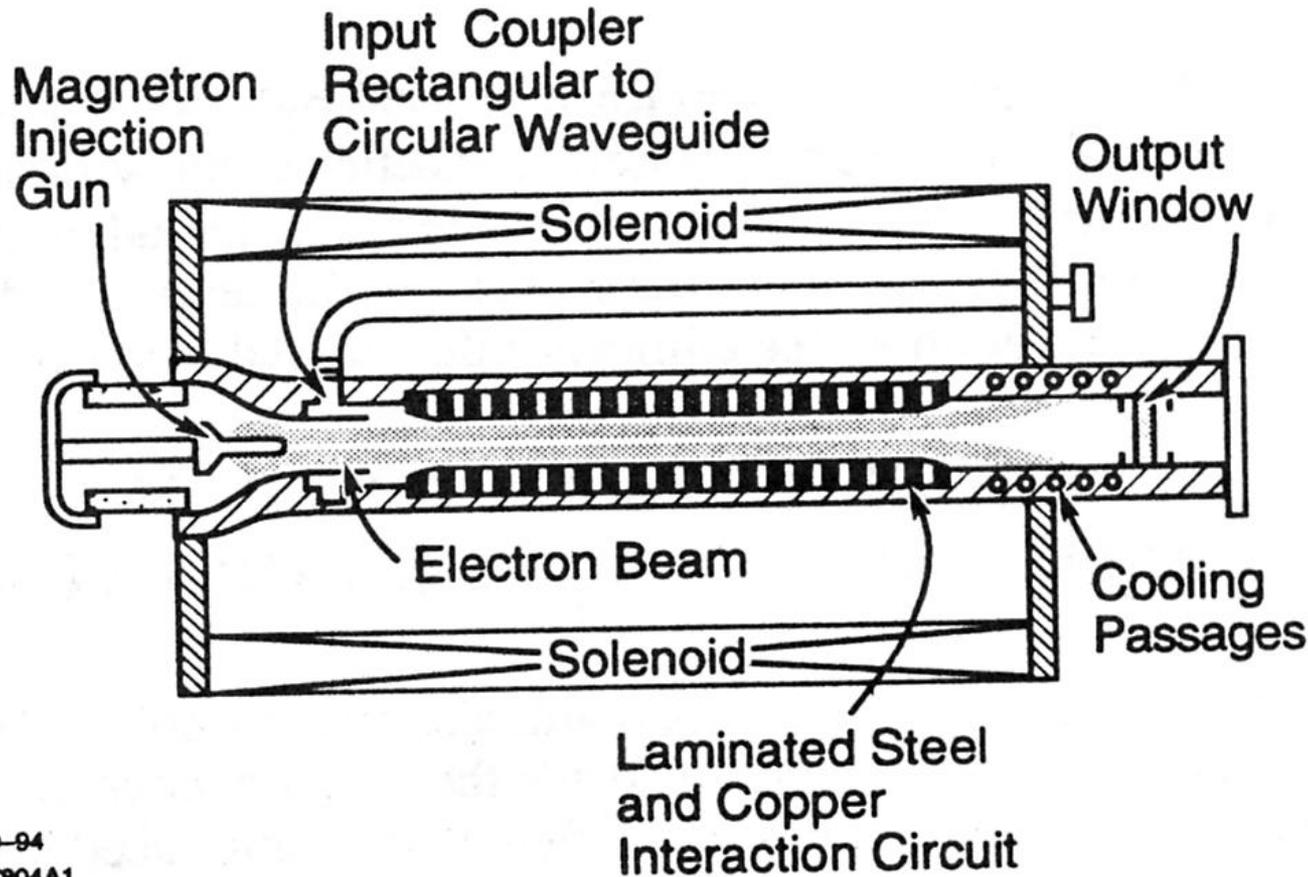


FIGURE 10. Computed steady-state evolution with z of radial orbit displacements (bottom) and particle energies (top). Also shown are the cavity outlines.

Transverse Deflection

Taken from O. Nezhevenko, *et al*, "High Power Pulsed Magnicon at 34-GHz", AIP Conf. Proc. 474, p. 195ff, (1998).

Ubitron or Free Electron Laser



9-94
7804A1

Taken from R. Phillips, "Conceptual Designs for NLC Ubitrons with Permanent-Magnet Wigglers", AIP Conf. Proc. 317, p.239, (1994).

LLNL 140 GHz FEL Amplifier

Measured Properties:

- Power: 1-2 GW (!)
@ 20 ns
- Efficiency: 14%
- Beam: 6 MV/2500A

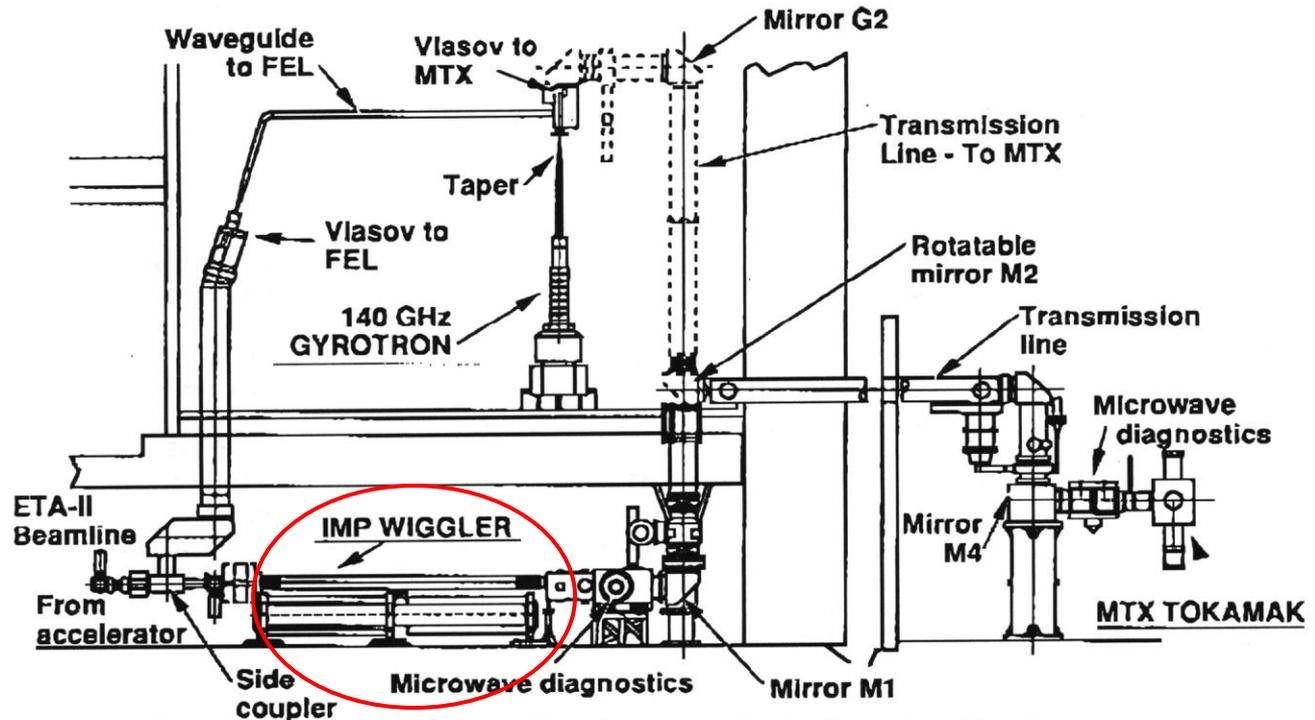


Fig. 1 The FEL is composed of the ETA-II accelerator (at left), the electron beamline, the master oscillator (gyrotron), and the quasi-optical microwave transport system to MTX.

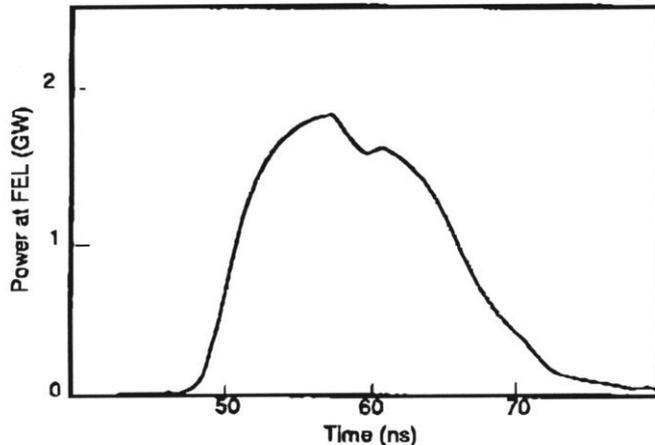


Fig. 2 A sample trace showing the output power versus time measured at the output of the FEL.

Taken from S. Allen, *et al*, "Generation of High Power 140 GHz Microwaves with an FEL for the MTX Experiment", in Proc. IEEE Part. Accel. Conf. (PAC 93), Washington, D.C., p.1551ff, (1993).

Compact Linear Collider (International collaboration centered at CERN)

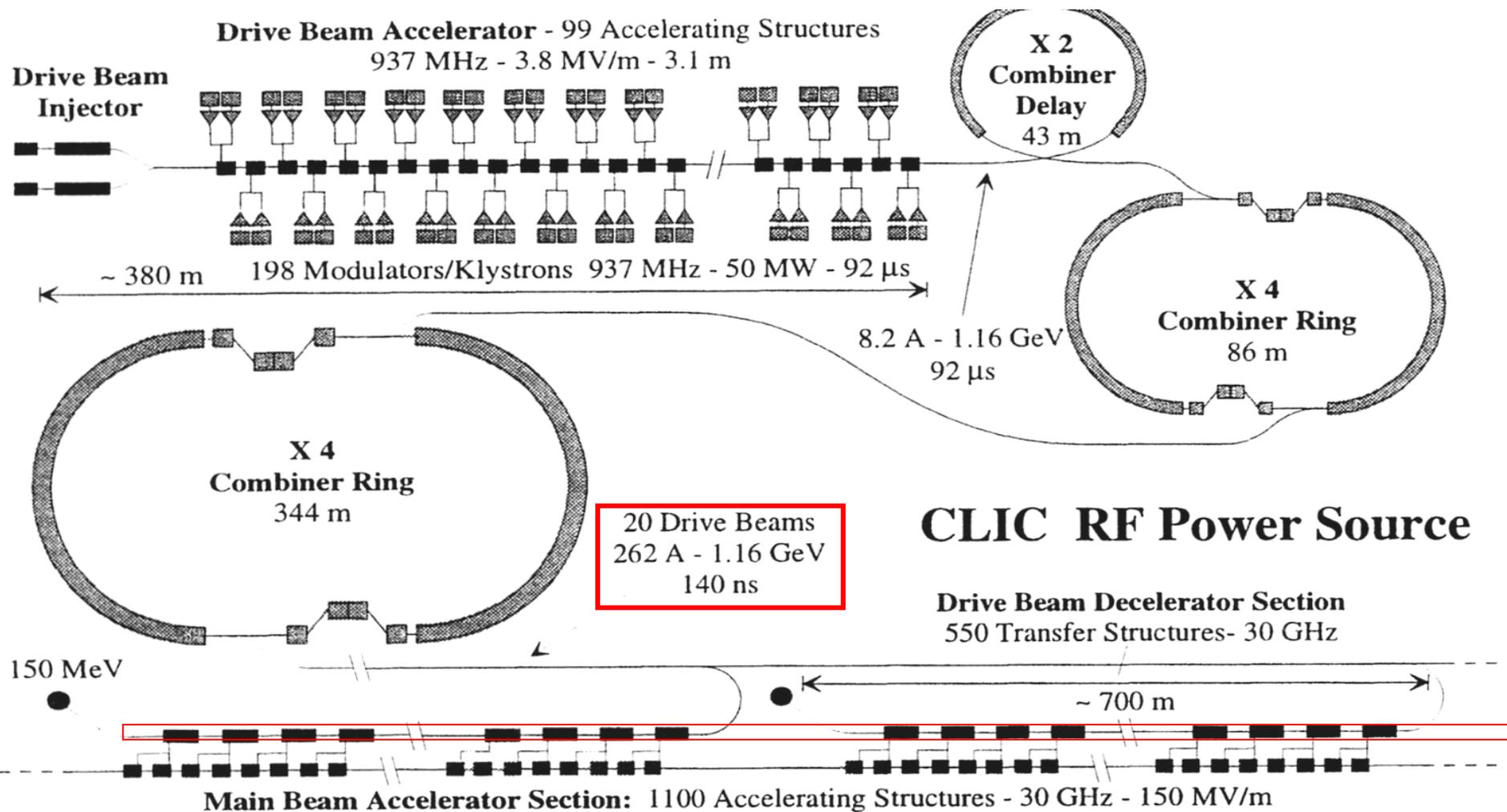


FIGURE 3. The 3 TeV RF Power Source.

Taken from H. Braun, *et al*, "A New Method for RF Power Generation for Two-Beam Linear Colliders", AIP Conf. Proc. 474, p.1ff, (1998).

CLIC Test Facility II

Inset photo taken from H. Braun, *et al*, “A New Method for RF Power Generation for Two-Beam Linear Colliders”, AIP Conf. Proc. 474, p. 1ff (1998).

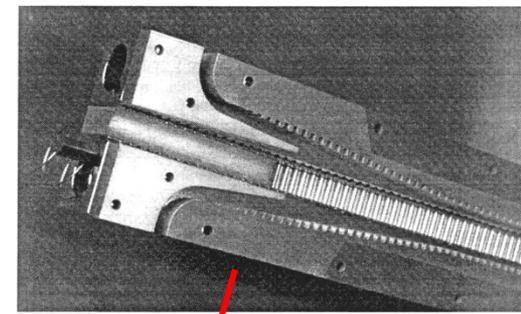


FIGURE 7. A four-channel decelerator structure.

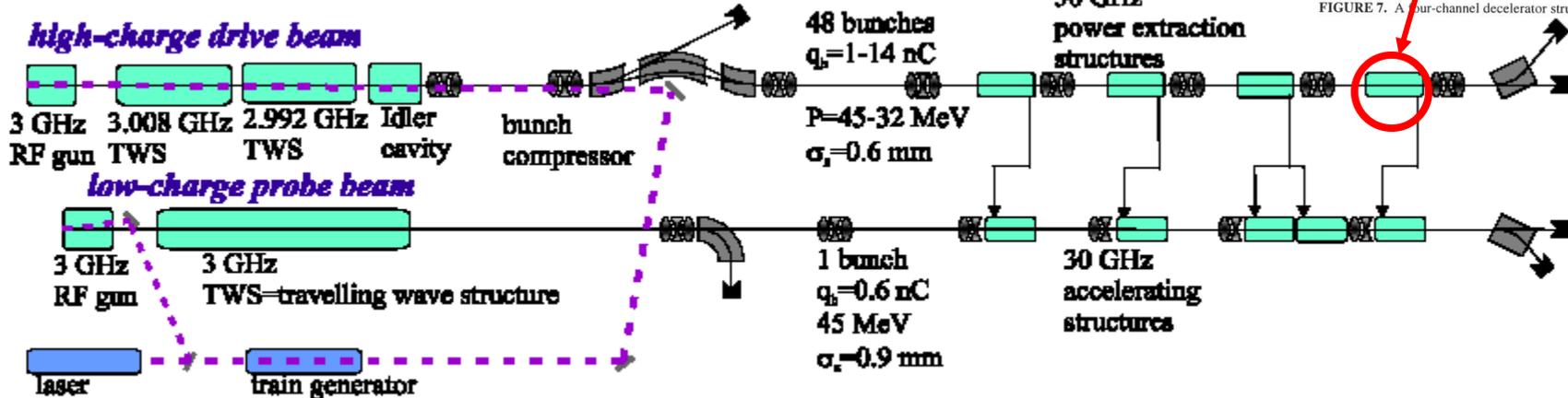


Figure 1: Configuration of CTF II during the 1999 run

Taken from: H. Braun, *et al*, “Experimental Results and Technical Research and Development at CTF II”, in Proc. Euro. Part. Accel. Conf., Vienna, Austria, p.48ff, (2000).

CTF-II Demonstrated:

- Generation and acceleration of 48x13.4 nC drive beam
- 120 MW of 30 GHz power generation in 16 ns pulses from the Power Extraction and Transfer Structures
- 59 MV/m acceleration of the probe beam

High Power Fiber Lasers

Analysis of the scalability of diffraction-limited fiber lasers and amplifiers to high average power

Jay W. Dawson, Michael J. Messerly, Raymond J. Beach, Miroslav Y. Shverdin, Eddy A. Stappaerts, Arun K. Sridharan, Paul H. Pax, John E. Heebner, Craig W. Siders and C.P.J. Barty

Lawrence Livermore National Laboratory, L-470, P.O. Box 808, Livermore, CA 94551
Corresponding author: rdawson17@llnl.gov

Abstract: We analyze the scalability of diffraction-limited fiber lasers considering thermal, non-linear, damage and pump coupling limits as well as fiber mode field diameter (MFD) restrictions. We derive new general relationships based upon practical considerations. Our analysis shows that if the fiber's MFD could be increased arbitrarily, 36 kW of power could be obtained with diffraction-limited quality from a fiber laser or amplifier. This power limit is determined by thermal and non-linear limits that combine to prevent further power scaling, irrespective of increases in mode size. However, limits to the scaling of the MFD may restrict fiber lasers to lower output powers.

©2008 Optical Society of America

OCIS codes: (140.3510) Lasers, fiber; (140.4480) Optical amplifiers

References and links

1. A. Tümmernann, T. Schreiber, F. Röser, A. Liem, S. Höfer, H. Zellmer, S. Nolte, and J. Limpert, "The renaissance and bright future of fibre lasers," *J. Phys. B* **38**, 681 (2005).
2. T. Simpson, F. Doff, P. Peterson, and A. Gavrielides, "Coherent combining of spectrally broadened fiber lasers," *Opt. Express* **15**, 11731-11740 (2007).
3. A. Shirakawa, K. Matsuo, and K. Ueda, "Fiber laser coherent array for power scaling, bandwidth narrowing, and coherent beam direction control," in Proc. of the SPIE: Fiber Lasers II: Technology, Systems, and Applications, **5709**, 165-174 (The International Society for Optical Engineering, 2005).
4. S. Chen, Y. Li, and K. Lu, "Branch arm filtered coherent combining of tunable fiber lasers," *Opt. Express* **13**, 7878-7883 (2005).
5. J. Bouillet, D. Sabourdy, A. D. Bertelemon, V. Kermène, D. Pagnoux, P. Roy, B. Dussardier, and W. Blanc, "Coherent combining in an Yb-doped double-core fiber laser," *Opt. Lett.* **30**, 1962-1964 (2005).
6. J. Xu, S. Zhao, S. Zhao, R. Hou, Y. Li, L. Shi, and S. Fang, "High power high brightness laser source for material processing through incoherent beam combination," Proc. of the SPIE **6722**, 67221D-1-5 (The International Society for Optical Engineering, 2007).
7. R. K. Huang, L. J. Missaggia, J. P. Donnelly, C. T. Harris, and G. W. Turner, "High-brightness slab-coupled optical waveguide laser arrays," *IEEE Photon. Tech. Lett.* **17**, 959-961 (2005).
8. S. J. Augst, J. K. Ranka, T. Y. Fan, and A. Sanchez, "Beam combining of ytterbium fiber amplifiers," *J. Opt. Soc. Am. B* **24**, 1707-1715 (2007).
9. T. H. Loftus, A. M. Thomas, P. R. Hoffman, M. Norsen, R. Roysse, A. Liu, and E. C. House, "Spectrally Beam-Combined Fiber Lasers for High-Average-Power Applications," *IEEE J. of Sel. Top. Quantum Electron.* **13**, 487-497 (2007).
10. D. Walton, S. Gray, J. Wang, M. Li, X. Chen, A. Liu, L. Zeeteno, and A. Crowley, "Kilowatt-Level, Narrow-Linewidth Capable Fibers and Lasers," *J. Lightwave Tech.* **24**, 4729-4749 (2006).
11. J. D. Hansryd, "Increase of the SBS threshold in a short highly nonlinear fiber by applying a temperature distribution," *J. Lightwave Tech.* **19**, 1691-1697 (2001).
12. J. C. Knight, "Photonic Crystal Fibers and Fiber Lasers," *J. Opt. Soc. Am. B* **24**, 1661-1668 (2007).
13. P. S. J. Russell, "Photonic-Crystal Fibers," *J. Lightwave Tech.* **24**, 4729-4749 (2006).
14. J. C. Baggott, T. M. Mouro, K. Furusawa, and D. J. Richardson, "Comparative study of large-mode hole and conventional fibers," *Opt. Lett.* **26**, 1045-1047 (2001).
15. Y. Jeong, J. Salun, D. Payne, and J. Nilsson, "Ytterbium-doped large-core fiber laser with 1.36 kW continuous-wave output power," *Opt. Express* **12**, 6088-6092 (2004).

#97711 - \$15.00 USD Received 20 Jun 2008; revised 1 Aug 2008; accepted 3 Aug 2008; published 13 Aug 2008
(C) 2008 OSA 18 August 2008 / Vol. 16, No. 17 / OPTICS EXPRESS 13240



Typical Specifications

Optical Parameters

| | Unit | YLR-1000* | YLR-3000 | YLR-5000 | YLR-10000 | YLR-20000 |
|--------------------------------|-----------|-----------|----------|----------|-----------|-----------|
| Nominal Output Power | W | 1000 | 3000 | 5000 | 10000 | 20000 |
| BPP after Feeding Fiber | mm * mrad | <2.0* | <2.0 | <4.0 | <5 | ~8 |
| BPP after Processing Fiber | mm * mrad | <4 | <4 | 6 | <8 | ~ |
| Feeding Fiber Core Diameter | um | 50 | 50 | 100 | 100 | 200 |
| Processing Fiber Core Diameter | um | 100 | 100 | 150 | 200 | ~ |

Electrical Parameters

| | Unit | YLR-1000 | YLR-3000 | YLR-5000 | YLR-10000 | YLR-20000 |
|---------------------------|------|---|----------|----------|-----------|-----------|
| Electrical Requirements | V AC | 360-528V, 3P+PE, 50/60Hz | | | | |
| Typical Power Consumption | kW | 3.5-4 | 10-12 | 17-20 | 35-40 | 70-80 |
| Standard Interfaces | | Ethernet, Digital I/O, Analog, PROFIBUS* **, DeviceNet* | | | | |
| Direct Modulation | kHz | 0.5 | 0.5 | 0.5 | 0.5 | 0.5 |

General Parameters

| | Unit | YLR-1000 | YLR-3000 | YLR-5000 | YLR-10000 | YLR-20000 |
|---------------------------------------|-------------------|-----------|-----------|------------|------------|------------|
| Max. Cooling Water Consumption (25°C) | m ³ /h | 0.6 | 1.4 | 2.5 | 4 | 7 |
| Cooling Water Temperature Range | °C | 20-25 | 20-25 | 20-25 | 20-25 | 20-25 |
| Dimensions (W x H x D) | cm | 86x120x81 | 86x120x81 | 110x150x81 | 110x150x81 | 210x150x81 |
| Weight | kg | 250 | 350 | 500 | 750 | 1200 |
| Ambient Temperature | °C | 10-50 | 10-50 | 10-50 | 10-50 | 10-50 |
| Enclosure Type | | IP54 | IP54 | IP54 | IP54 | IP54 |

* - 1 and 2kW lasers also available in with a single mode output versions YLR-1000SM and YLR-2000SM respectively. In this case BPP is ~0.4 mm * mrad from a 20um feeding fiber and ~2 mm * mrad from a 50um processing fiber. All other electrical and general parameters are the same as for YLR-1000 and YLR-2000 lasers.

** - optional, on request.



http://www.ipgphotonics.com/documents/documents/HP_Brochure.pdf

Carrier Phase-Locked Lasers

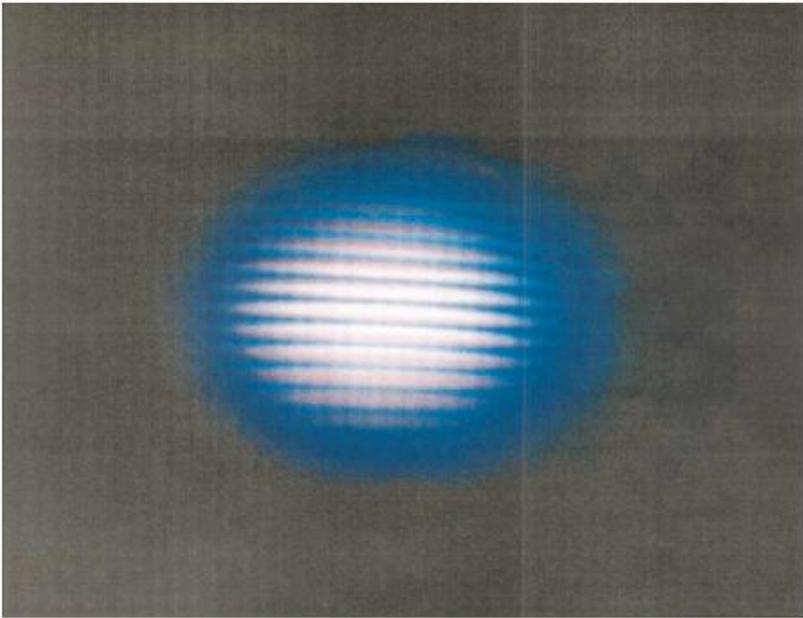


Fig. 2. White-light fringes resulting from the interference of the two continua generated by the two phase-locked IR laser pulses when the relative delay is properly adjusted to zero.

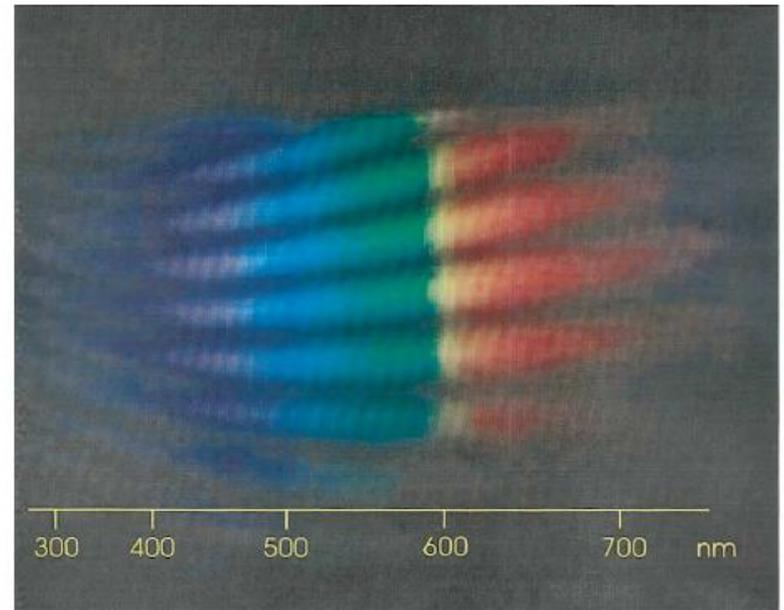


Fig. 3. Spectrally dispersed white-light fringes. Clear and well-defined fringes indicate that a stable phase relationship is conserved across all the generated visible spectrum.

Interference fringes of carrier phase-locked white light continua generated from Ti:Sapphire laser.

M. Bellini, T Hansch, *Optics Letters*, **25** (14), p.1049, (2000).

Sources by Power Density and Wavelength

