



Basics of photon detection and statistics

Jeff Corbett^a, Alan Fisher^a, Walter Mok^a, Weixing Cheng^b

^a - *SLAC National Accelerator Laboratory*

^b - *Brookhaven National Laboratory*

*Beam Diagnostics Using Synchrotron Radiation:
Theory and Practice*

US Particle Accelerator School
University of California, Santa Cruz
San Francisco — 2010 January 18 to 22



Outline

- Pulsed LED to simulate the synchrotron radiation light
- Photodiode used as sensor to detect the light
- Statistics basics



Synchrotron Radiation Pulse

- Several ps to 100ps pulse length (RMS) from most of third generation light sources
- Broadband from IR to x-ray
- Huge number of photons per pulse (depends on current, magnet field strength, aperture, filters etc)
- Spatial distribution (σ , π – mode from dipole)

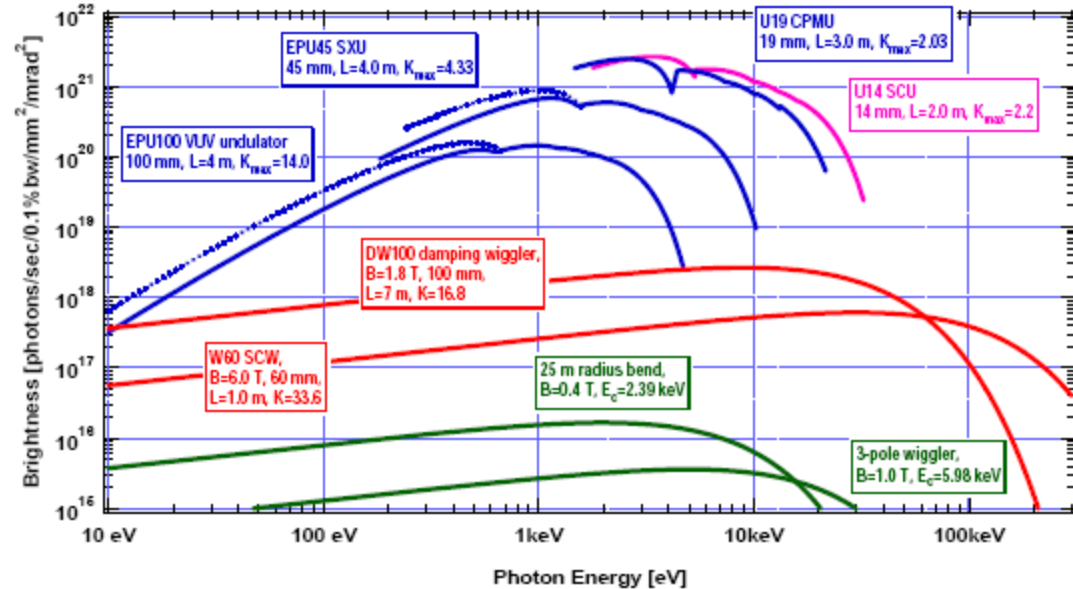


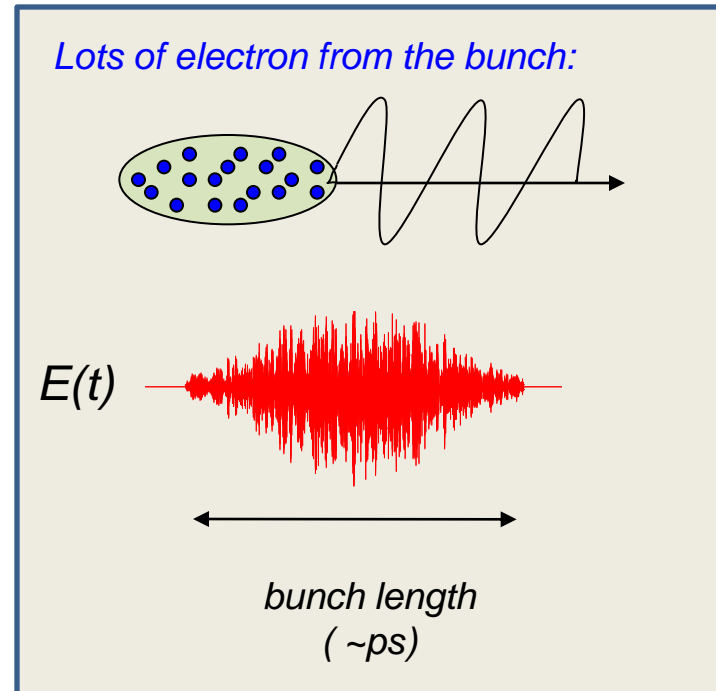
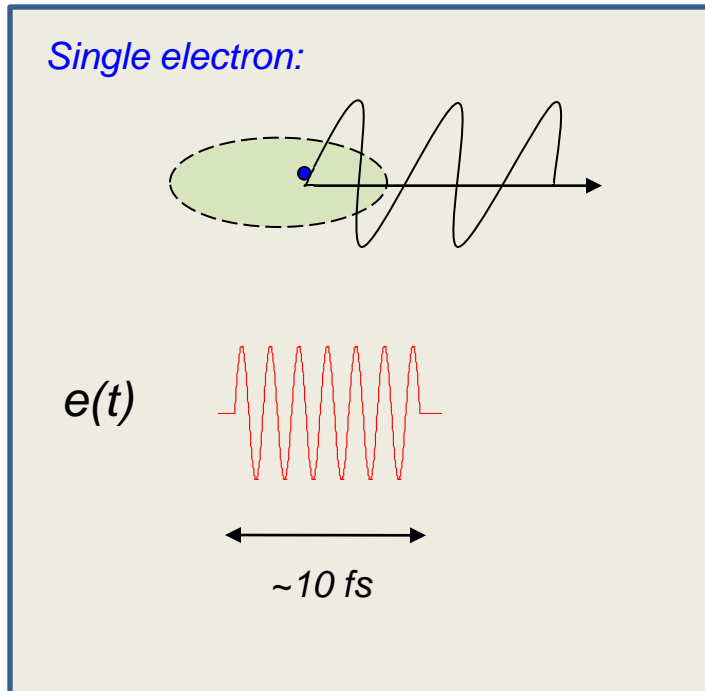
Figure 1. Brightness vs. photon energy for various devices at NSLS-II.

*These broadband spectrum are **average effect** of many electrons' radiation, or one electron radiation for many photons in long time scale.*



Radiation Theory (time-domain)

- One electron (j -th) in the bunch radiates an electromagnetic pulse $\mathbf{e}(t_j)$
- Total radiated field is
$$E(t) = \sum_{j=1}^N e(t - t_j)$$





Radiation Theory (frequency-domain)

$$\hat{E}(\omega) = \int_{-\infty}^{\infty} E(t) e^{i\omega t} dt = \hat{e}(\omega) \sum_{k=1}^N e^{i\omega t_k}$$

Radiation power per passage:

$$P(\omega) \propto |\hat{E}(\omega)|^2 = |\hat{e}(\omega)|^2 \sum_{k=1}^N \sum_{l=1}^N e^{i\omega(t_k - t_l)}$$

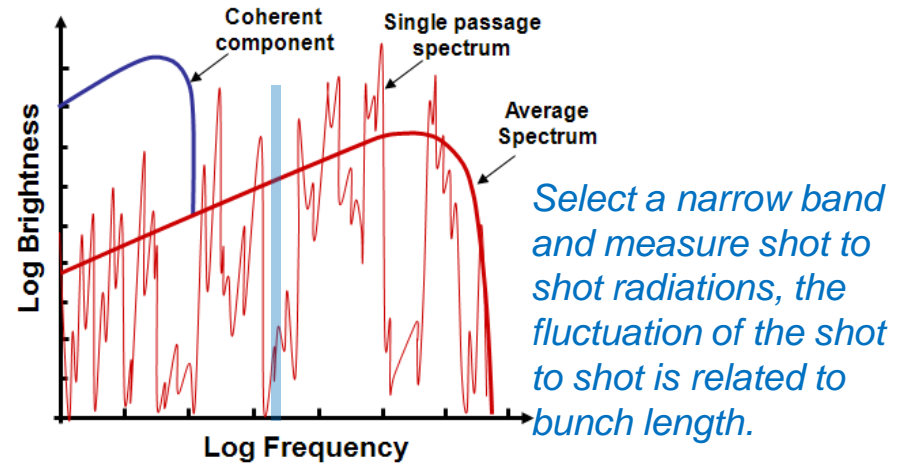
Average power:

$$\langle P(\omega) \rangle \propto |\hat{e}(\omega)|^2 \sum_{k,l=1}^N \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} dt_k dt_l f(t_k) f(t_l) e^{i\omega(t_k - t_l)} = |\hat{e}(\omega)|^2 \left[N + N(N-1) |\hat{f}(\omega)|^2 \right]$$

$f(t)$ is bunch distribution

$$f(t) = \frac{1}{\sqrt{2\pi}\sigma_z} \exp\left(-\frac{z^2}{2\sigma_z^2}\right)$$

$$\hat{f}(\omega) = \int_{-\infty}^{\infty} f(t) e^{i\omega t} dt = \exp\left(-2\pi^2 \left(\frac{\sigma_z}{\lambda}\right)^2\right)$$



Incoherent term

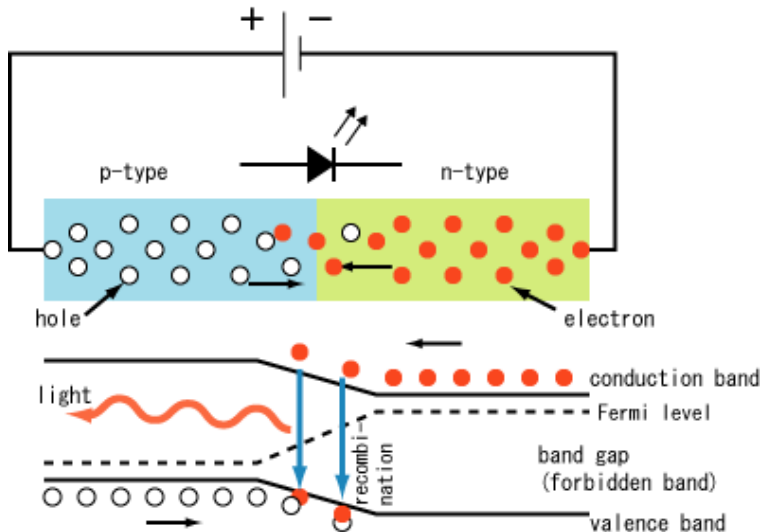
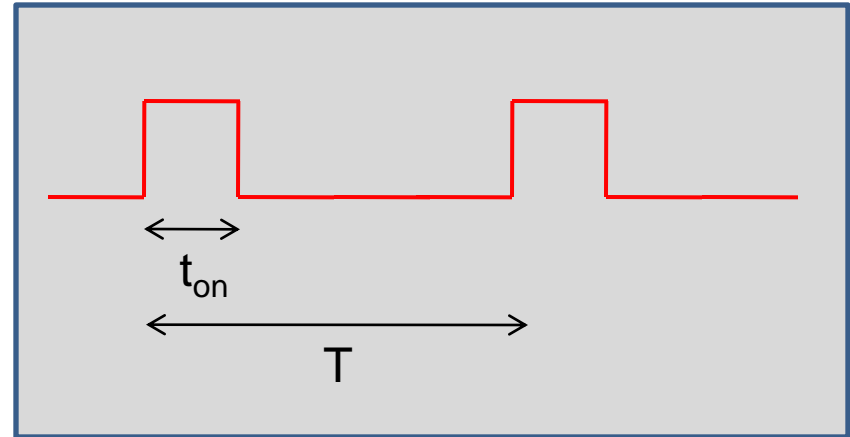
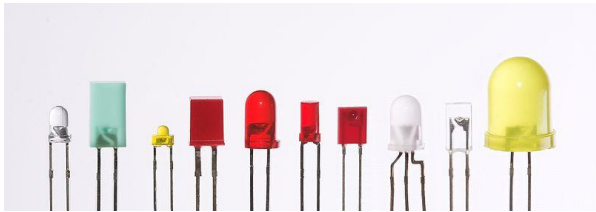
Coherent



Simulated pulse from LED

LED – Light Emitting Diode

- Semiconductor light source
- Visible, UV, IR wavelength
- High brightness



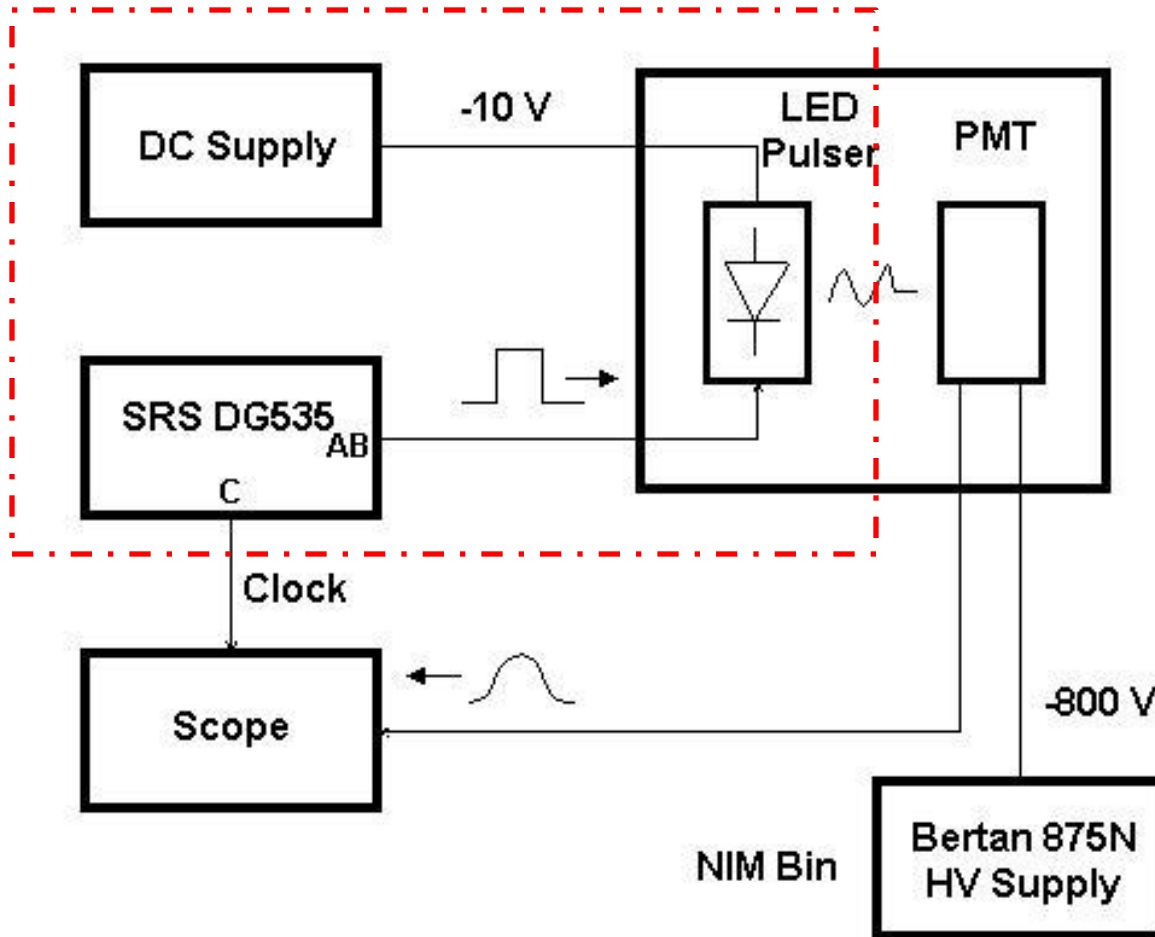
Pulse driver



- 50ps FWHM minimum laser pulse
- Repetition rate 31.25kHz to 80MHz
- Pulse energy adjustable
- Laser heads from 375nm to 1550nm
- LED heads from 255nm to 600nm
- External trigger / sync output



Pulse LED with DG535



Typical rise time 2-3ns,
hard to get pulse less than
1ns

With fast rise time option,
could get 100ps (20% -
80% amplitude)

Use ~10ns minimum due to
LED limitations



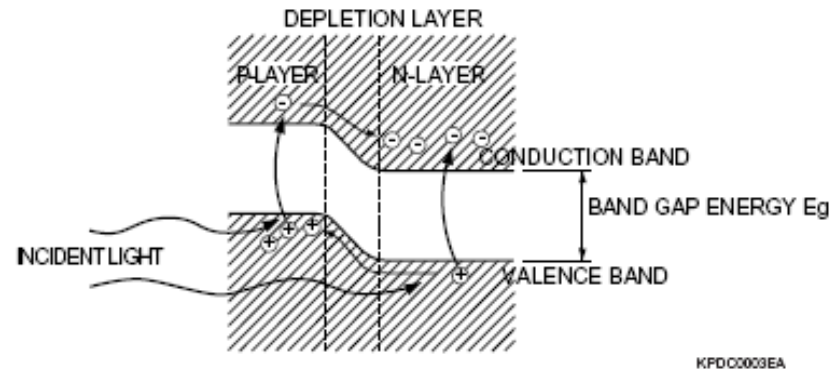
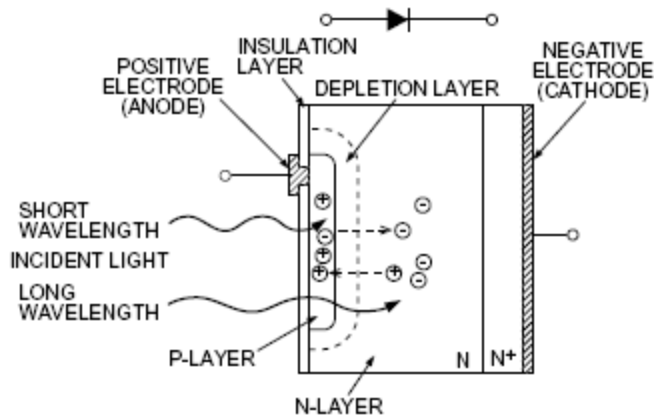
Photodiode

- Convert light to current/voltage signal
- Similar to regular semiconductor diode
- Photon hit the diode => create electron/hole pair
- Positive charge collected at P-layer, negative charge in the N-layer

LED symbol



Photodiode symbol





APD principle

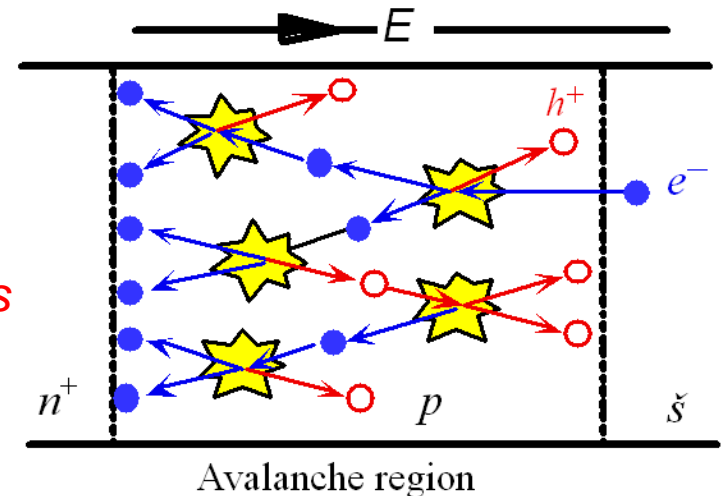
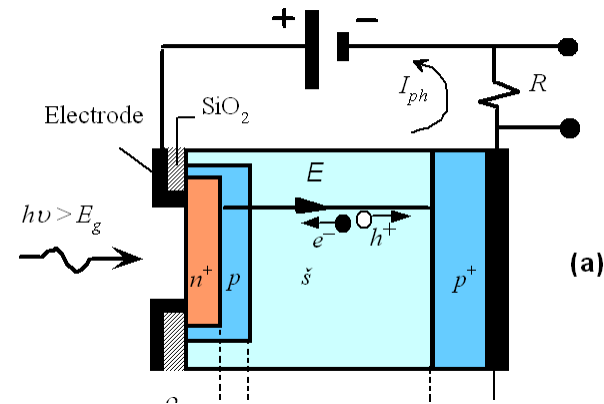
Diode operates in reverse bias mode

Electrons drift to N-layer with large kinetic energy, holes drift to P-layer
Impact ionization will generate more electron-hole pair

Avalanche of impact ionization

Internal gain (typical gain 100 ~1000)

A single electron entering the p-layer can generate a large number of electron-hole pairs which contribute to an observed photocurrent





SPAD (Single Photon Avalanche Diode)

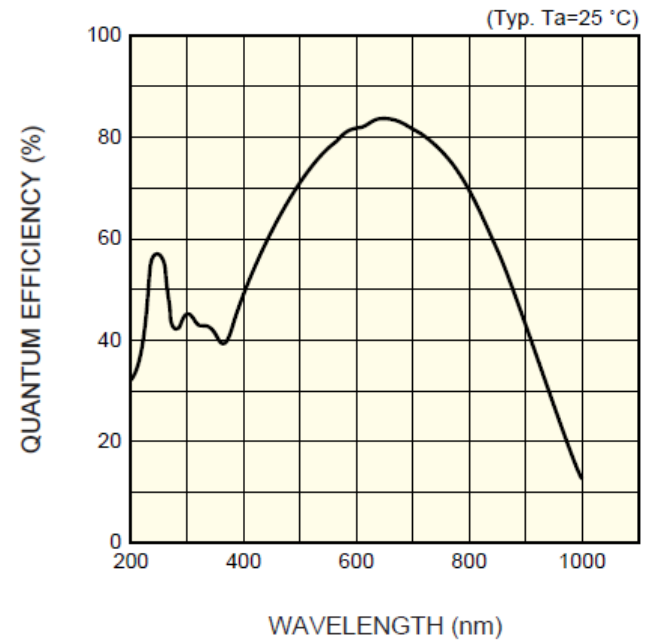
- SPAD is APD working in Geiger mode
- Reverse bias voltage well above the breakdown voltage. (APDs operate at a bias lesser than the breakdown voltage)
- Able to detect low intensity signals (down to the single photon)
- APD a single photon produces only tens or few hundreds of electrons, but in a SPAD a single photon triggers a current in the mA region (billions of billions of electrons per second) that can be easily "counted".
- APD is a linear amplifier for the input optical signal with limited gain (~hundreds), the SPAD is a trigger device so the gain concept is meaningless.
- SPAD used for bunch purity monitoring



APD Examples



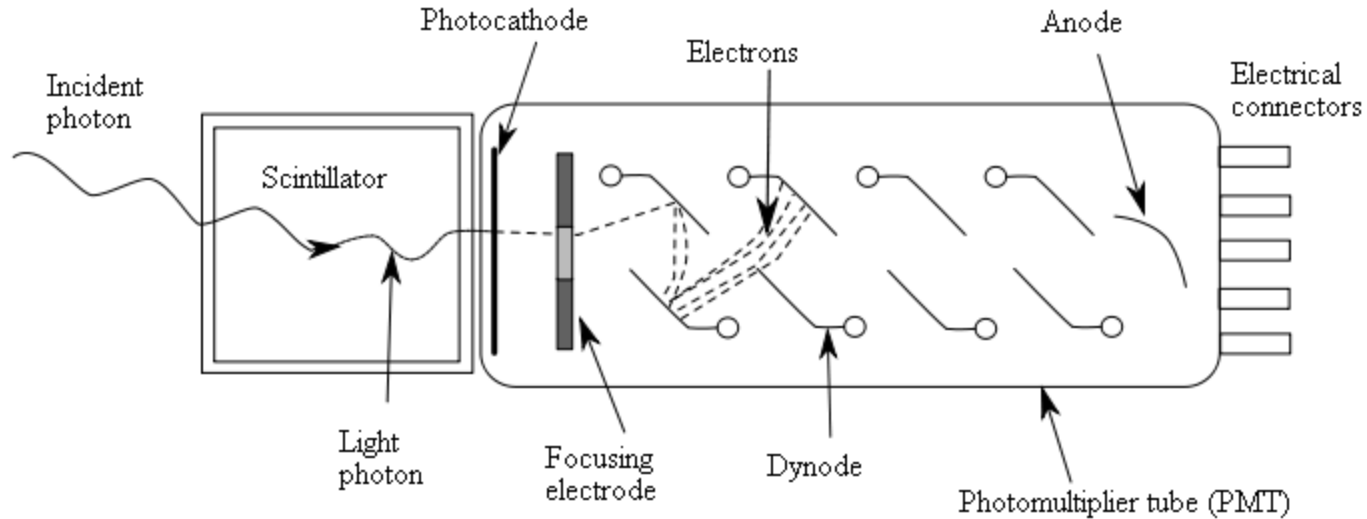
Hamamatsu Photonics S5343



200nm – 1000nm, peak at 620nm



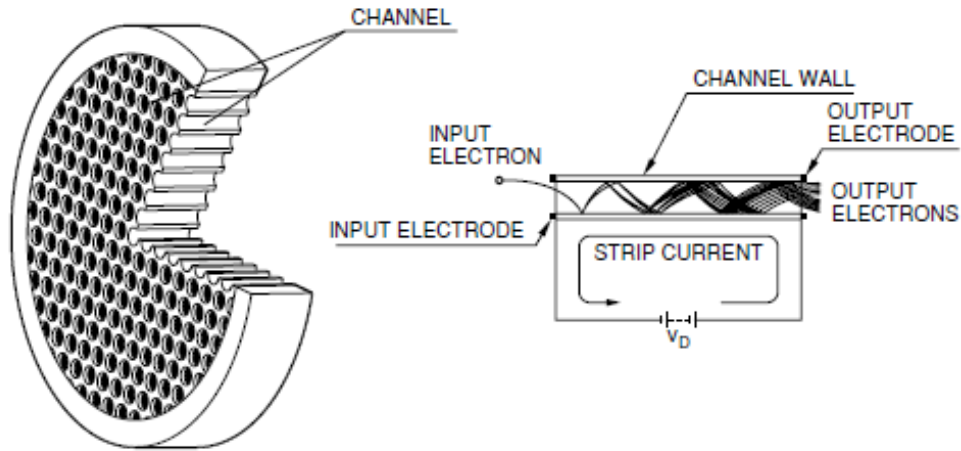
Photomultiplier tube



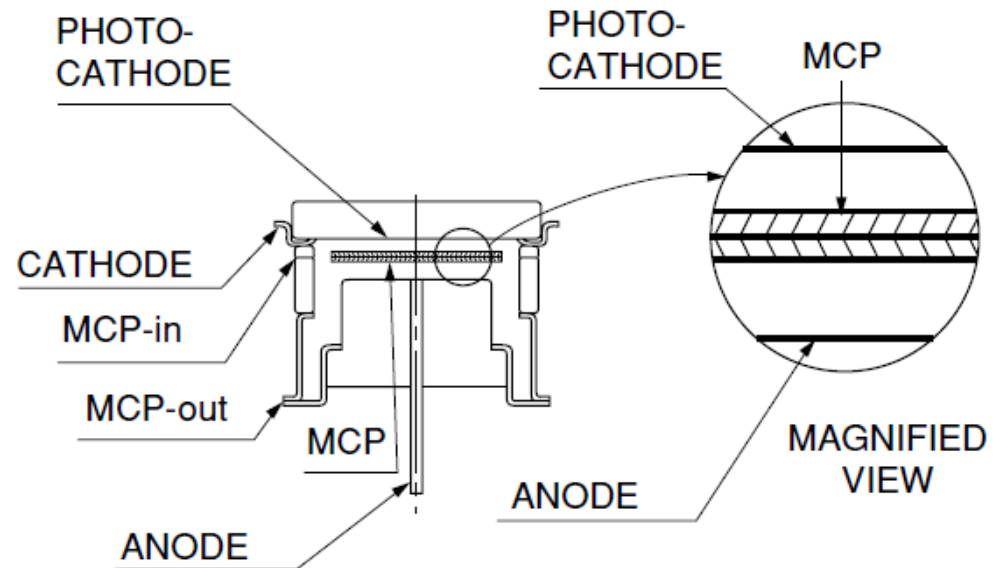
Photodiode	Photomultiplier
Good linearity	High gain ($\sim 10^8$)
Wide response spectral	Fast response time
Low noise	High sensitivity (detecting low level light, 30 photons per second)
Low cost	



MCP-PMT

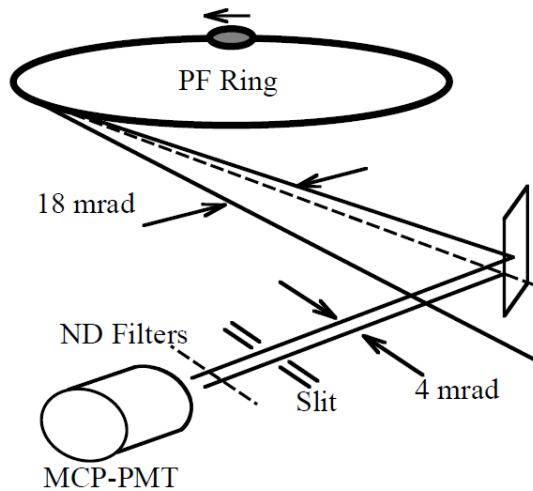


High gain
High spatial resolution
Fast time response



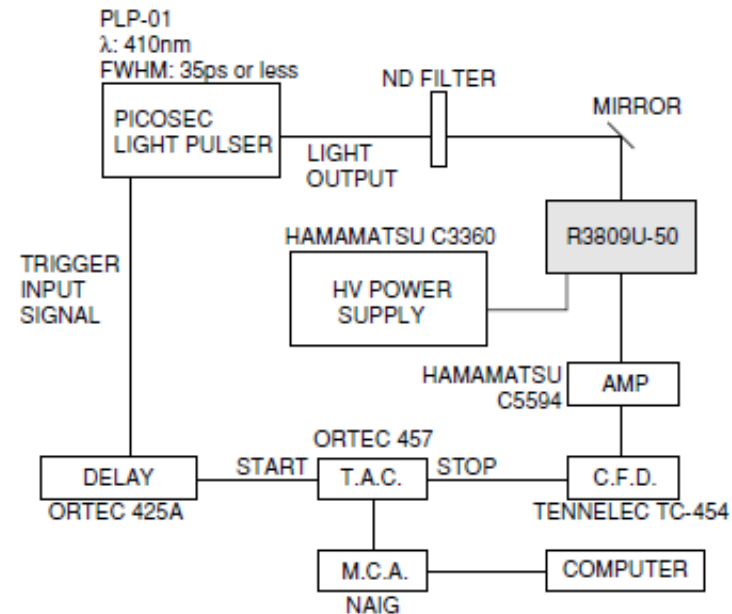


MCP-PMT used for bunch purity and filling pattern measurement



Hamamatsu R2809U-06 (or R3809U-52), used at PF, Spring-8, Diamond, ESRF etc. for the bunch purity and filling pattern monitor

Transit Time Spread = 25ps, 160nm to 650nm peak at 400nm

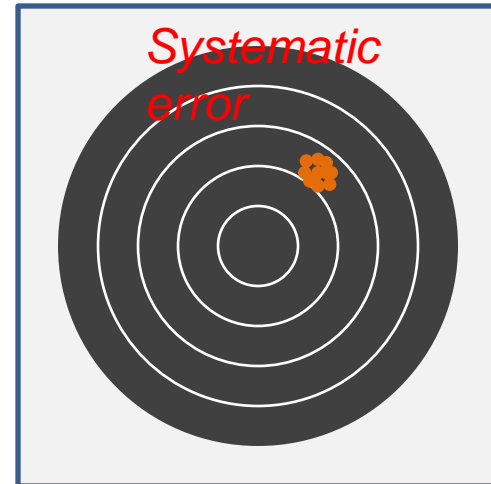
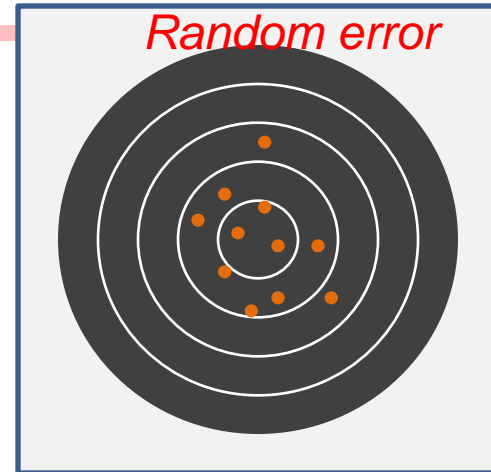
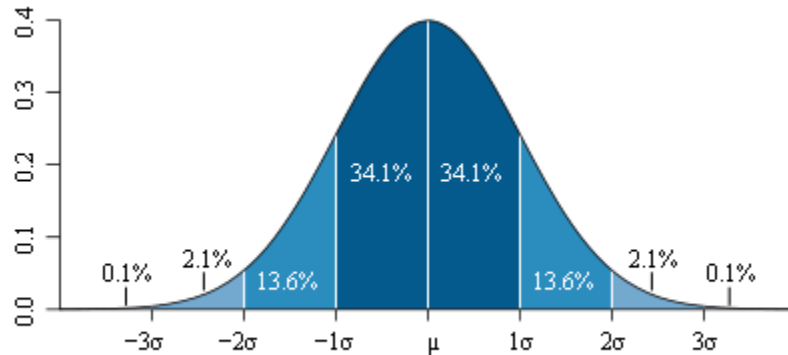


Time Correlated Single Photon Counting (TCSPC)



Statistic of sampled data

Gaussian distribution



N samples,

Mean:
$$\bar{x} = \frac{x_1 + x_2 + \dots + x_N}{N} = \frac{1}{N} \sum_{i=1}^N x_i$$

Std:
$$\sigma = \sqrt{\frac{1}{N} \sum_{i=1}^N (x_i - \bar{x})^2}$$

$$\sigma = \sqrt{\frac{1}{N} \left(\left(\sum_{i=1}^N x_i^2 \right) - N\bar{x}^2 \right)} = \sqrt{\frac{1}{N} \left(\sum_{i=1}^N x_i^2 \right) - \bar{x}^2}$$

Std of the mean: $\sigma_{\text{mean}} = \frac{\sigma}{\sqrt{N}}$ N – number of samples

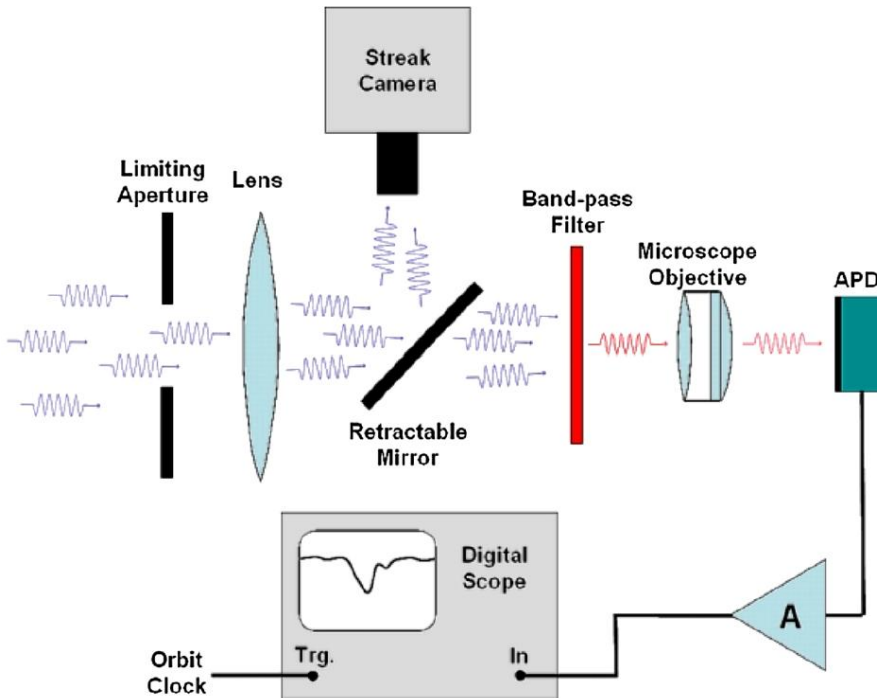
$$\sigma_{\text{mean}} = \frac{\sigma}{\sqrt{N}}$$

$$\delta x_{\text{rand}} = \sigma_{\text{mean}}$$

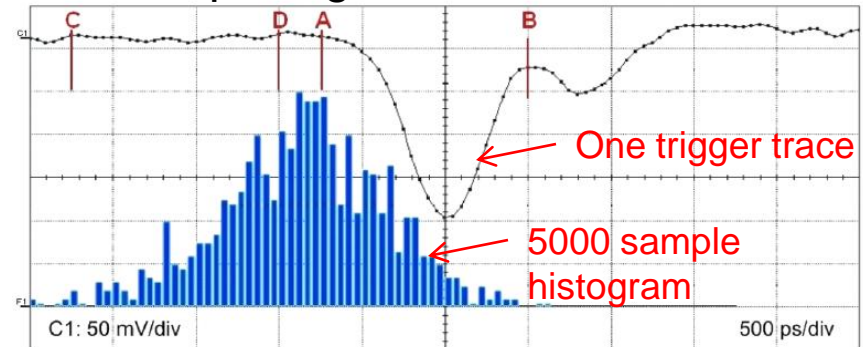
$$\delta x_{\text{total}} = \sqrt{(\delta x_{\text{rand}})^2 + (\delta x_{\text{sys}})^2}$$



Statistics in the fluctuation bunch length measurement



APD output signal



$$\sigma_{\delta^2} / \delta^2 = \sqrt{2 / (N_s - 1)}$$

N_s – number of sample, > 5000
to get <2% measurement error of δ^2

ALS Fluctuation Bunch Length
Measurement
(1nm FWHM band pass filter at 632.8nm)

$$\delta^2 \approx \frac{1}{\sqrt{1 + 4\sigma_\omega^2 \sigma_t^2}}$$

$$\delta^2 = \frac{\sigma_W^2}{\langle W \rangle^2}$$

W – radiation energy of per passage,
represented by APD pulse amplitude or area