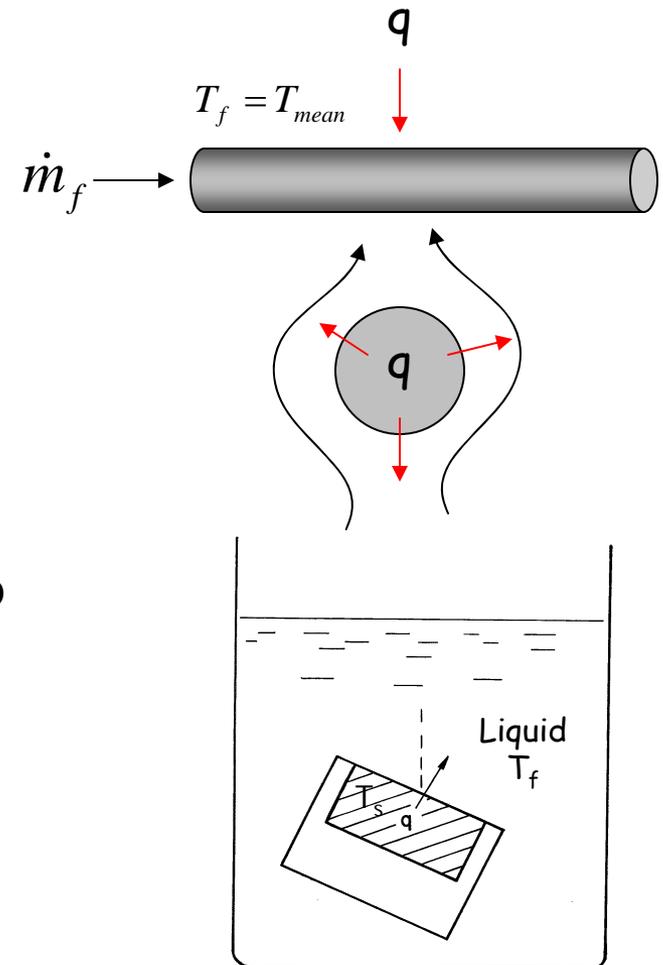


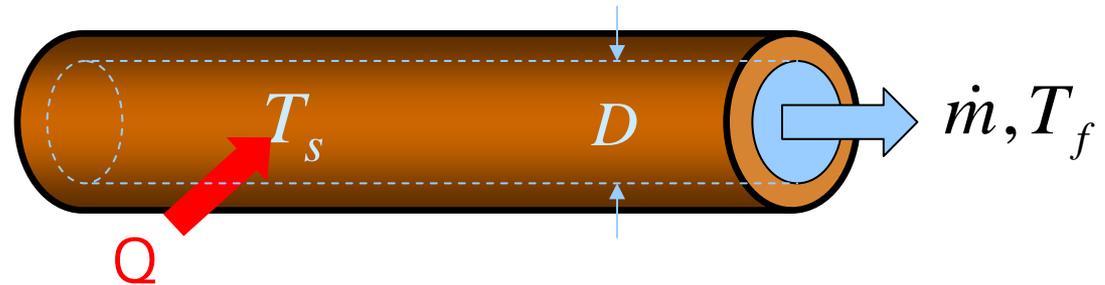
3.2 Cryogenic Convection Heat Transfer

- Involves process of heat transfer between solid material and adjacent cryogenic fluid
- Classic heat transfer problem (Newton's law)
$$q(\text{kW/m}^2) = h (T_s - T_f)$$
- Configurations of interest
 - Internal forced flow (single phase, $T_f = T_{mean}$)
 - Free convection (single phase, $T_f = T_\infty$)
 - Internal two phase flow
 - Pool boiling (two phase)
- Understanding is primarily empirical leading to correlations based on dimensionless numbers
- Issue is relevant to the design of:
 - Heat exchangers
 - Cryogenic fluid storage
 - Superconducting magnets
 - Low temperature instrumentation



Single phase internal flow heat transfer

Forced
Convection



Classical fluid correlations

- The heat transfer coefficient in a classical fluid system is generally correlated in the form where the Nusselt number, where, $Nu_D \equiv hD/k_f$ and D is the characteristic length
- For laminar flow, $Nu_D = \text{constant} \sim 4$ (depending on b.c.)
- For turbulent flow ($Re_D > 2000$)

$$Nu_D = f(Re_D, Pr) = C Re_D^n Pr^m \quad \text{and} \quad Pr \equiv \frac{\mu_f C_p}{k_f} \quad (\text{Prandtl number})$$

- Dittus-Boelter Correlation for classical fluids (+/- 15%)

$$Nu_D = 0.023 Re_D^{4/5} Pr^{2/5}$$

Note that fluid properties should be computed at T_f (the "film temperature"):

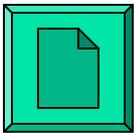
$$T_f \equiv \frac{T_s + T_f}{2}$$

Johannes Correlation (1972)

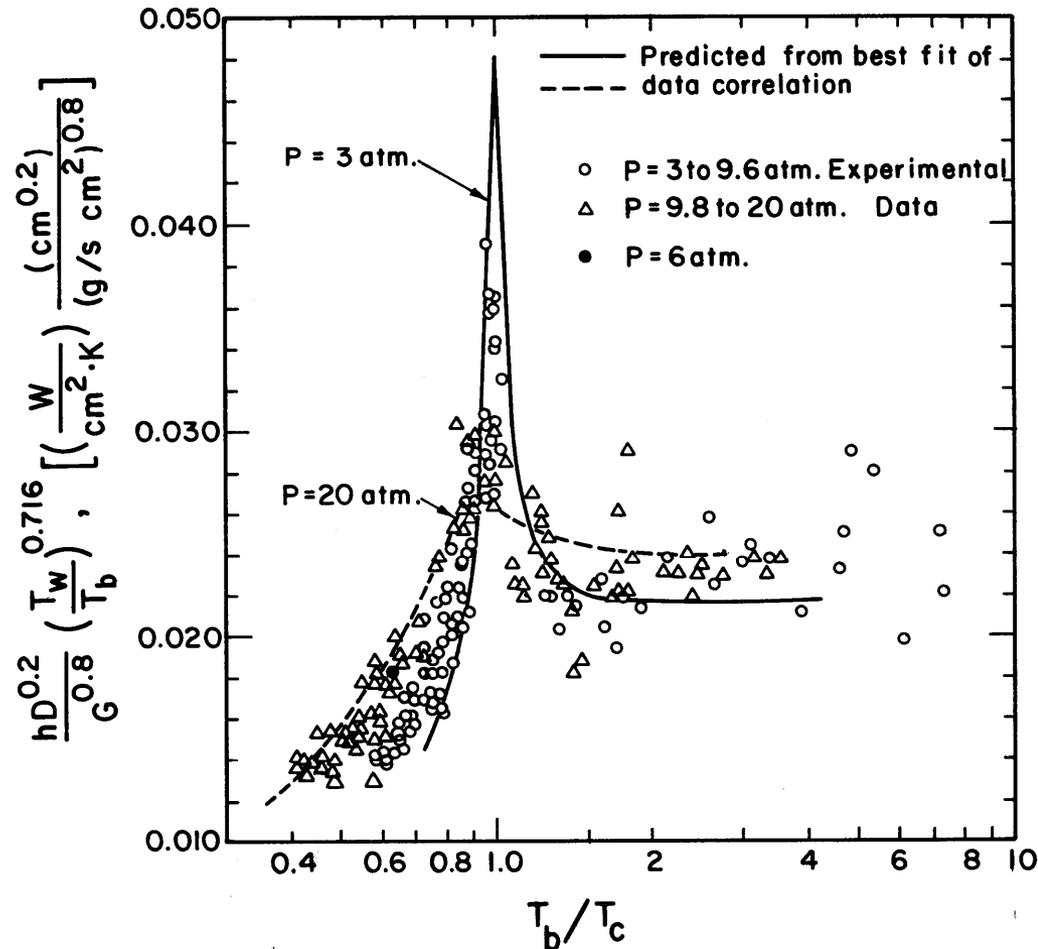
- Improved correlation specifically for helium (+/- 8.3%)

$$Nu_D = 0.0259 Re_D^{4/5} Pr^{2/5} \left(\frac{T_s}{T_f} \right)^{-0.716}$$

- Last factor takes care of temperature dependent properties
- Note that one often does not know T_f , so iteration may be necessary.

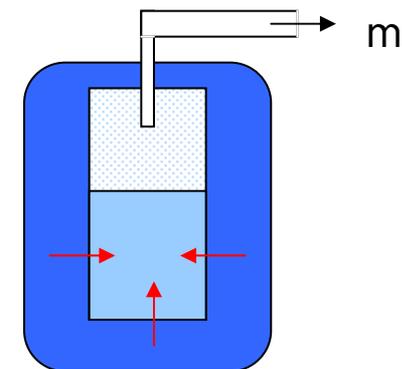
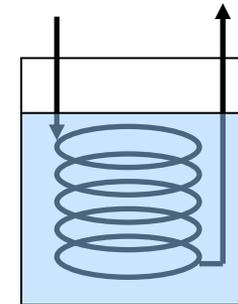
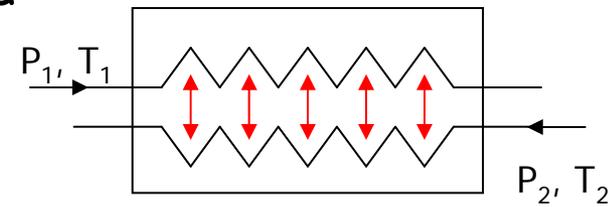


Example

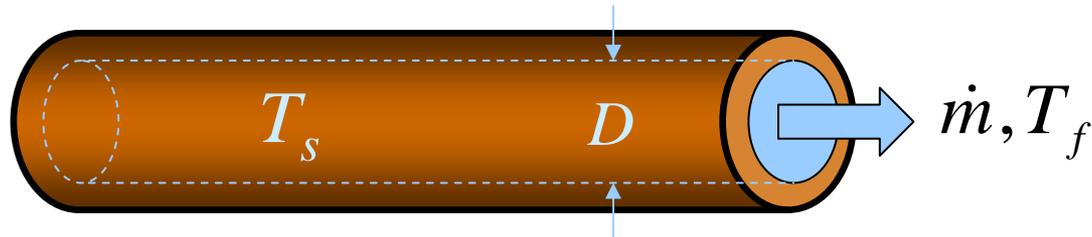


Application: Cryogenic heat exchangers

- Common types of heat exchangers used in cryogenic systems
 - Forced flow single phase fluid-fluid
 - E.g. counterflow heat exchanger in refrigerator/liquefier
 - Forced single phase flow - boiling liquid (Tube in shell HX)
 - E.g. LN₂ precooler in a cooling circuit
 - Static boiling liquid-liquid
 - E.g. Liquid subcooler in a magnet system



Simple 1-D heat exchanger

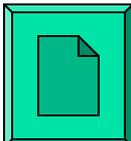
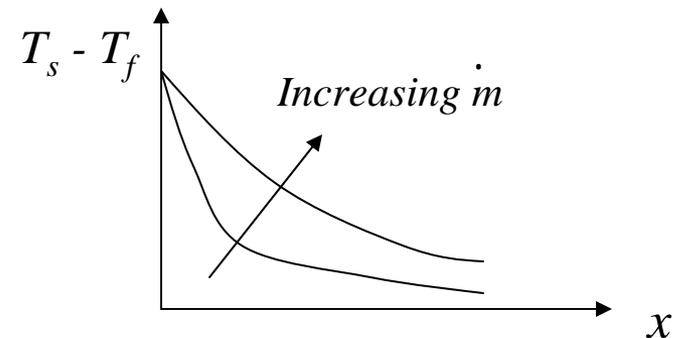


- Differential equation describing the temperature of the fluid in the tube:

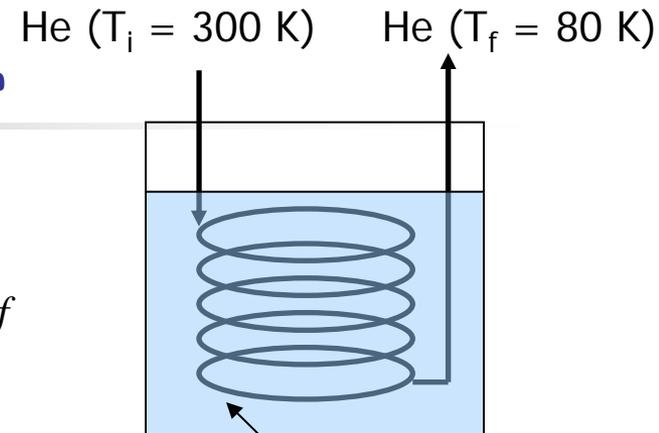
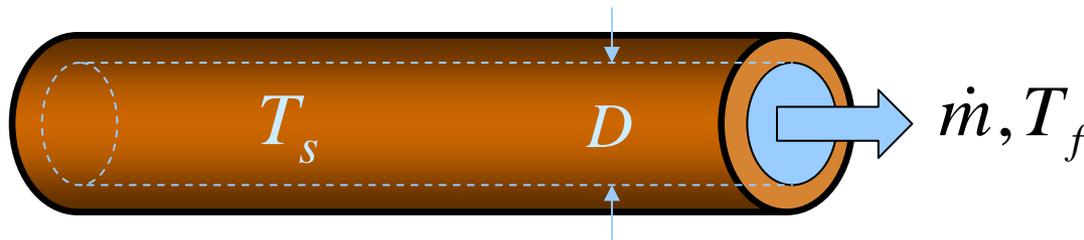
$$\dot{m}C \frac{dT_f}{dx} + hP(T_f - T_s) = 0$$

- For constant T_s , the solution of this equation is an exponential

$$T_s - T_f = (T_s - T_f)_0 \exp\left(-\frac{hP}{\dot{m}C} x\right)$$



Liquid nitrogen precooler



Assumptions & givens

- T_s is a constant @ 77 K (NBP of LN₂)
- Helium gas ($C_p = 5.2$ kJ/kg K; $\mu = 15 \times 10^{-6}$ Pa s; $\rho = 0.3$ kg/m³, $k = 0.1$ W/m K)
- Allowed pressure drop, $\Delta p = 10$ kPa
- Helium mass flow rate = 1 g/s

Properties are average values between 300 K and 80 K

Find the length and diameter of the HX (copper tubing)

Total heat transfer:

$$Q = \dot{m}C_p(T_{in} - T_{out}) = 1 \text{ g/s} \times 5.2 \text{ J/g K} \times 220 \text{ K} = 1144 \text{ W}$$

Log mean ΔT :

$$\Delta T_{lm} = \frac{\Delta T_f(x=0) - \Delta T_f(x=L)}{\ln\left(\frac{\Delta T_f(x=0)}{\Delta T_f(x=L)}\right)} = \frac{(300-77) - (80-77) \text{ K}}{\ln\left[\frac{300-77}{80-77}\right]} = 51 \text{ K}$$

Liquid nitrogen precooler (continued)

$$UA = h\pi DL = \frac{Q}{\Delta T_{lm}} = 1144 \text{ W/51 K} = 22.4 \text{ W/K}$$

- The heat transfer coefficient is a function of Re_D and $Pr = 0.67$
- Assuming the flow is turbulent and fully developed, use the Dittus Boelter correlation

$$Nu_D = \frac{hD}{k_f} = 0.023 Re_D^{0.8} Pr^{0.3} \quad \text{and} \quad Re_D = \frac{4\dot{m}}{\pi D \mu}$$

Substituting the Re_D and solving for h

$$h = 0.023 \frac{k_f}{D} \left(\frac{4\dot{m}}{\pi D \mu} \right)^{0.8} Pr^{0.3} = \frac{0.0247 \times 0.1 \text{ W/m K} \times (10^{-3} \text{ kg/s})^{0.8}}{(15 \times 10^{-6} \text{ Pa s})^{0.8} \times D^{1.8}}$$
$$= 0.07/D(\text{m})^{1.8}$$

$$h\pi DL = 0.224 \times (L/D^{0.8}) = 22.4 \text{ W/K}$$

or $L/D^{0.8} = 100 \text{ m}^{0.2}$ ←

1 equation for two unknowns

Liquid nitrogen precooler (continued)

- Pressure drop equation provides the other equation for L & D

$$\Delta p = f \frac{L}{2\rho D} \left(\frac{\dot{m}}{A_{flow}} \right)^2 \quad ; \quad A_{flow} = \pi D^2/4 \text{ and } f \sim 0.02 \text{ (guess)}$$

$$Re_D = \frac{4\dot{m}}{\pi D \mu}$$

$$\frac{4 \times 0.001 \text{ kg/s}}{\pi \times 0.0062 \text{ m} \times 15 \times 10^{-6} \text{ Pa s}}$$

Substituting for Re_D and f

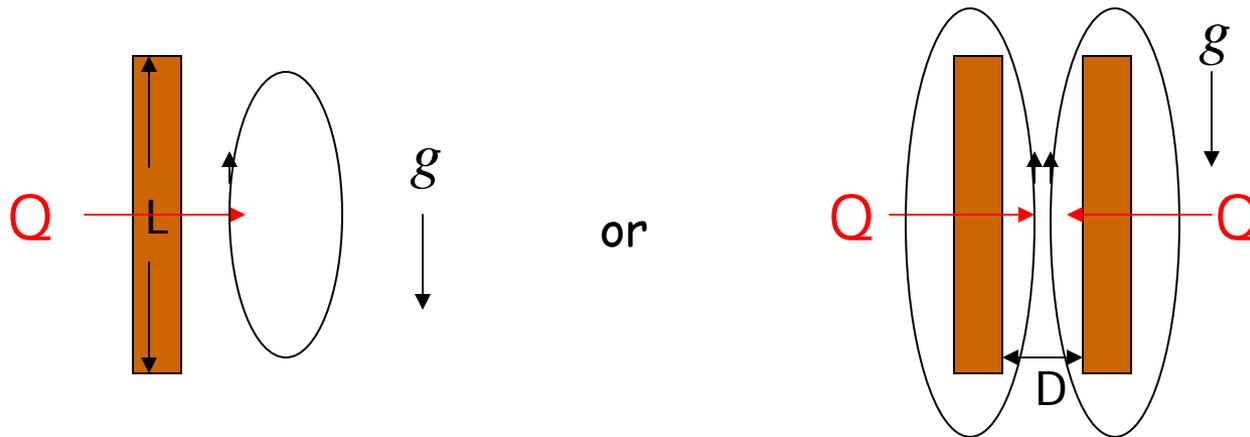
$$\Delta p \approx 0.016 \frac{\dot{m}^2}{\rho} \frac{L}{D^5} = \frac{0.016 \times (10^{-3} \text{ kg/s})^2}{0.3 \text{ kg/m}^3} \left(\frac{L}{D^5} \right) \quad \sim 13,700$$

Substitute: $\Delta p = 5.33 \times 10^{-8} \left(\frac{L}{D^5} \right)$ ← 2nd equation for two unknowns

Eq. 1: $L = 100 D^{0.8} \rightarrow \Delta p = 5.33 \times 10^{-6}/D^{4.2}$ with $\Delta p = 10,000 \text{ Pa}$

$D = [5.33 \times 10^{-6}/\Delta p]^{1/4.2} = 6.2 \text{ mm}$ and $L = 100 \times (0.0062 \text{ m})^{0.8} = 1.7 \text{ m}$

Single phase free convection heat transfer



- Compressible fluid effect: Heat transfer warms the fluid near the heated surface, reducing density and generating convective flow.
- Free convection heat transfer is correlated in terms of the Rayleigh number,

$$Nu_L = f(Gr, Pr) \sim CRa_L^n \quad \text{where} \quad Ra_L \equiv Gr Pr = \frac{g\beta\Delta TL^3}{D_{th}\nu}$$

where g is the acceleration of gravity, β is the bulk expansivity, ν is the kinematic viscosity (μ/ρ) and D_{th} is the thermal diffusivity ($k/\rho C$). L (or D) are the scale length of the problem (in the direction of \vec{g})

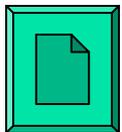
Free convection correlations

- For very low Rayleigh number, $Nu = 1$ corresponding to pure conduction heat transfer
- For $Ra < 10^9$, the boundary layer flow is laminar (conv. fluids)

$$Nu_L \approx 0.59 Ra_L^{0.25}$$

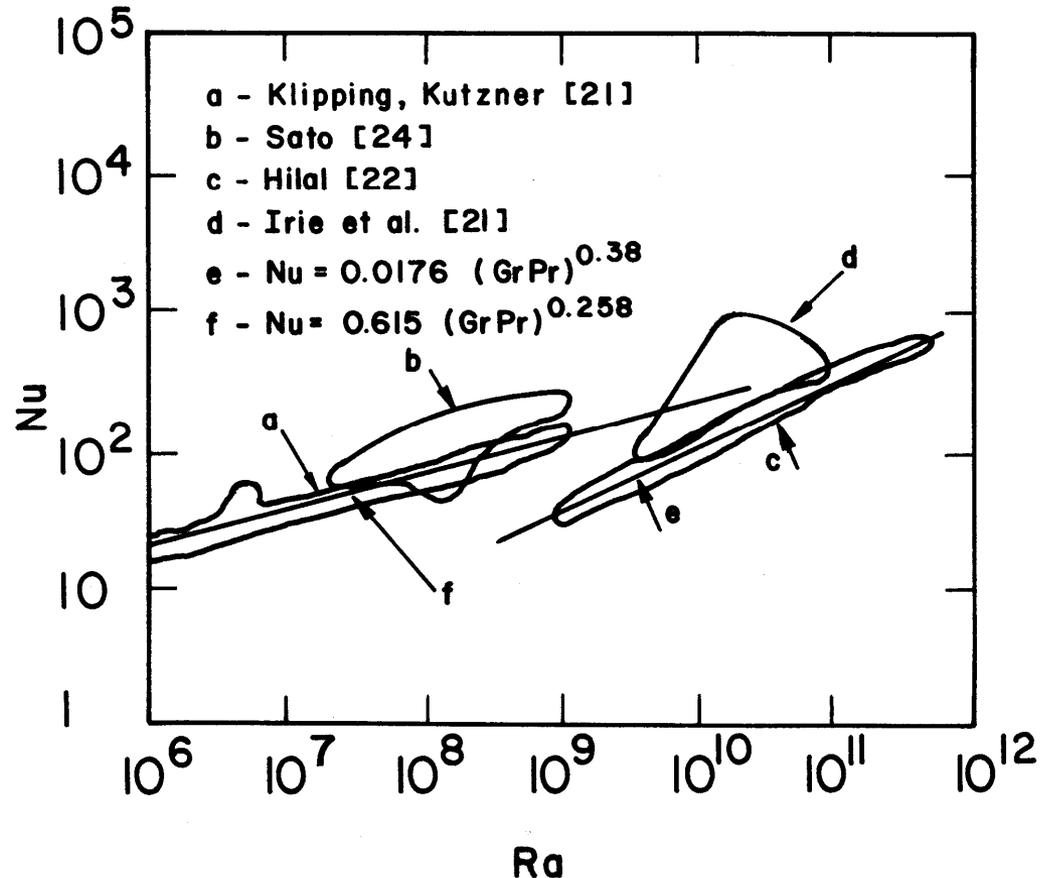
- For $Ra > 10^9$, the boundary layer is turbulent

$$Nu_L \approx 0.1 Ra_L^{0.33}$$



Example

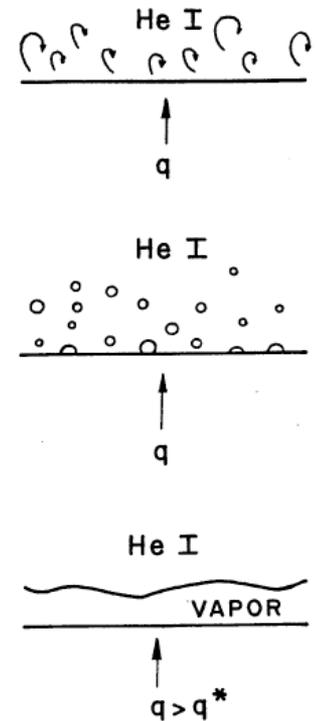
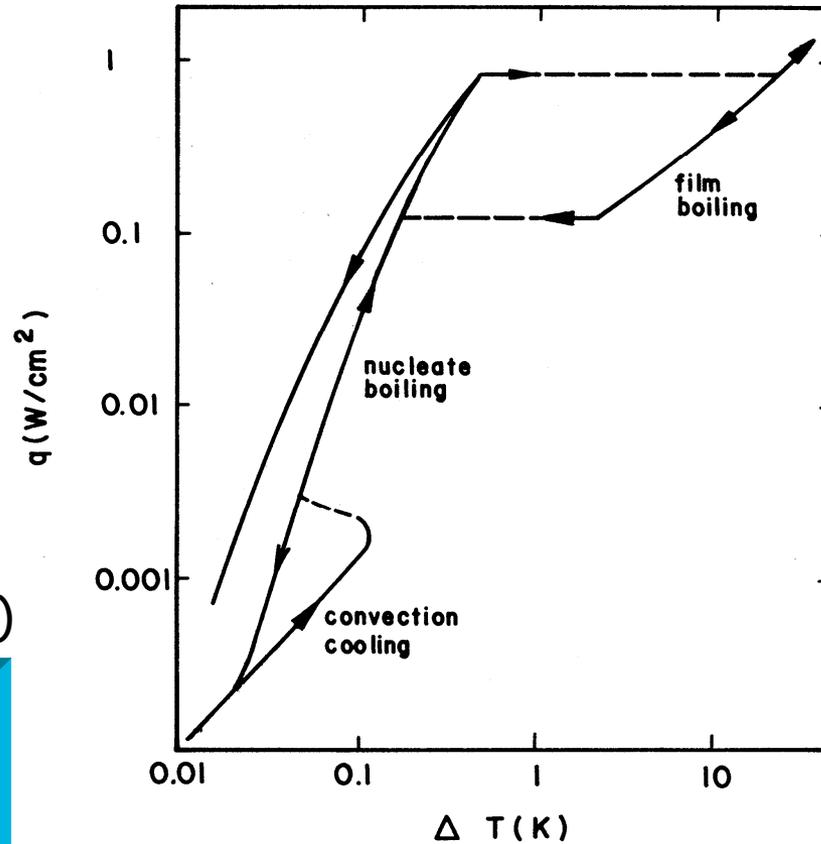
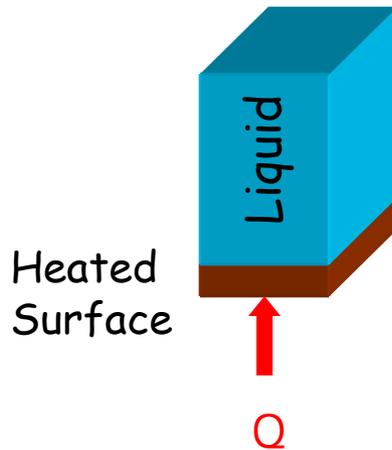
Free convection correlation for low temperature helium



Pool Boiling Heat Transfer (e.g. Helium)

Factors affecting heat transfer curve

- Surface condition (roughness, insulators, oxidation)
- Orientation
- Channels (circulation)
- Time to develop steady state (transient heating)



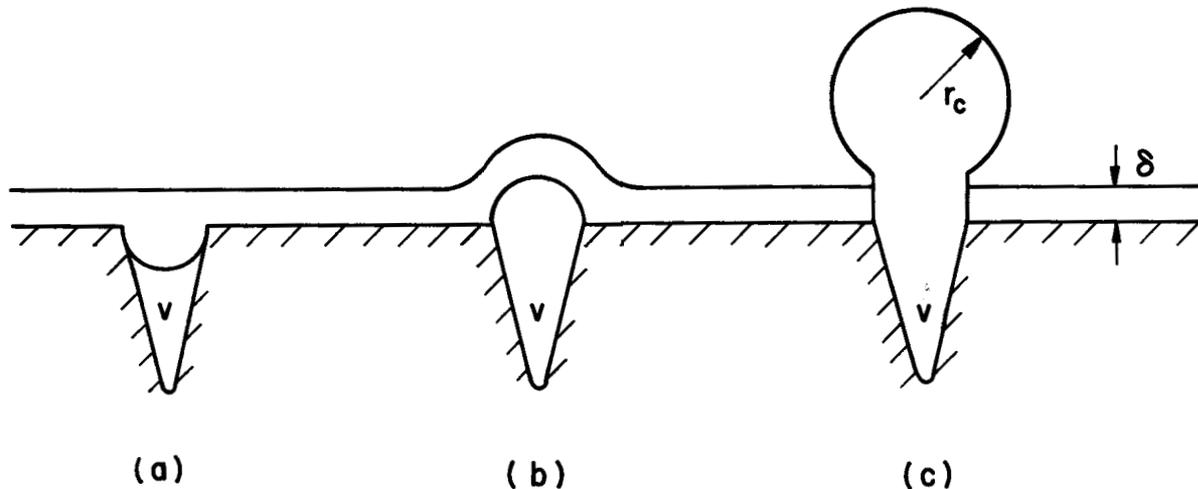
Note: Other cryogenic fluids have basically the same behavior, although the numerical values of q and ΔT are different.

Nucleate Boiling Heat Transfer

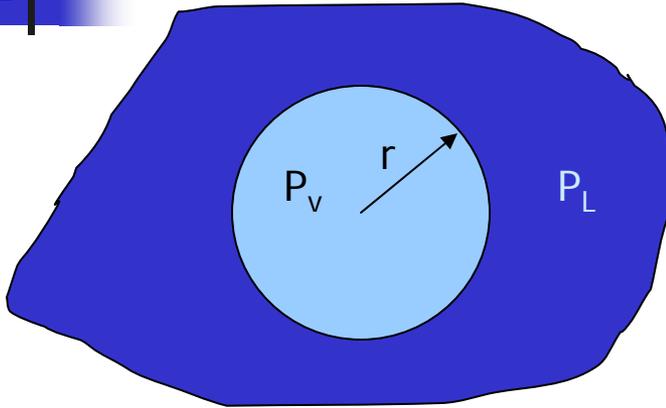
- Nucleate boiling is the principal heat transfer mechanism for static liquids below the peak heat flux ($q^* \sim 10 \text{ kW/m}^2$ for helium)
- Requirements for nucleate boiling
 - Must have a thermal boundary layer of superheated liquid near the surface

$$\delta_{th} = \frac{k_f \Delta T}{q} \sim 1 \text{ to } 10 \mu\text{m for helium}$$

- Must have surface imperfections that act as nucleation sites for formation of vapor bubbles.



Critical Radius of Vapor Bubble



$$P_v = P_L + \frac{2\sigma}{r} \quad \sigma \text{ is the surface tension}$$

- Critical radius: For a given T, p , the bubble radius that determines whether the bubble grows or collapses
 - $r > r_c$ and the bubble will grow
 - $r < r_c$ and the bubble will collapse
- Estimate the critical radius of a bubble using thermodynamics
 - Clausius Clapeyron relation defines the slope of the vapor pressure line in terms of fundamental properties

$$\left. \frac{dp}{dT} \right)_{sat} = \frac{\Delta s}{\Delta v} = \frac{h_{fg}}{T(v_v - v_L)} \approx \frac{h_{fg} p}{RT^2}$$

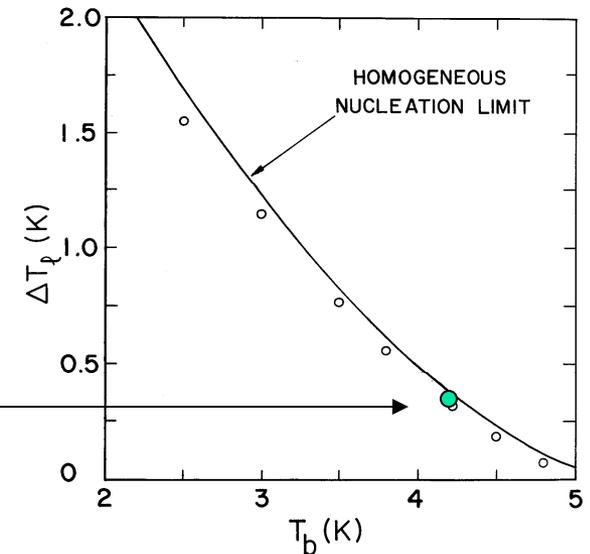
If the gas can be approximated as ideal and $v_v \gg v_L$

Critical radius calculation

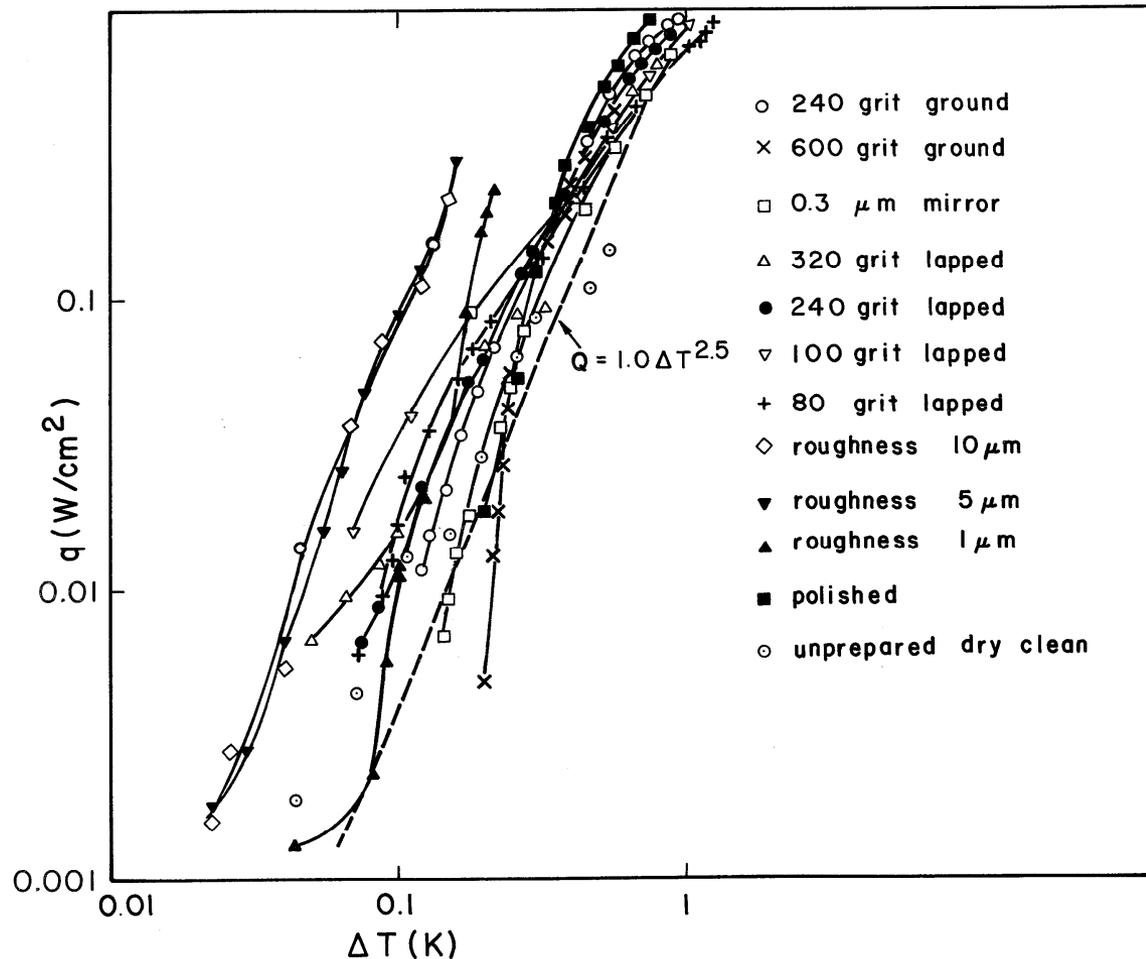
- Integrating the Clausius Clapeyron relation between p_s and $p_s + 2\sigma/r$

$$r_c = \frac{2\sigma}{p_s} \left(e^{h_{fg}\Delta T/RT^2} - 1 \right)^{-1} \approx \frac{2\sigma RT_s^2}{h_{fg} p_s \Delta T}$$

- Example: helium at 4.2 K (NBP)
 - Empirical evidence indicates that $\Delta T \sim 0.3$ K
 - This corresponds to $r_c \sim 17$ nm
 - Number of helium molecules in bubble $\sim 10,000$
 - Bubble has sufficient number of molecules to be treated as a thermodynamic system
- Actual nucleate boiling heat transfer involves heterogeneous nucleation of bubbles on a surface. This is more efficient than homogeneous nucleation and occurs for smaller ΔT .



Nucleate Boiling Heat Transfer He I



Note that h_{nb} is not constant because $Q \sim \Delta T^{2.5}$

Nucleate Boiling Heat Transfer Correlations

- The mechanism for bubble formation and detachment is very complex and difficult to model
- Engineering correlations are used for analysis
- Kutateladse correlation

$$\frac{h}{k_l} \left(\frac{\sigma}{g\rho_l} \right)^{1/2} = 3.25 \times 10^{-4} \left[\frac{q C_{pl} \rho_l}{h_{fg} \rho_v k_l} \left(\frac{\sigma}{g\rho_l} \right)^{1/2} \right]^{0.6} \times \left[g \left(\frac{\rho_l}{\mu_l} \right)^2 \left(\frac{\sigma}{g\rho_l} \right)^{3/2} \right]^{0.125} \left(\frac{p}{(\sigma g \rho_l)^{1/2}} \right)^{0.7}$$

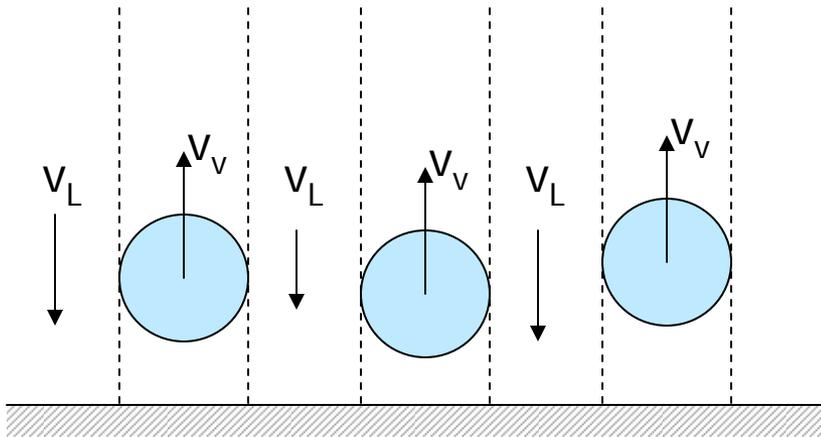
- Rearranging into a somewhat simpler form,

$$q = 1.90 \times 10^{-9} \left[g \left(\frac{\rho_l}{\mu_l} \right)^2 \chi^3 \right]^{0.3125} \left(\frac{p\chi}{\sigma} \right)^{1.75} \left(\frac{\rho_l}{\rho_v} \right)^{1.5} \times \left(\frac{C_p}{h_{fg}} \right)^{1.5} \left(\frac{k_l}{\chi} \right) (T_s - T_b)^{2.5}$$

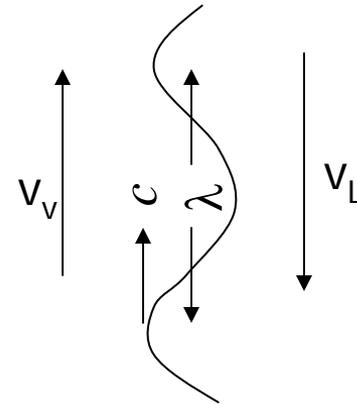
where $\chi = \left(\frac{\sigma}{g\rho_l} \right)^{1/2}$ $q(W / cm^2) = 5.8\Delta T^{2.5}$ For helium at 4.2 K

Peak Heat Flux (theory)

- Understanding the peak nucleate boiling heat flux is based on empirical arguments due to instability in the vapor/liquid flow



Instability in the vapor-liquid boundary



- Instability due to balance between surface energy and kinetic energy

$$c^2 = \frac{2\pi\sigma}{\lambda(\rho_L + \rho_v)} - \frac{\rho_L\rho_v}{(\rho_L + \rho_v)^2} (v_v - v_L)^2$$

Transition to unstable condition when $c^2 = 0$

Peak Heat Flux Correlations

- Zuber correlation:

$$q^* \approx Kh_{fg} \rho_v \left[\frac{\sigma(\rho_l - \rho_v)g}{\rho_v^2} \right]^{1/4} \left[\frac{\rho_l}{\rho_l + \rho_v} \right]^{1/2}$$

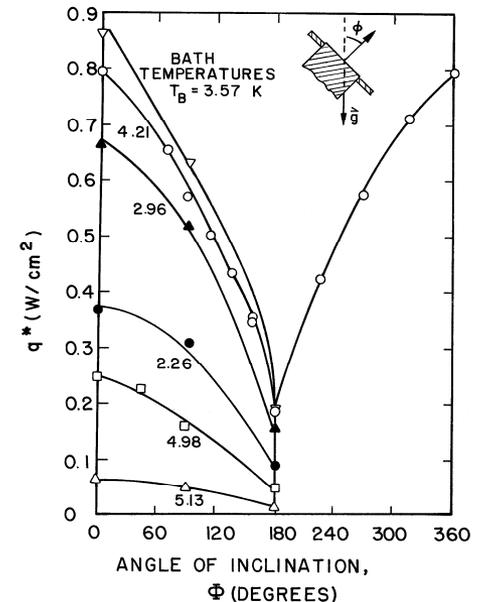
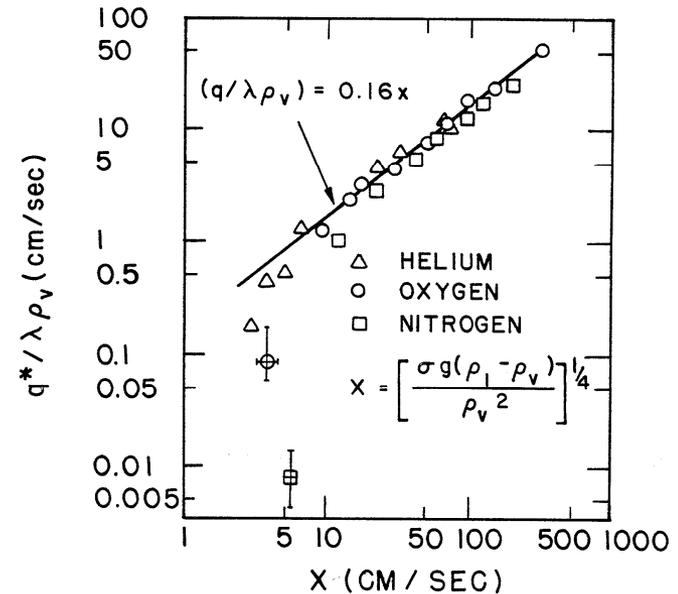
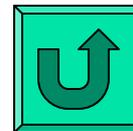
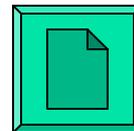
- Empirical based on Zuber Correlation

$$q^* \approx 0.16h_{fg} \rho_v^{1/2} [\sigma g(\rho_l - \rho_v)]^{1/4}$$

~ 8.5 kW/m² for He I at 4.2 K

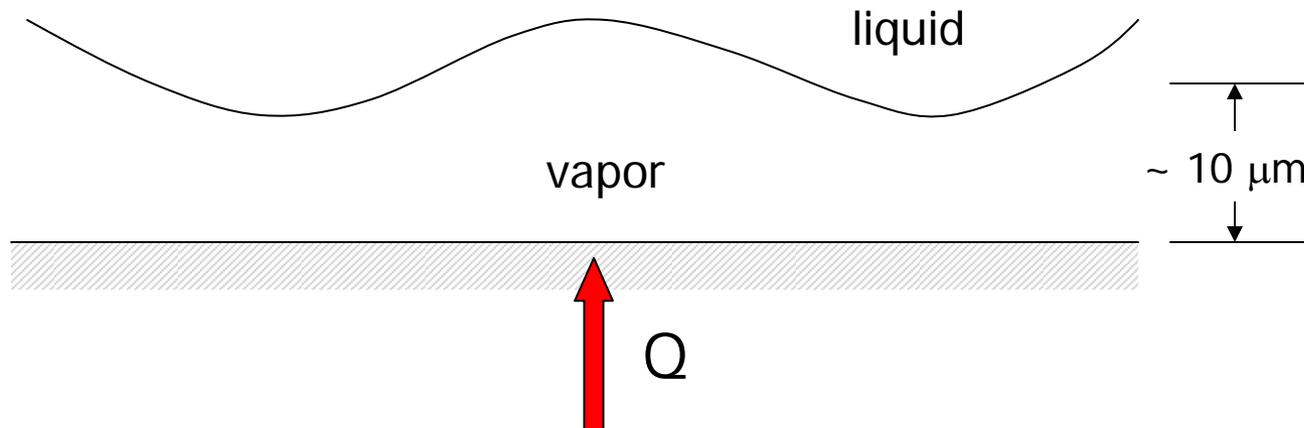
- Limits:

- $T \rightarrow T_c$; $q^* \rightarrow 0$ since $h_{fg} \rightarrow 0$ and $\sigma \rightarrow 0$
- $T \rightarrow 0$; $q^* \sim \rho_v^{1/2}$ (decreases)
- q^*_{max} near 3.6 K for LHe



Film Boiling

- Film boiling is the stable condition when the surface is blanketed by a layer of vapor
 - Film boiling heat transfer coefficient is generally much less than that in nucleate boiling
 - Minimum film boiling heat flux, q_{mfb} is related to the stability of the less dense vapor film under the more dense liquid



"Taylor Instability" governs the collapse of the vapor layer

Film Boiling Heat Transfer Correlations

- Factors affecting the process
 - Fluid properties: C_p , h_{fg} , σ , ρ_l , ρ_v
 - Fluid state: saturated or pressurized (subcooled)
 - Heater geometry (flat plate, cylinder, etc.)
- Breen-Westwater correlation

$$h_{fb} \left(\frac{\sigma}{g(\rho_l - \rho_v)} \right)^{1/8} \left(\frac{\mu_v(T_s - T_b)}{k_v^3 \rho_v (\rho_l - \rho_v) g \lambda'} \right)^{1/4} = 0.37 + 0.28 \left(\frac{\sigma}{g D^2 (\rho_l - \rho_v)} \right)^{1/2}$$

Where,
$$\lambda' = \frac{[h_{fg} + 0.34 C_{pv} (T_s - T_b)]^2}{h_{fg}}$$

- Simplified form for large diameter

$$h_{fb} = 0.37 \left(\frac{g(\rho_l - \rho_v)}{\sigma} \right)^{1/8} \left(\frac{k_v^3 \rho_v (\rho_l - \rho_v) g \lambda'}{\mu_v (T_s - T_b)} \right)^{1/4} \rightarrow q \approx (T_s - T_f)^{3/4}$$

Minimum film boiling heat flux

- Minimum film boiling heat flux is less than the peak heat flux
- Recovery to nucleate boiling state is associated with Taylor Instability.

$$q_{mfb} = 0.16 h_{fg} \rho_v \left[\frac{g \sigma (\rho_l - \rho_v)}{(\rho_l + \rho_v)^2} \right]^{1/4}$$

- Dimensionless ratio:

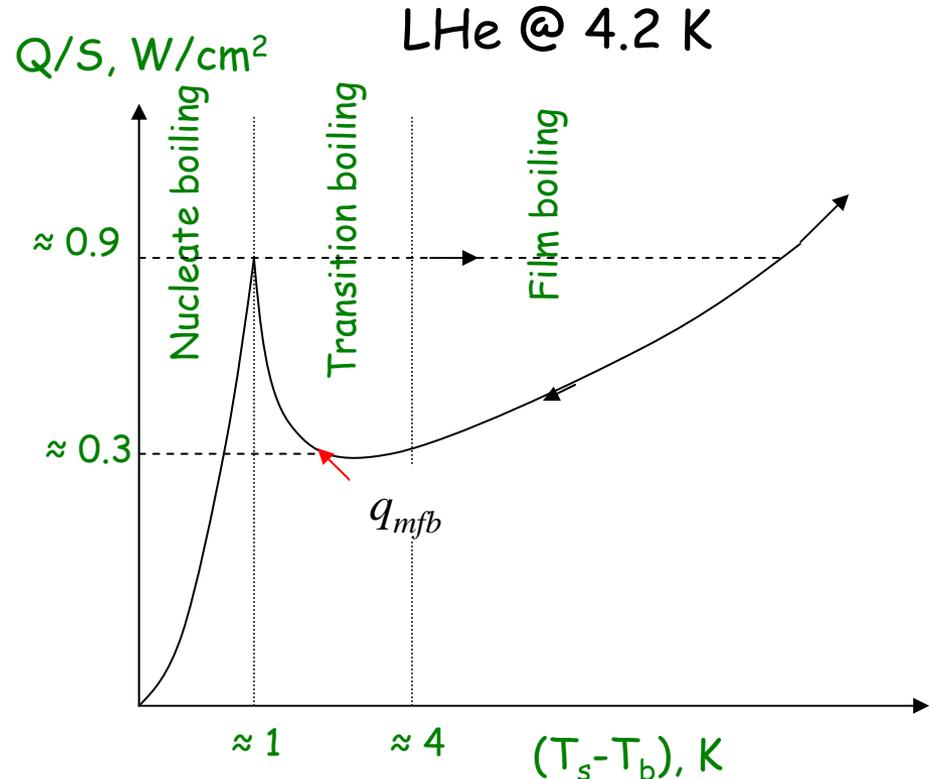
$$\frac{q_{mfb}}{q^*} = \left[\frac{\rho_v}{\rho_l + \rho_v} \right]^{1/2}$$

~ 0.36 for LHe @ 4.2 K

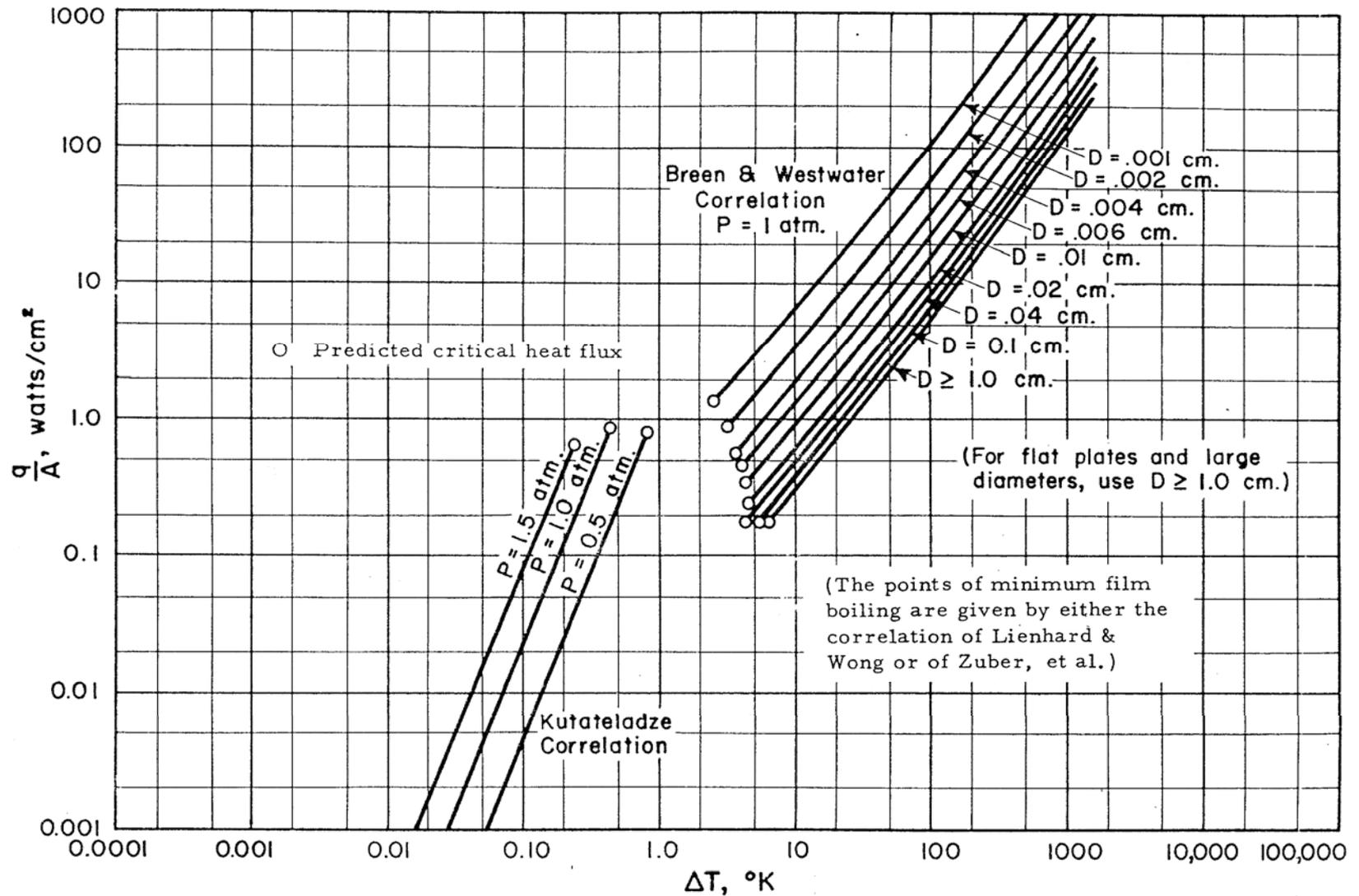
~ 0.1 @ 2.2 K

~ 1 @ $T_c = 5.2$ K

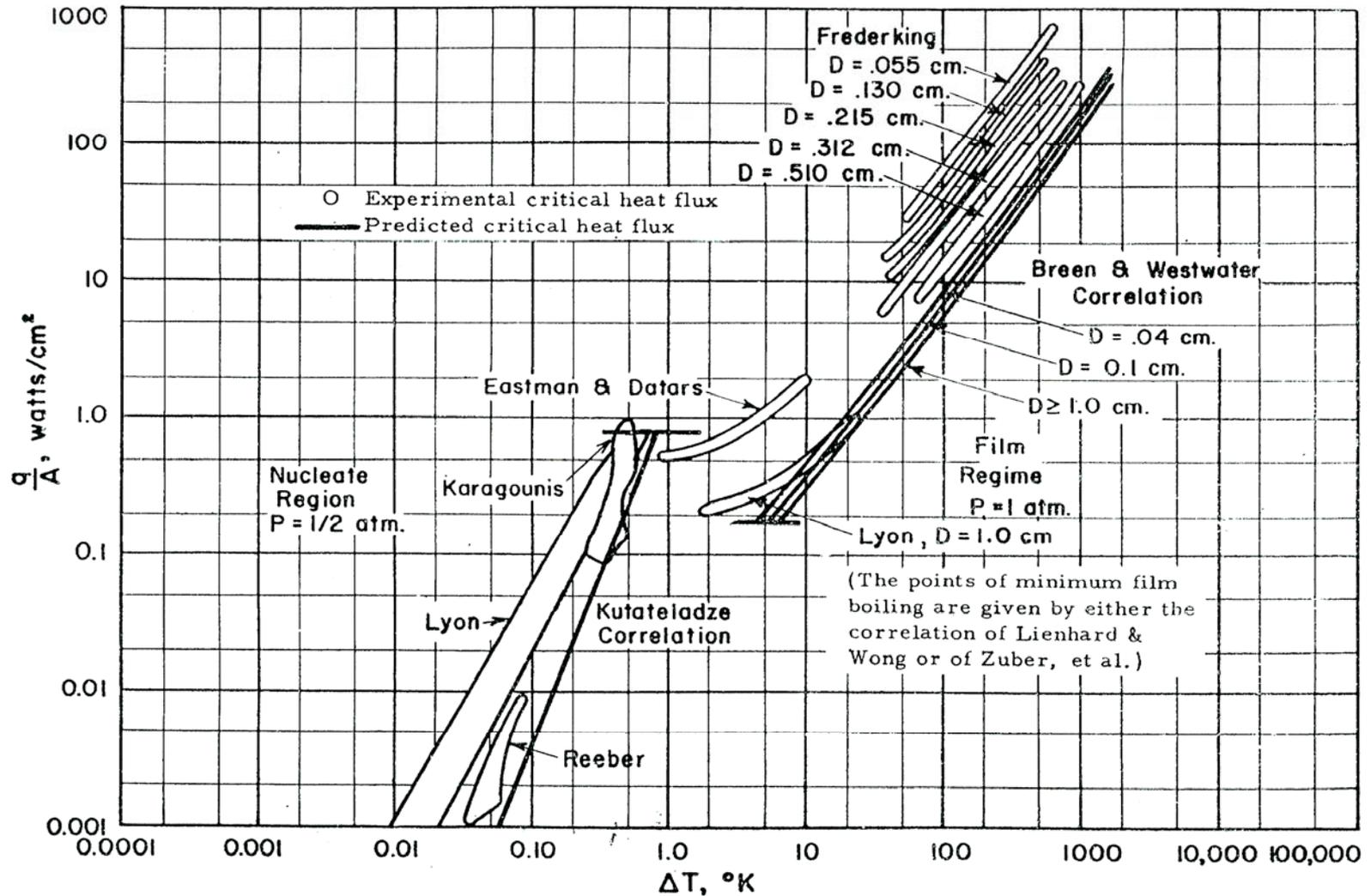
~ 0.09 for LN₂ @ 80 K where $q^* \sim 200$ kW/m²



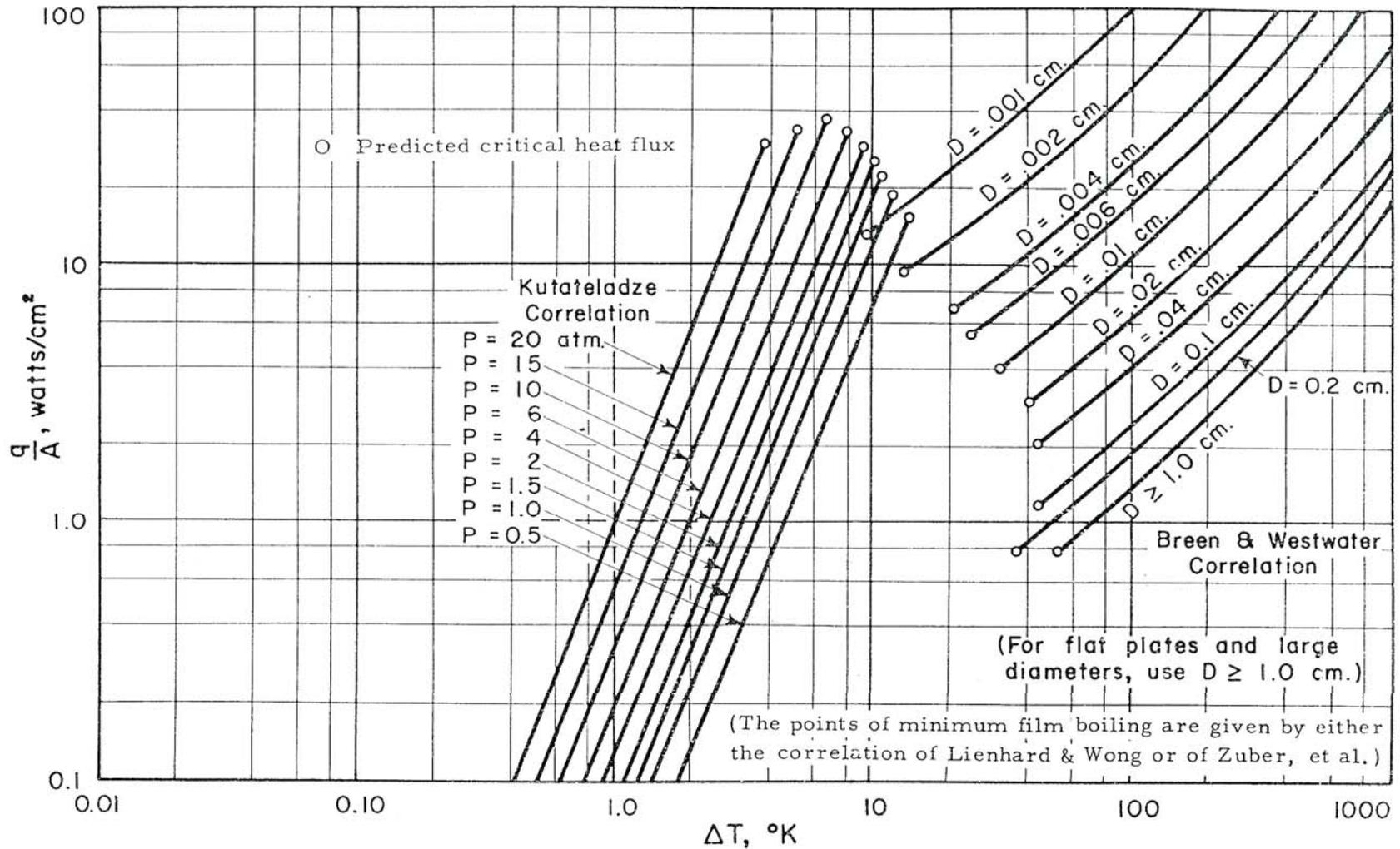
Prediction of Nucleate/Film Boiling for Helium



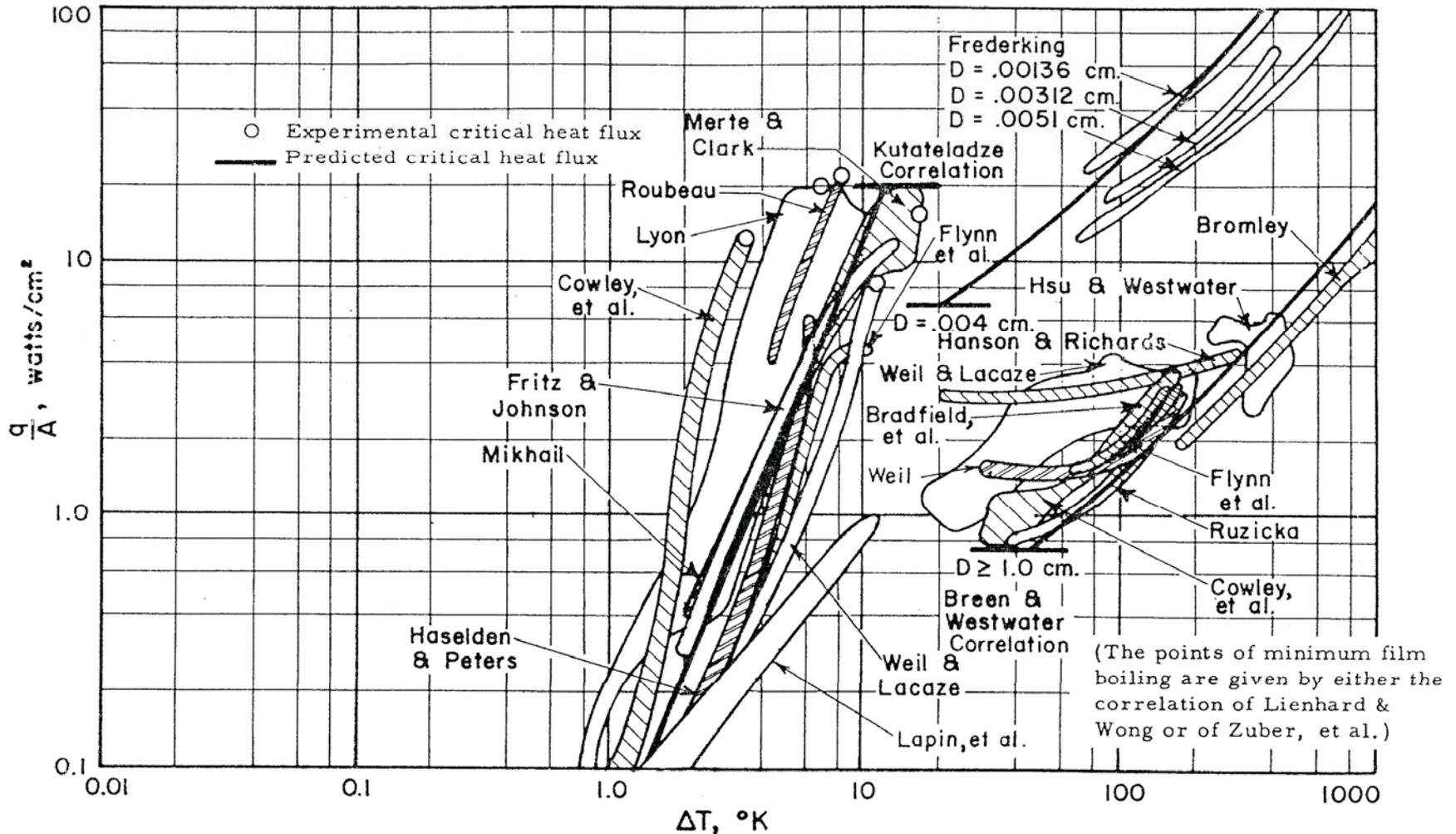
Experimental Heat Transfer (Helium)



Prediction of Nucleate/Film Boiling for Nitrogen



Experimental Heat Transfer (Nitrogen)



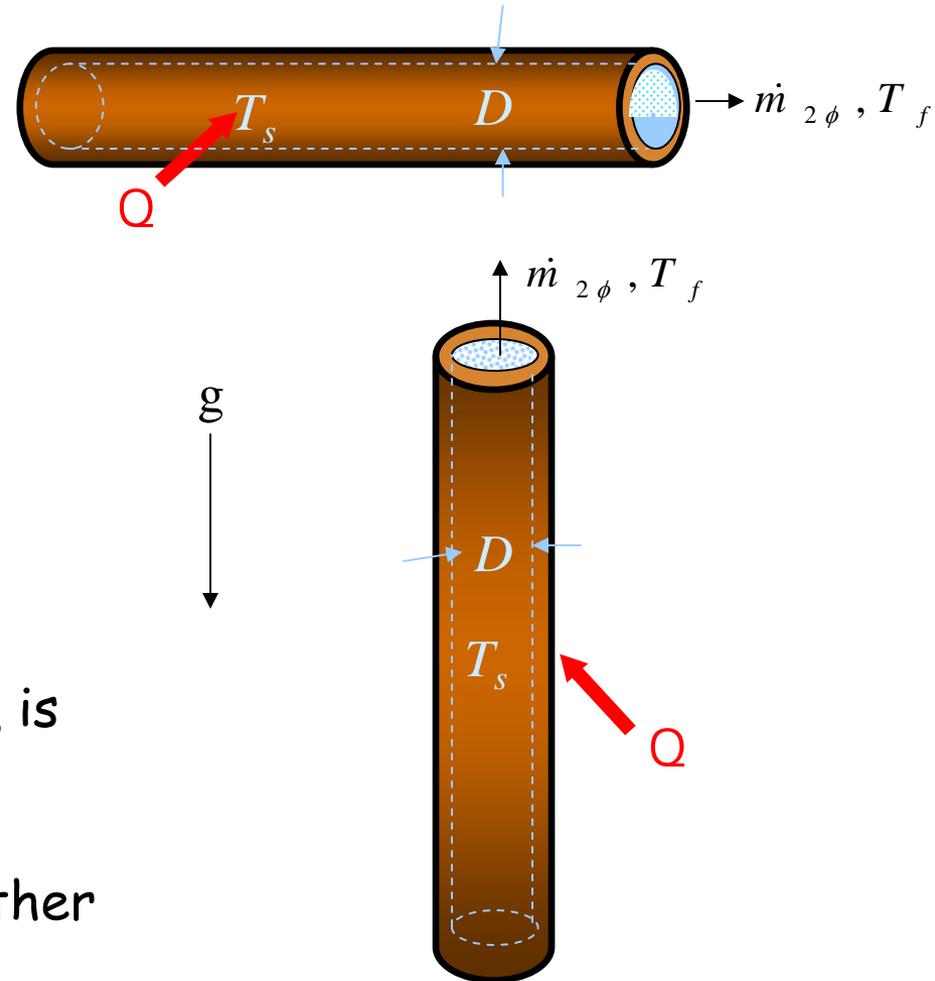
Internal two phase flow

- Heat transfer depends on various factors
 - Mass flow rate
 - Orientation w/r/t gravity
 - Flow regime
 - Quality (χ)
 - Void fraction (α)
- Total heat transfer rate

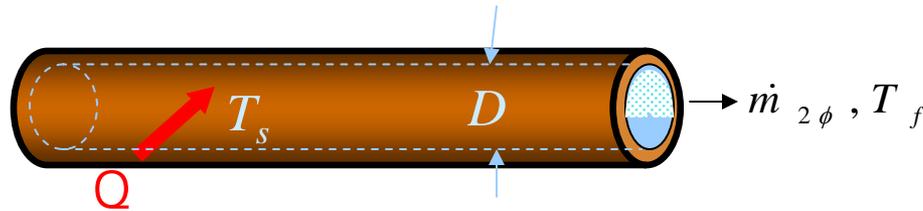
$$Q_T = Q_{fc} + Q_b$$

where: Q_{fc} is convective and Q_b is gravity enhanced boiling.

Depending on factors above, either contribution may dominate



Horizontal flow two phase heat transfer



- Consider the case where gravitational effects are negligible
 - Horizontal flow at moderate Re so that inertial forces dominate
- Correlation based on enhanced Nusselt number

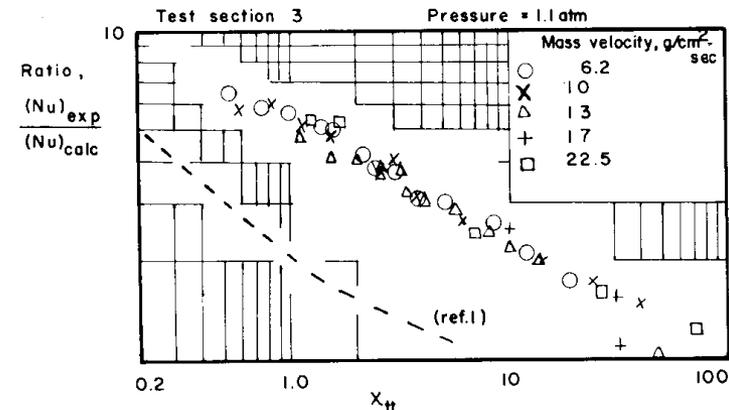
$$\frac{Nu_{2\phi}}{Nu_L} = f(\chi_{tt}) \quad \text{where} \quad \chi_{tt} = \left(\frac{1-x}{x}\right)^{0.9} \left(\frac{\mu_L}{\mu_v}\right)^{0.1} \left(\frac{\rho_v}{\rho_L}\right)^{0.5}$$

$$\text{and} \quad Nu_L = 0.023 Re_L^{0.8} Pr_L^{0.4} ; \quad Re_L = \frac{\dot{m}(1-x)}{\mu_L A_{flow}}$$

Typical correlation (de la Harpe):

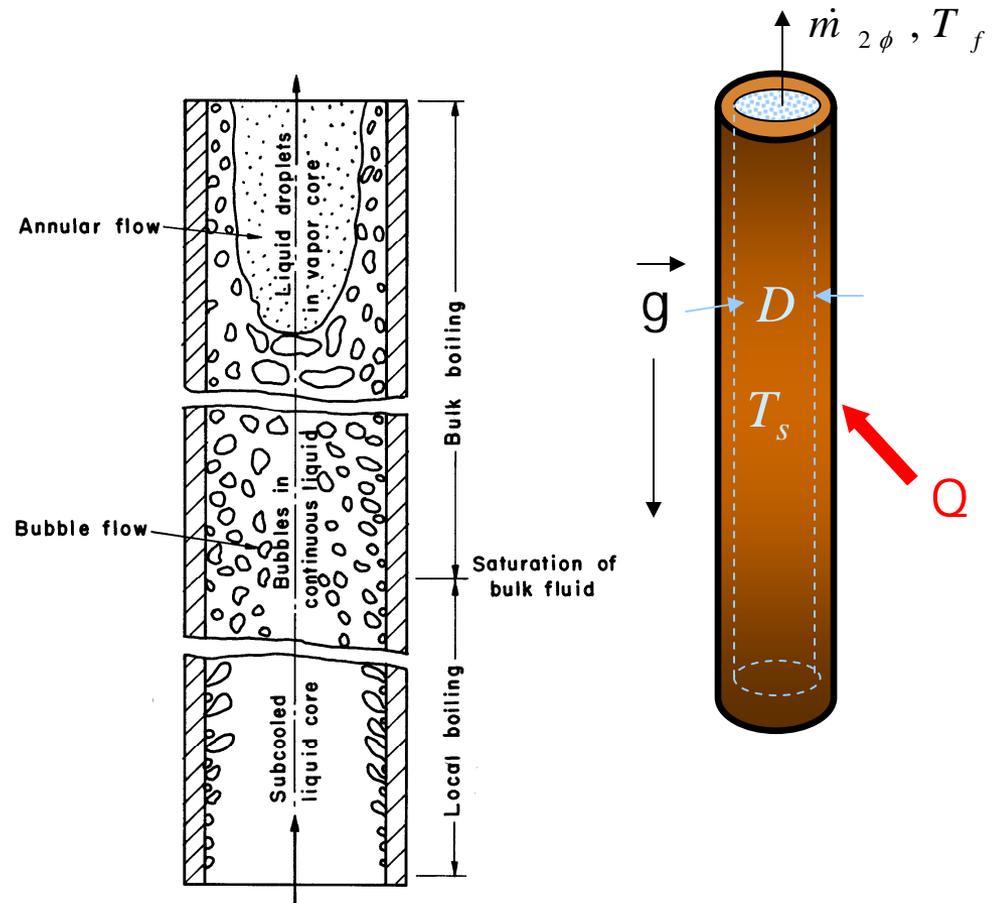
$$\frac{Nu_{2\phi}}{Nu_L} \approx A \chi_{tt}^{-m} \rightarrow 1 \text{ for } \chi_{tt} \text{ large}$$

with $m \sim 0.385$ and $A \sim 5.4$

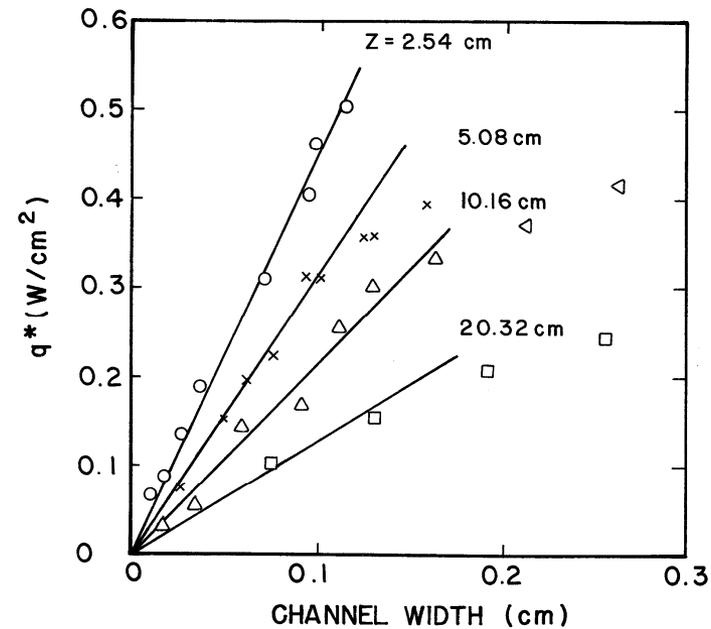
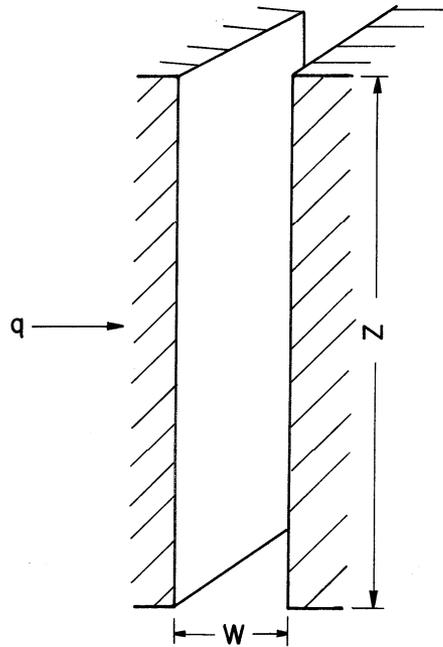


Vertical channel heat transfer

- Main difference between this problem and pool boiling is that the fluid is confined within channel
- At low mass flow rate and self driven flows (natural circulation) the heat transfer is governed by buoyancy effects
- Process is correlated against classical boiling heat transfer models
- In the limit of large D the correlation is similar to pool boiling heat transfer (vertical surface)



Vertical channel maximum heat flux



Critical flow in an evaporator

Empirical observations

- $q^* \sim w$ for small w
- $q^* \sim z^{-1/2}$ for $w/z < 0.1$
- $q^* \sim \text{constant}$ for $w/z \gg 1$

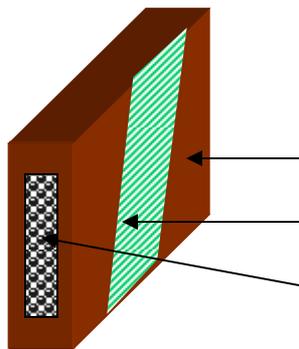
$$q^* = \frac{w}{\sqrt{z}} \frac{h_{fg} \rho_v}{2} \left[\frac{g \chi_c}{\beta - 1} \left(1 - \frac{\ln[1 + \chi_c (\beta - 1)]}{\chi_c (\beta - 1)} \right) \right]^{1/2} \quad \beta = \frac{\rho_l}{\rho_v}$$

and $\chi_c = \left(\frac{\dot{m}_v}{\dot{m}} \right)_{\max}$ "Critical quality"
 ~ 0.3 for helium



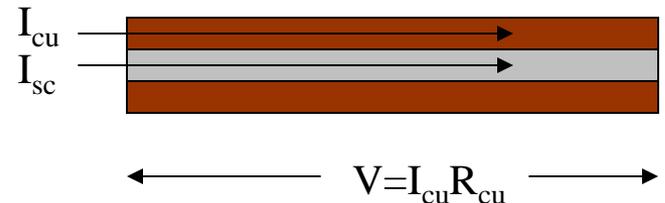
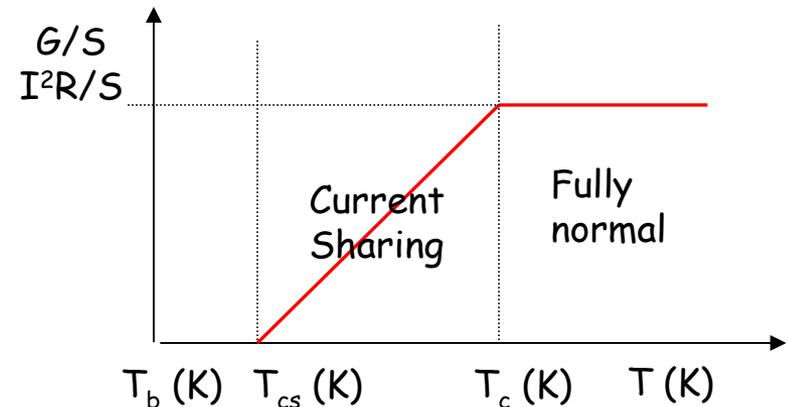
Example: Cryogenic Stability of Composite Superconductors (LTS in LHe @ 4.2 K)

- Used in large magnets where flux jumping and other small disturbances are possible and must be arrested
- General idea: in steady state ensure that cooling rate exceeds heat generation rate ($Q > G$)
- Achieved by manufacturing conductor with large copper (or aluminum) fraction and cooling surface
- Lower overall current density
- Potentially high AC loss (eddy currents)



Composite Conductor

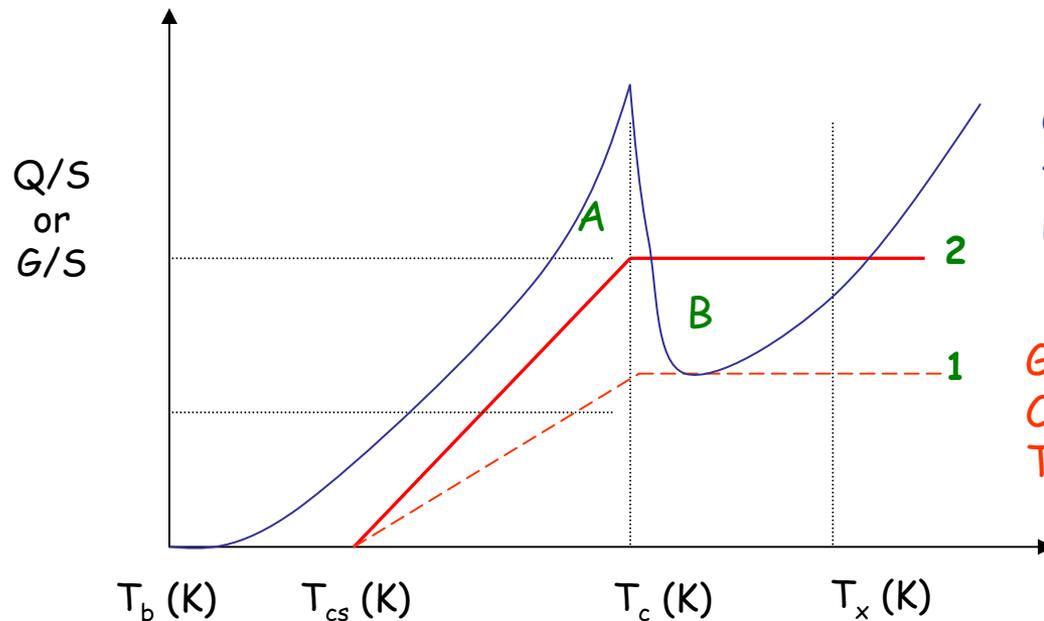
Copper/Aluminum stabilizer
Insulating spacer (G-10)
NbTi/Nb₃Sn



Cryogenic Stability (LHe @ 4.2 K)

Case 1: Unconditional stability, recovery to fully superconducting state occurs uniformly over length of normal zone

Case 2: Cold End recovery (Equal area criterion): Excess cooling capacity (area A) > Excess heat generation (area B)

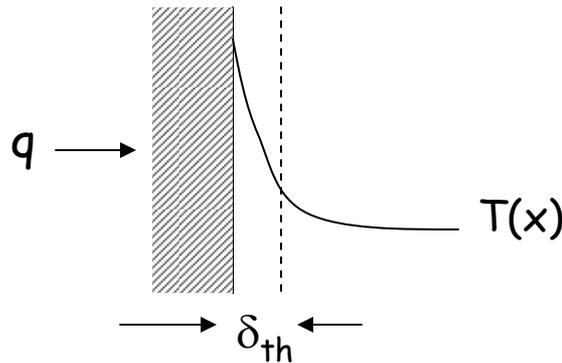


Q/S is the LHe boiling heat transfer curve for bath cooling normalized per surface area

G/S is the two part heat generation Curve for a composite superconductor
 T_{cs} is the temperature at which $T_{op} = T_c$

Transient Heat Transfer

- Heat transfer processes that occur on time scale short compared to boundary layer thermal diffusion. Why is this important in cryogenics? (D_{th} (copper) $\sim 10^{-4} \text{ m}^2/\text{s}$ @ 300 K; $\sim 1 \text{ m}^2/\text{s}$ @ 4 K)
- Normal liquid helium has a low thermal conductivity and large heat capacity



$$D_f = \frac{k_f}{\rho C} \sim 3 \times 10^{-8} \text{ m}^2/\text{s} \text{ (LHe @ 4.2 K)}$$

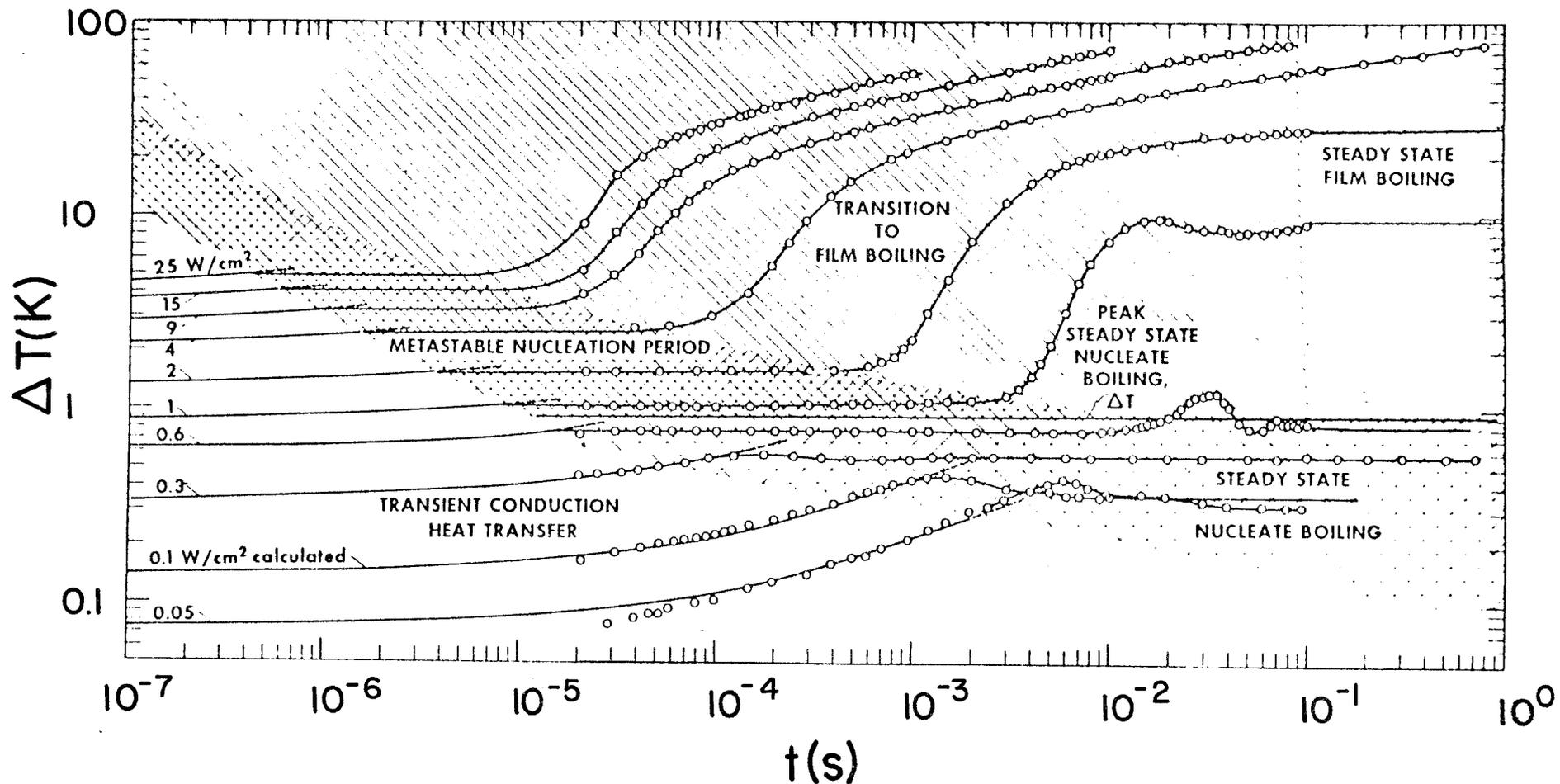
$$\delta_{th} = \frac{\pi}{2} (D_f t)^{1/2} \sim 3 \times 10^{-4} t^{1/2} \text{ [m] for LHe @ 4.2 K}$$

$$\sim 1.5 t^{1/2} \text{ [m] for copper @ 4.2 K}$$

- Lumped capacitance condition: $Bi = \frac{hL}{k} \ll 1 \sim 10L \text{ [m]}$
- Note that this subject is particularly relevant to cooling superconducting magnets, with associated transient thermal processes
- Important parameters to determine
 - $\Delta E = q \Delta t^*$ (critical energy)
 - T_s - surface temperature during heat transfer

Transient Heat Transfer to LHe @ 4.2 K

Time evolution of the temperature difference following a step heat input:
Steady state is reached after ~ 0.1 s



Critical Energy for Transition to Film Boiling

Hypothesis: The "critical energy" is determined by the amount of energy that must be applied to vaporize a layer of liquid adjacent to the heated surface.

- Energy required
- Layer thickness determined by diffusion

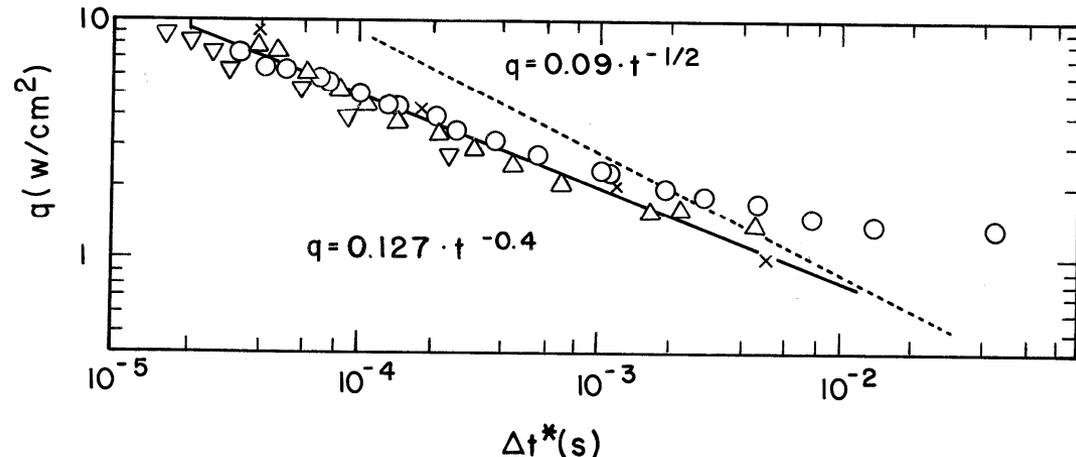
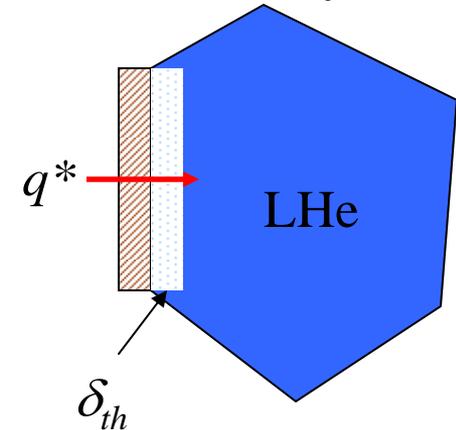
$$\Delta E = \rho_l h_{fg} \delta_{th}$$

$$\delta_{th} = \frac{\pi}{2} (D_f t)^{1/2}$$

- Critical flux based on heat diffusion:

$$q^* = \frac{\pi}{2} \rho_l h_{fg} \left(\frac{k_l}{\rho_l C \Delta t^*} \right)^{1/2}$$

$$\sim 0.09 \Delta t^{1/2} \text{ [W/cm}^2\text{] @ 4.2 K}$$



Surface Temperature (Transient HT)

- During transient heat transfer, the surface temperature will be higher than the surrounding fluid due to two contributions:
 - Fluid layer diffusion: transient conduction in the fluid layer will result in a finite temperature difference
 - Kapitza conductance: At low temperatures, there can be a significant temperature difference, ΔT_k , due to thermal impedance mismatch (more on this subject later). This process is dominant at very low temperatures, but is small above ~ 4 K, so is normally only important for helium systems.
- Fluid layer diffusion equation:

$$\frac{\partial^2 \Delta T_f}{\partial x^2} = \frac{1}{D_f} \frac{\partial \Delta T_f}{\partial t}; D_f = \frac{k}{\rho C_p}$$

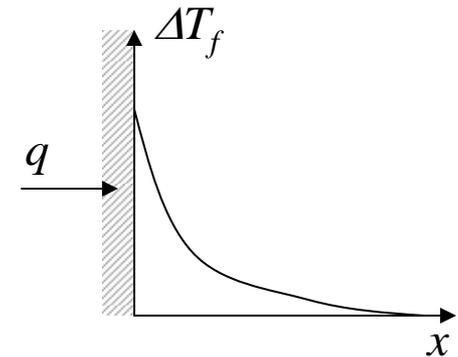
- Boundary conditions:
 - $\Delta T_f(x,0) = 0$; initial condition
 - $\Delta T_f(\text{infinity}, t) = 0$; isothermal bath
 - $q = -k_f d\Delta T_f/dt)_{x=0}$; heat flux condition
 - Solid is isothermal ($Bi = hL/k \ll 1$)

Solution is a standard second order differential equation with two spatial and one time boundary condition.

Transient diffusion solution

- Integrating the diffusion equation:

$$\Delta T_f = \frac{q}{k_f} \left[2 \left(\frac{D_f t}{\pi} \right)^{1/2} \exp\left(-\frac{x^2}{4D_f t}\right) - x \operatorname{erfc}\left(\frac{x}{2(D_f t)^{1/2}}\right) \right]$$



- Evaluating at $x = 0$ (surface of heater)

$$\Delta T_f(x=0) = \frac{2q}{\sqrt{\pi}} \left(\frac{t}{\rho C_p k_f} \right)^{1/2}$$

- The transient heat transfer coefficient can then be defined

$$h = \frac{q}{\Delta T_f} = \frac{\sqrt{\pi}}{2} \left(\frac{\rho C_p k_f}{t} \right)^{1/2}$$

$$h \approx \frac{0.1}{\sqrt{t}} \quad ; \text{ kW/m}^2\text{K for helium near 4 K}$$

At $t = 10 \mu\text{s}$; $h \sim 30 \text{ kW/m}^2 \text{ K}$ and for $q = 10 \text{ kW/m}^2$; $\Delta T_f \sim 0.3 \text{ K}$

Note: this value of $h \gg h_{nb}$ (nucleate boiling HT coefficient)

Summary Cryogenic Heat Transfer

- Single phase heat transfer correlations for classical fluids are generally suitable for cryogenic fluids
 - Free convection
 - Forced convection
- Two phase heat transfer also based on classical correlations
 - Nucleate boiling
 - Peak heat flux
 - Film boiling
- Transient heat transfer is governed by diffusive process for ΔT and onset of film boiling