



Fluctuation-Based Bunch Length Experiments

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- Motivation
- Time Domain Measurements
- Frequency-based Measurements
- Interferometer-based Measurements
- Introduction to USPAS Simulator



Motivation

- Alan derived theoretical basis for using statistical fluctuations to measure pulse length
- Each electron is an independent 'radiator' with a random, granular distribution along the bunch (shot noise)
- Sometimes the phase of wave packets overlap, sometimes they don't
- The *mean and variance* (moments) in the signal yields pulse length
- Measurements can be made in the time domain or frequency domain
- We will review some experiments and introduce the USPAS simulator



Time Domain View

Sum electric field emission from individual electrons

$$E(t) = \sum_{k=1}^N e(t - t_k)$$

where emission times t_k are random, Gaussian-distributed numbers

$$f(t) = \frac{1}{\sqrt{2\pi\sigma_t}} e^{-t^2/2\sigma_t^2}$$

Each wavepacket $e(t)$ is centered at random time t_k

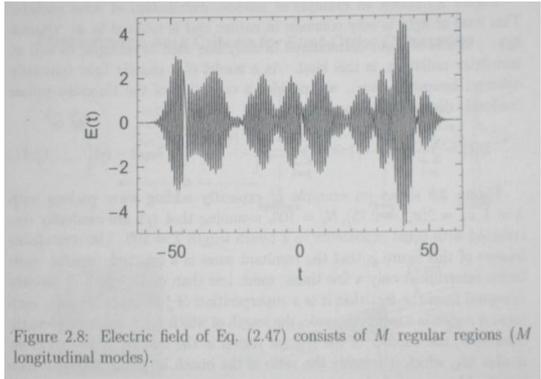
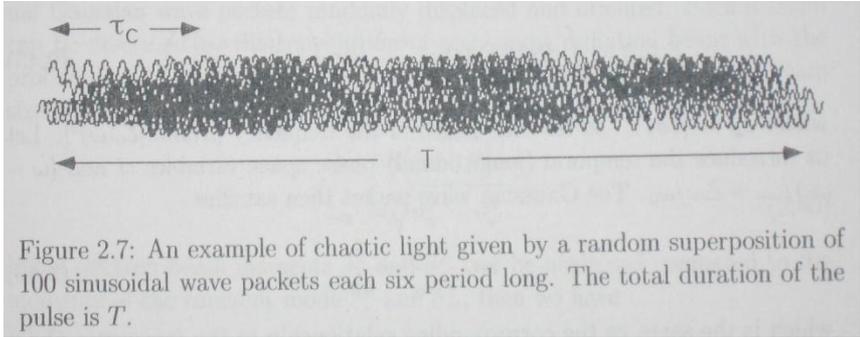
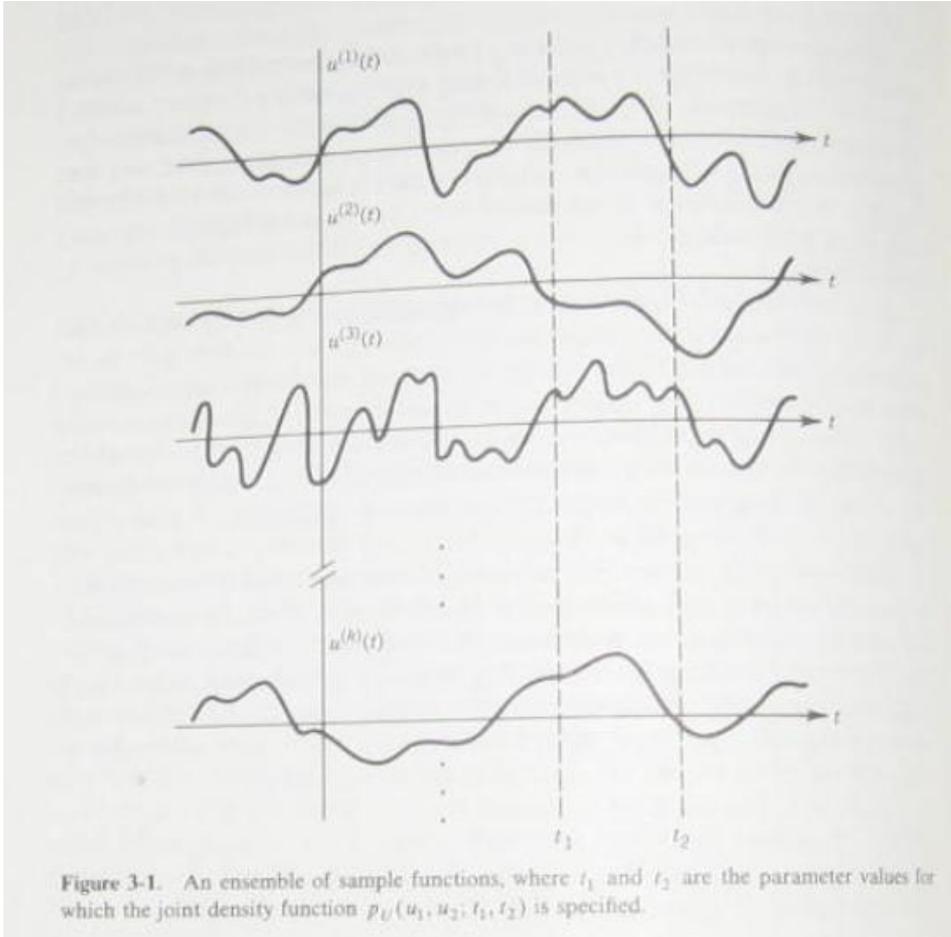
Wavepackets superimpose to produce more or less field at time t

The electromagnetic field intensity is E^*E

Total pulse energy $\int E^* E dt$ is therefore random in time.



Statistically Random Function in Time



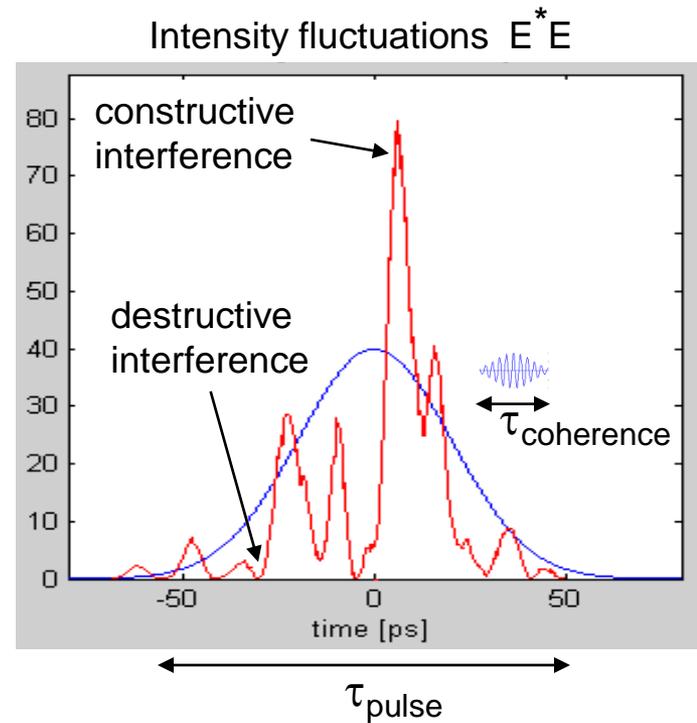
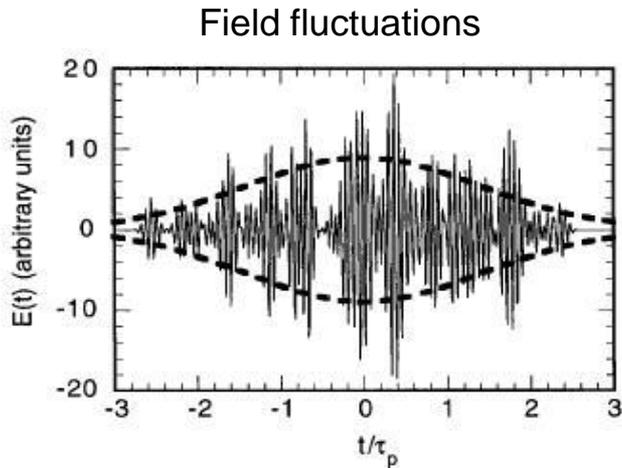
M~10

an incoherent electric field is often not what we were lead to believe -



Fluctuations in Electric Field and Intensity

Each light pulse from the synchrotron has statistical structure



wave packets emitted from individual electrons statistically add or cancel

the *correlation length* corresponds to the wavepacket *coherence length*

$$T_{\text{pulse}}/\tau_{\text{coh}} = M \text{ (mode number)}$$



Coherence Length and Coherence Time

To increase correlation length band-limit the radiation

This increases the coherence length of the individual wave packets

$$f = c / \lambda$$

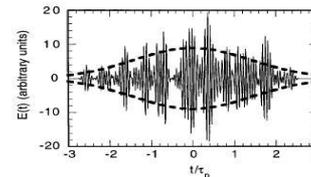
$$\delta f = -\delta \lambda c / \lambda^2$$

$$\delta t = 1 / \delta f = \lambda^2 / c \delta \lambda$$

For 633nm light and a 1nm band pass filter

$$\delta t = \lambda^2 / c \delta \lambda = \frac{(633 * 10^{-9})^2}{(3 * 10^8)(1 * 10^{-9})} = 1.3 \text{ ps}$$

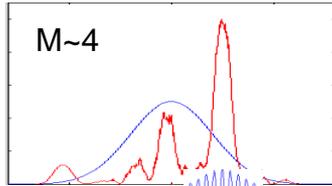
For a 15ps bunch, the 'mode number' $M \sim 15$.



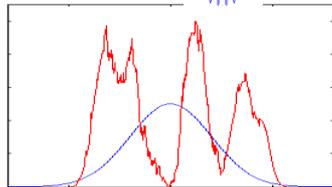


Energy Fluctuation Statistics

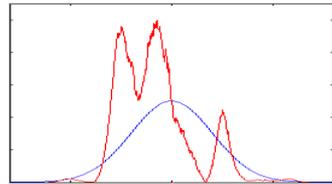
Narrow band filter



$$E_1 = \int P_1(t) dt$$

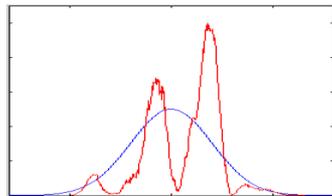


$$E_2 = \int P_2(t) dt$$



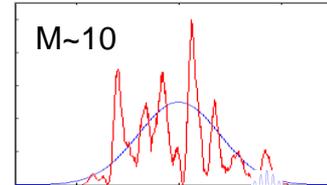
$$E_3 = \int P_3(t) dt$$

...

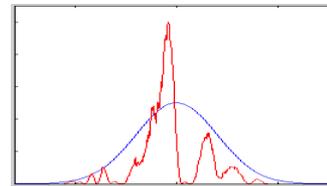


$$E_N = \int P_N(t) dt$$

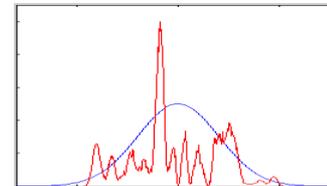
Wide band filter



$$E_1 = \int P_1(t) dt$$

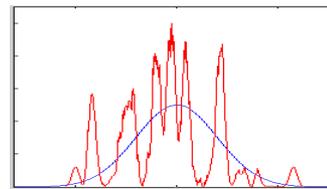


$$E_2 = \int P_2(t) dt$$

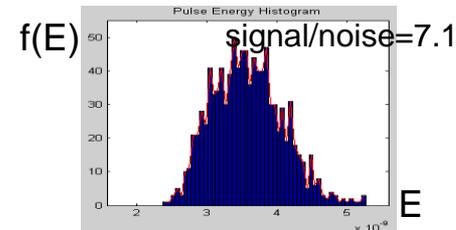
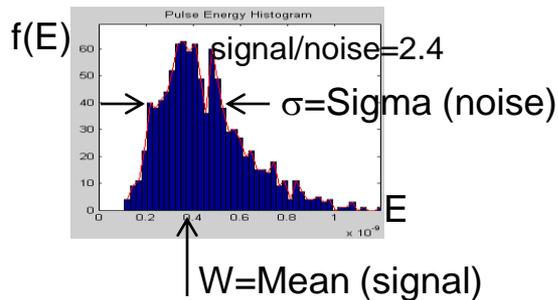


$$E_3 = \int P_3(t) dt$$

...



$$E_N = \int P_N(t) dt$$





Intensity Fluctuation Derivation

Goodman, *Statistical Optics* Chapter 6

Average Value $\bar{W} = \int_{-T}^T \bar{I}(t) dt$

Variance
$$\begin{aligned} \sigma_W^2 &= E \left[\left(\int_{-T}^T I(t) dt \right)^2 \right] - \bar{W}^2 \\ &= \int_{-T}^T \int_{-T}^T \overline{I(t)I(t')} dt dt' - \bar{W}^2 \\ \sigma_W^2 &= \int_{-T}^T \int_{-T}^T \Gamma_I(t-t') dt dt' - \bar{W}^2 \end{aligned}$$

where Γ is the autocorrelation function of $I(t)$

in terms of fields $\Gamma_I(\tau) = E \{ e(t) e^*(t) e(t+\tau) e^*(t+\tau) \}$
'fourth order correlation'



Intensity Fluctuations (cont'd)

$$\sigma_W^2 = \int_{-T}^T \int_{-T}^T \Gamma_I(t-t') dt dt' - \bar{W}^2$$

$$\Gamma_I(\tau) = E\{e(t)e^*(t)e(t+\tau)e^*(t+\tau)\} \quad \text{'fourth order correlation'}$$

But from interferometry $\Gamma_I(\tau) = I^2 \cdot (1 + |\gamma(\tau)|^2)$

Then
$$\sigma_W^2 = \bar{W}^2 \frac{1}{T} \int |\gamma(\tau)|^2 d\tau$$

$$\frac{\bar{W}^2}{\sigma_W^2} = \left(\frac{1}{T} \int |\gamma(\tau)|^2 d\tau \right)^{-1} = M \quad \text{(same as before)}$$

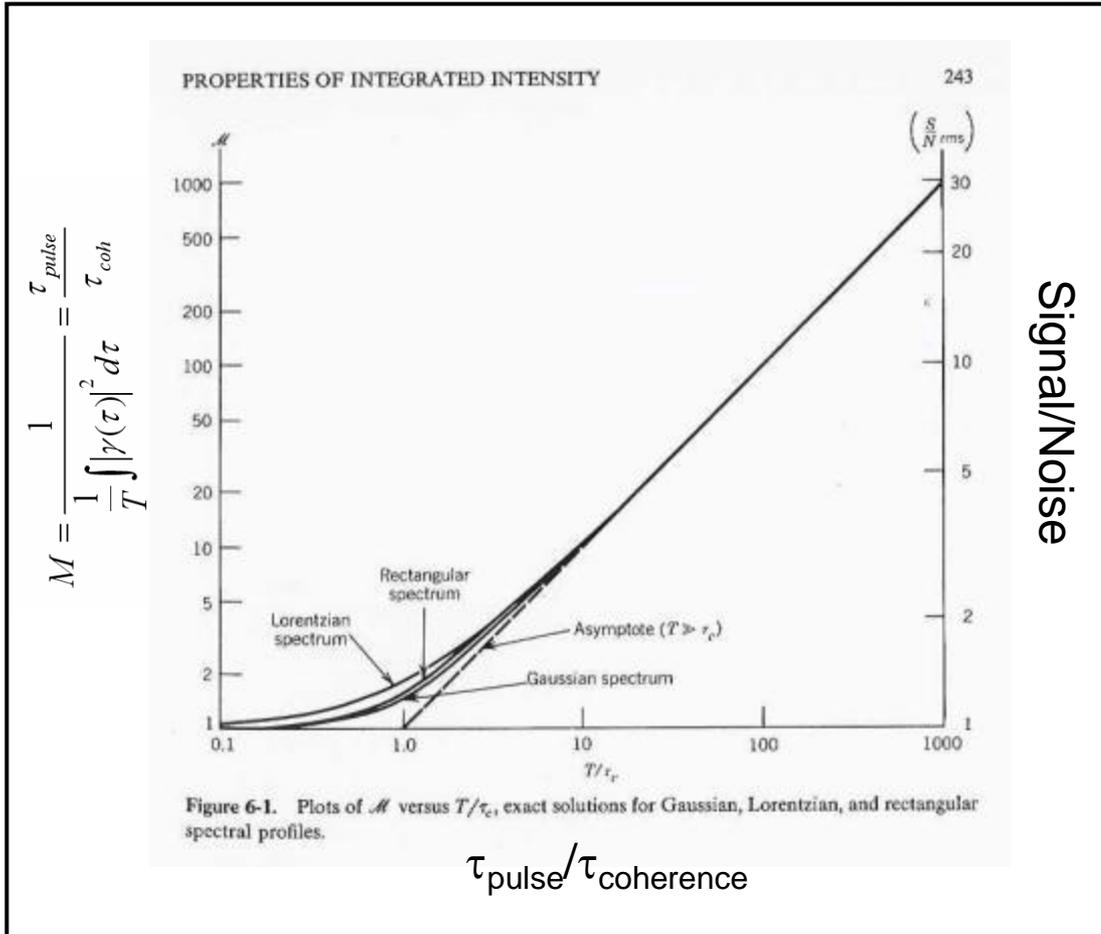
$$M = \frac{1}{\frac{1}{T} \int |\gamma(\tau)|^2 d\tau} = \frac{\tau_{pulse}}{\tau_{coh}} \quad \text{is the number of modes-per-pulse!}$$

→ measurement of W , σ_W with known τ_c yields τ_{pulse}



M: The ratio of Pulse Time to Coherence Time

Goodman, *Statistical Optics* Chapter 6

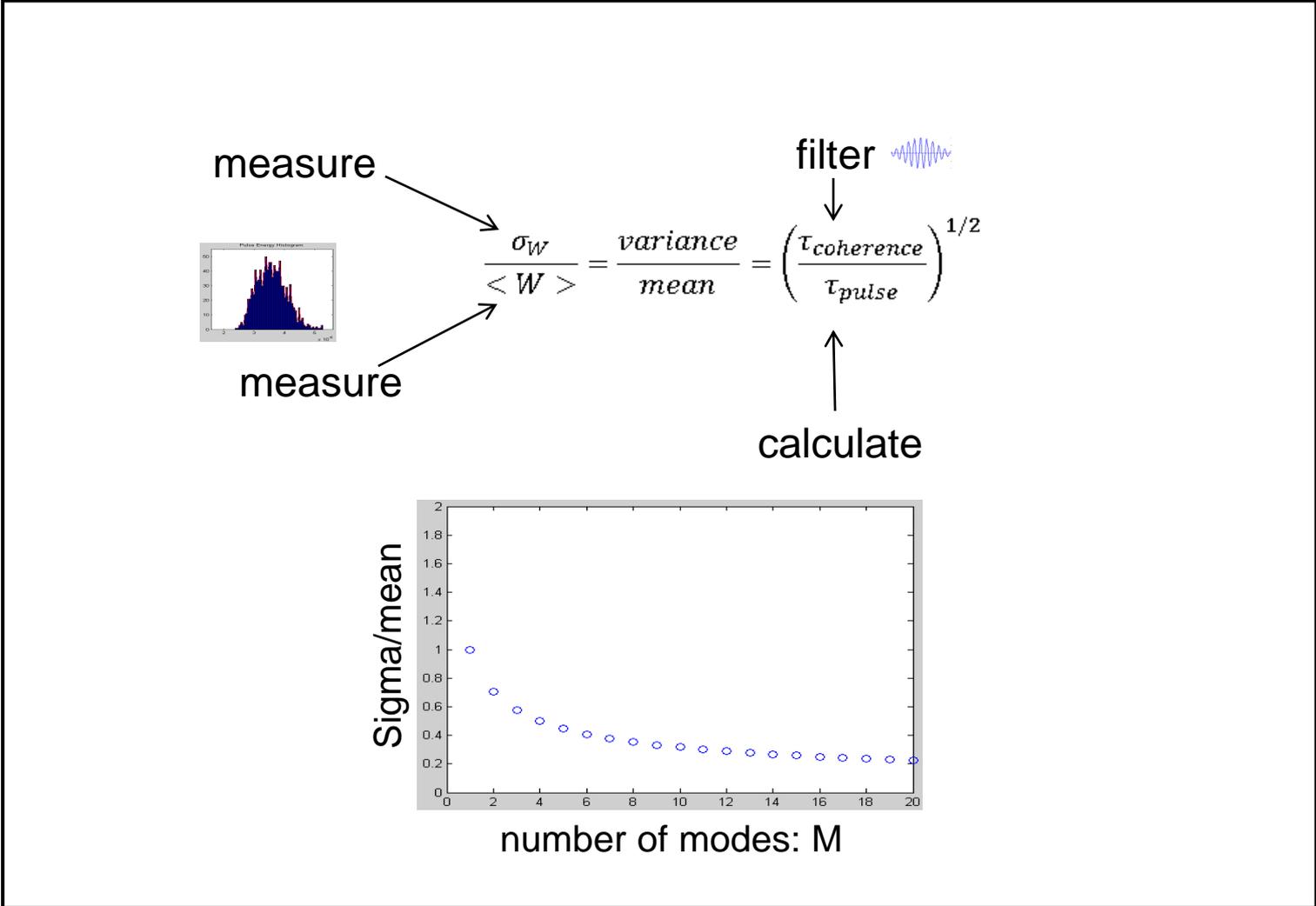


PROPERTIES OF INTEGRATED INTENSITY 243

Figure 6-1. Plots of M versus T/τ_c , exact solutions for Gaussian, Lorentzian, and rectangular spectral profiles.



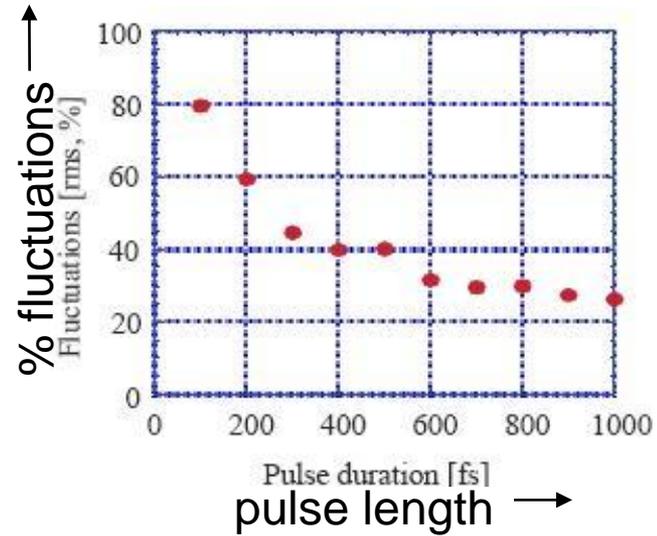
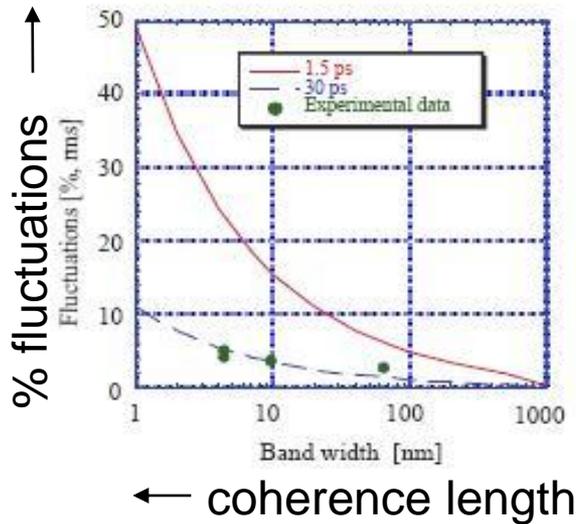
Relation between physics and measurement





Modes-per-pulse: Experimental Evidence, U. Tokyo

$$\frac{\sigma_W}{\langle W \rangle} = \frac{\text{variance}}{\text{mean}} = \left(\frac{\tau_{\text{coherence}}}{\tau_{\text{pulse}}} \right)^{1/2}$$





Time-Domain Measurements (cont'd)

In the simplest form...

$$\delta^2 = \frac{\sigma_W^2}{\overline{W}^2} = \int_{-T}^T \int_{-T}^T I(t)I(t') dt dt' \quad \text{fluctuations proportional to intensity correlation}$$

For Gaussian statistics and band pass filter $\delta^2 = \frac{1}{\sqrt{1 + 4\sigma_\tau^2 \sigma_\omega^2}}$

Expanding $\delta^2 \approx \frac{1}{2\sigma_\tau \sigma_\omega}$

For $\sigma_{coh} \approx 1/\sigma_\omega$

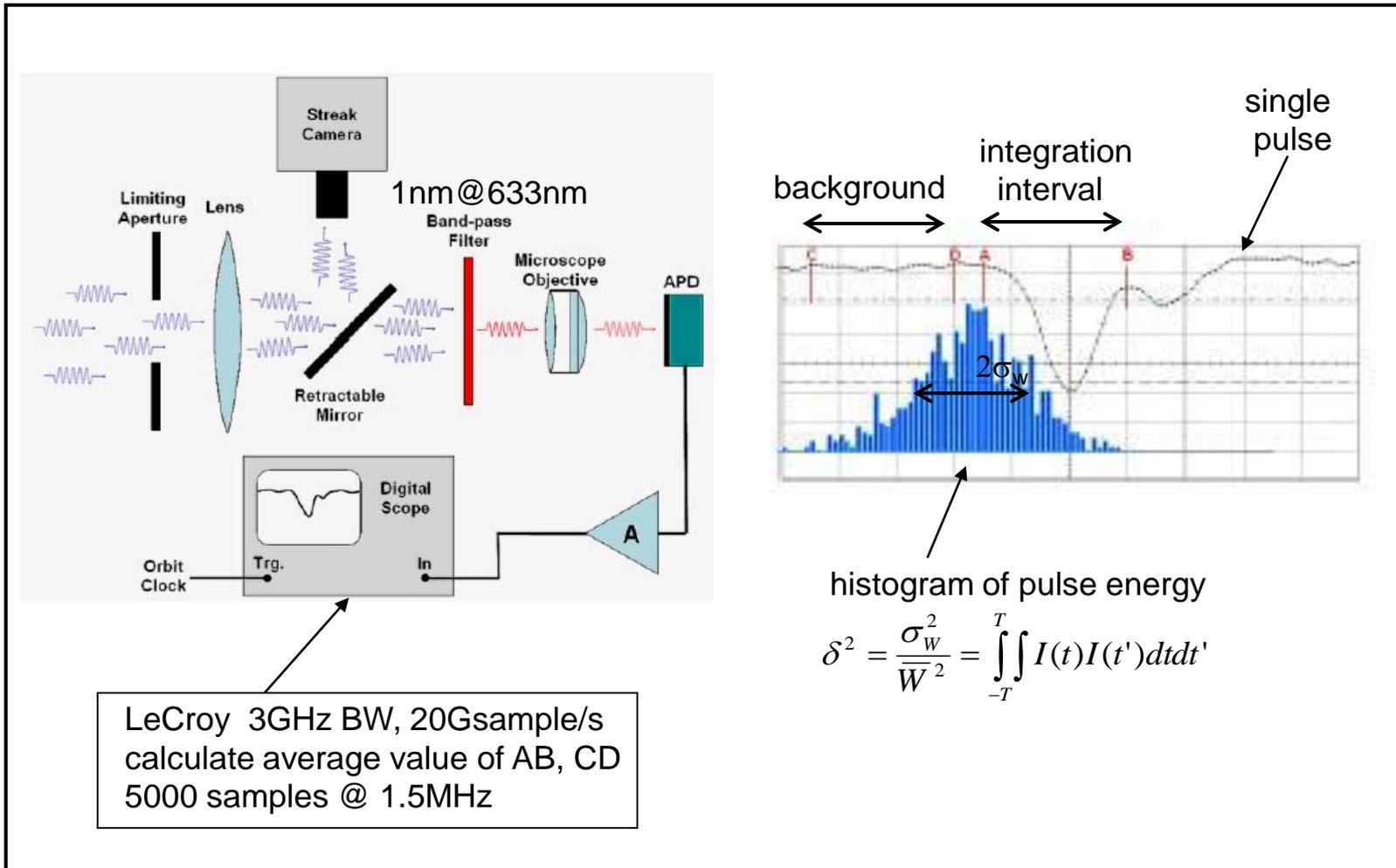
We get $\delta^2 \approx \frac{\sigma_{coh}}{\sigma_\tau} = \frac{1}{M}$

Make the coherence length long to reduce the number of modes M



Time-Domain Measurements at Berkeley

Intensity fluctuations, F. Sannibale, et al





Calibration against Streak Camera

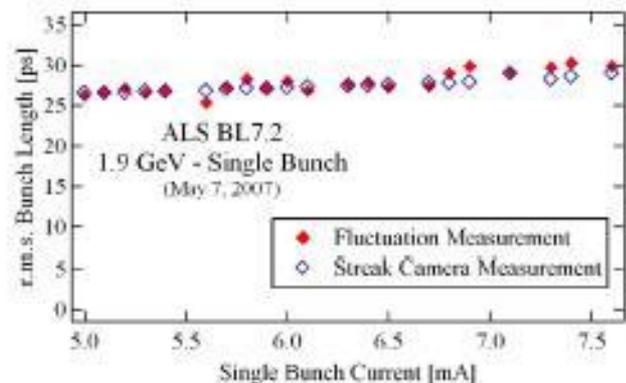
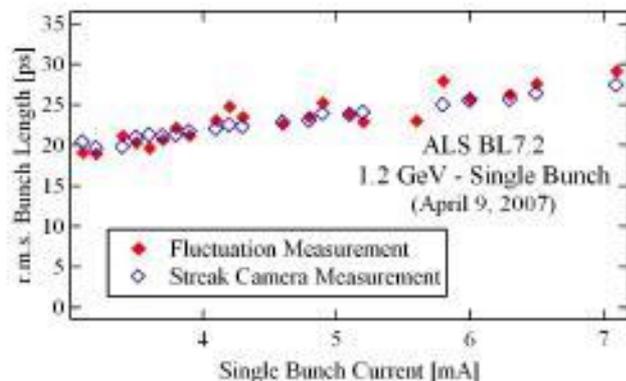


Figure 3: Examples of fluctuation and streak-camera bunch length measurements at the ALS for different beam parameters.

$$\delta^2 = \sqrt{1 + \frac{\sigma_\tau}{\sigma_{\tau,c}}} \sqrt{1 + \frac{\sigma_x}{\sigma_{x,c}}} \sqrt{1 + \frac{\sigma_y}{\sigma_{y,c}}}$$

$\sigma_{x/y,c}$ are transverse coherence sizes
-related to transverse EM modes at 633nm
-radiation process, including diffraction
-ratios about 2 and 0.1

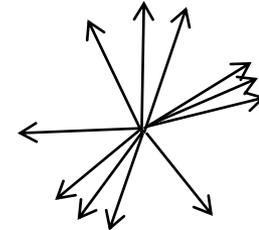
- also shot noise, photodiode noise



Frequency Domain View

Total electric field has a spectral content

$$f(t) = e(t) \sum \delta(t - t_i) \quad \tilde{E}(\omega) = \tilde{e}(\omega) \sum_{k=1}^N e^{i\omega t_k}$$



Phasors add up to 'spike' at frequencies ω

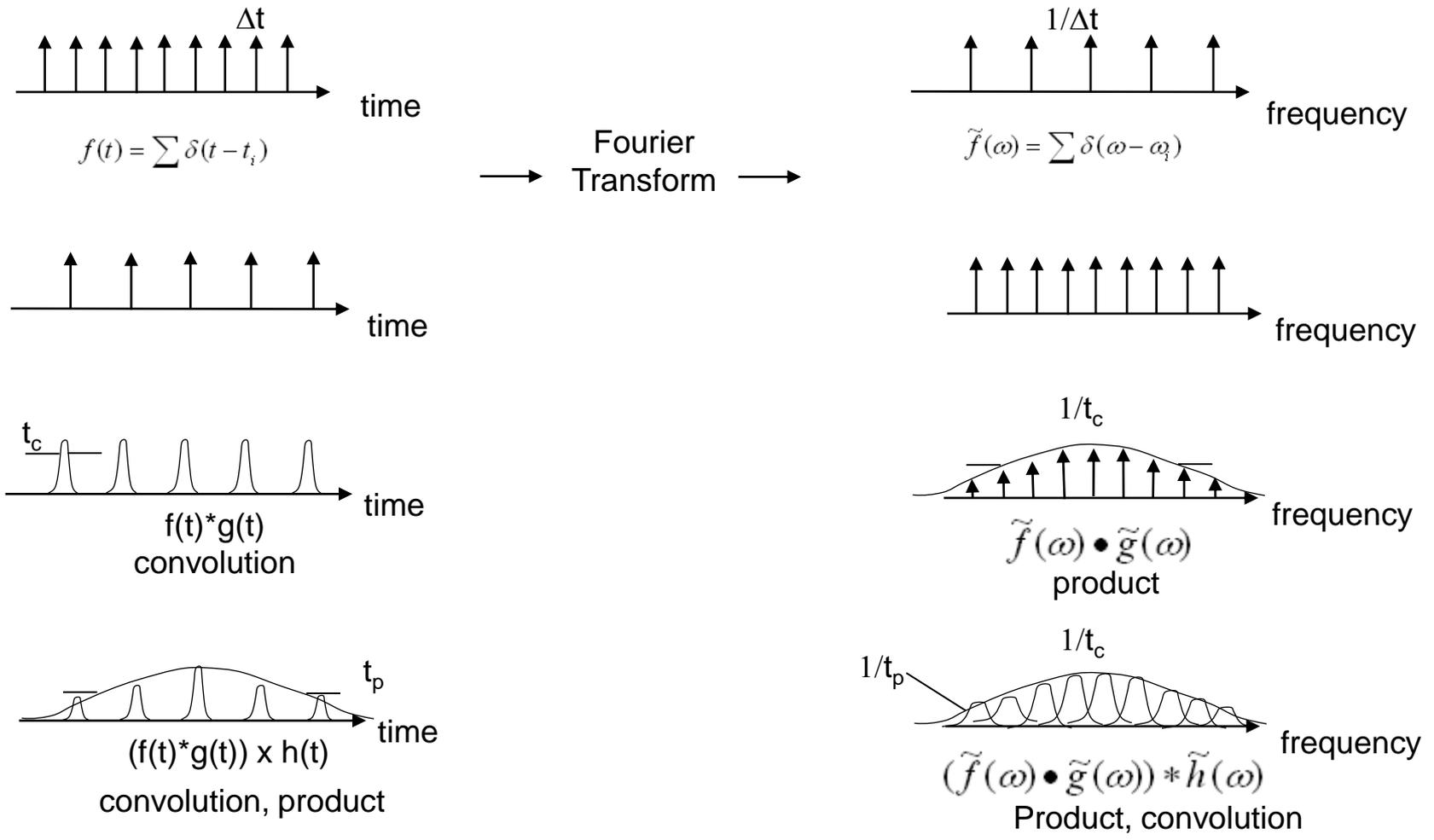
Shot-noise in wavepacket emission causes the spikes

In the frequency domain still have shot-to-shot fluctuations

Width of each spike is inversely proportional to the bunch length

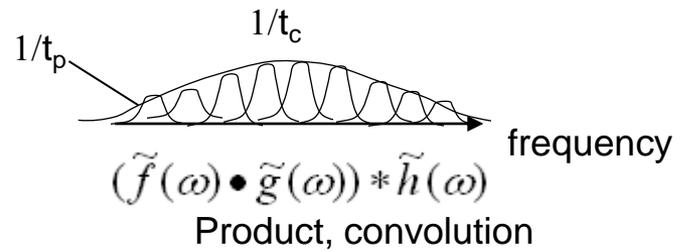
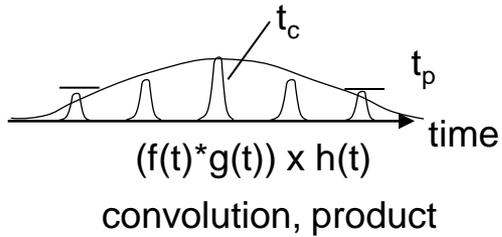


Frequency Domain: An Empirical Argument

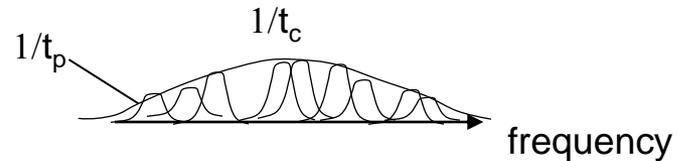
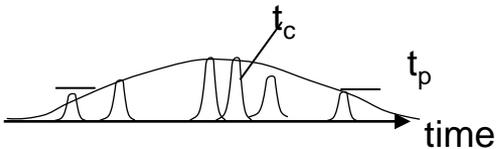




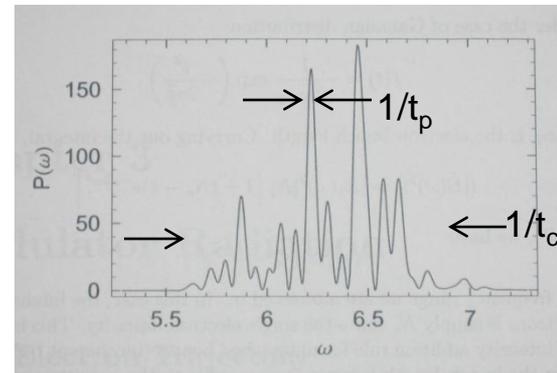
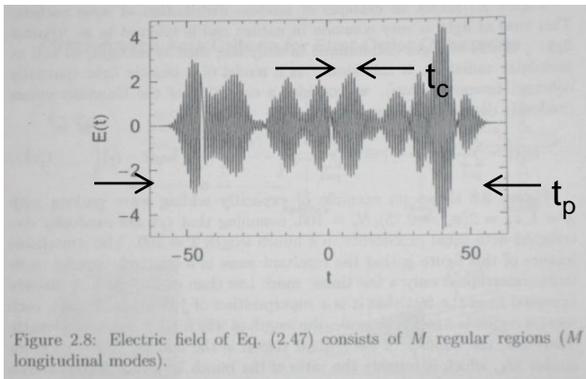
Empirical Argument (cont'd)



Now make a leap of faith to random emission...



In both domains we have constructive and destructive interference





Frequency Statistical Analysis

Experimentally can also analyze fluctuations in the frequency domain

Integrate the power spectrum of each pulse over frequency to find energy

$$\varepsilon = \int P(\omega) d\omega$$

The average energy is $\langle \varepsilon \rangle = \int (P(\omega)) d\omega$

And the variance is $\frac{\langle \Delta \varepsilon^2 \rangle}{\langle \varepsilon \rangle^2} = \frac{1}{\langle \varepsilon \rangle^2} \iint \langle [P - \langle P \rangle] \cdot [P' - \langle P' \rangle] \rangle d\omega d\omega'$

or

$$\frac{\langle \Delta \varepsilon^2 \rangle}{\langle \varepsilon \rangle^2} = \frac{1}{\langle \varepsilon \rangle^2} \iint \langle |PP'| \rangle - \langle P \rangle \langle P' \rangle d\omega d\omega'$$

Need to compute $\langle P \rangle$ and 4th order field correlation $\langle PP' \rangle$ to evaluate variance



Frequency Domain Analysis (cont'd)

$$\frac{\langle \Delta \varepsilon^2 \rangle}{\langle \varepsilon \rangle^2} = \frac{1}{\left(\int T d\omega \right)^2} \iint TT' |f(\omega - \omega')|^2 d\omega d\omega'$$

filter

bunch transform

For a Gaussian filter and a Gaussian bunch

$$\frac{\langle \Delta \varepsilon^2 \rangle}{\langle \varepsilon \rangle^2} = \frac{1}{\sqrt{1 + 4\Delta\omega^2 \sigma_b^2}}$$

$$\frac{\langle \Delta \varepsilon^2 \rangle}{\langle \varepsilon \rangle^2} = \frac{\tau_c}{\tau_b}$$

(same as time-domain analysis)

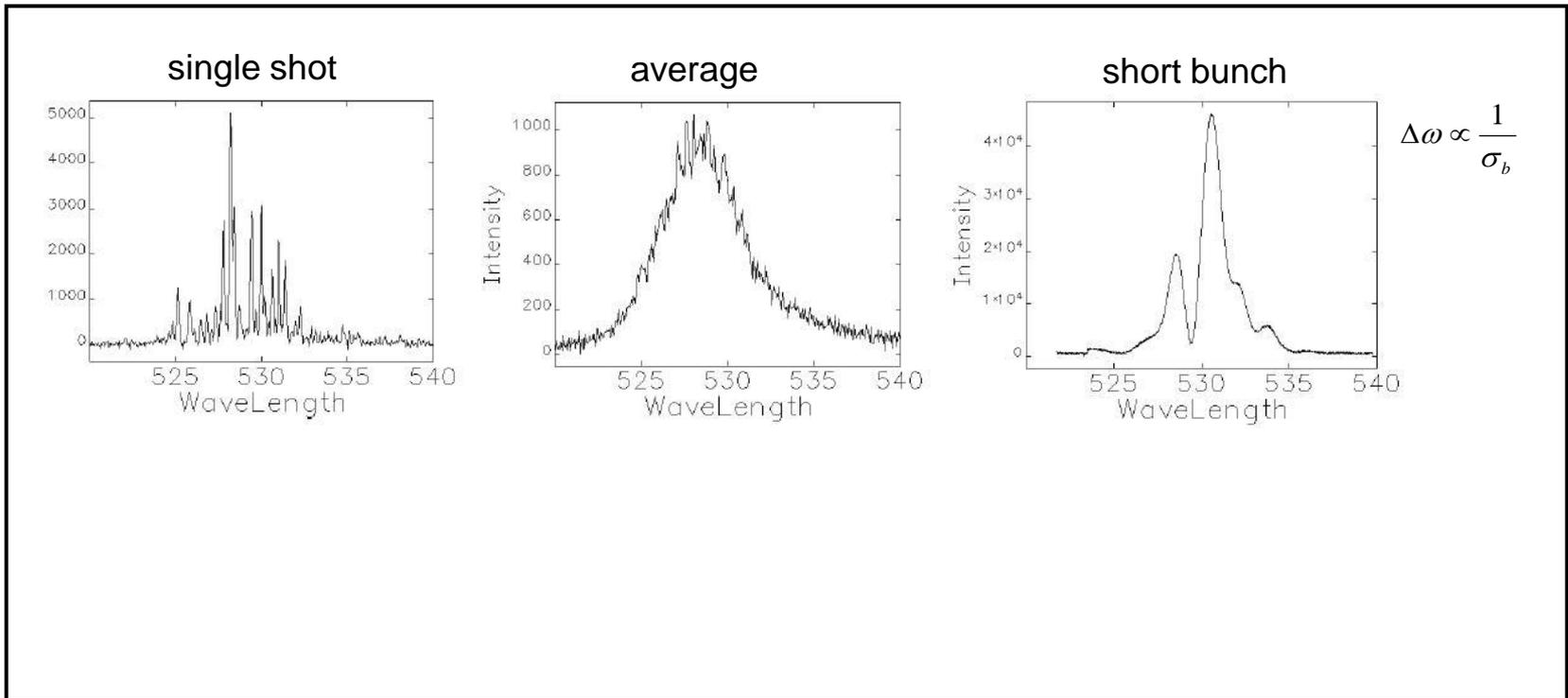
filter characteristic

bunch characteristic



Frequency Domain Experiments at APS

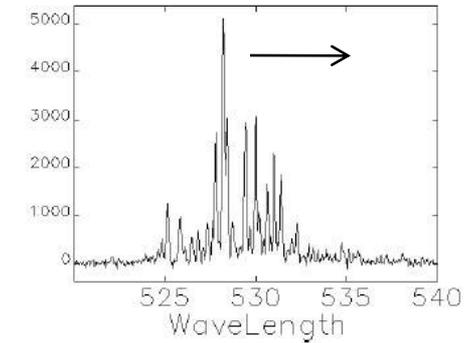
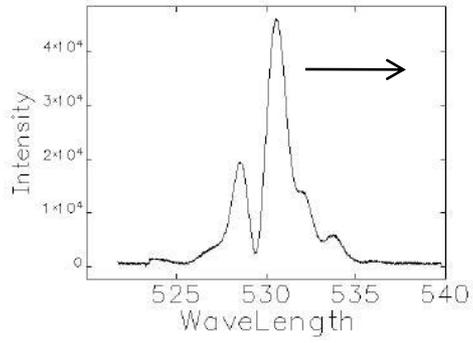
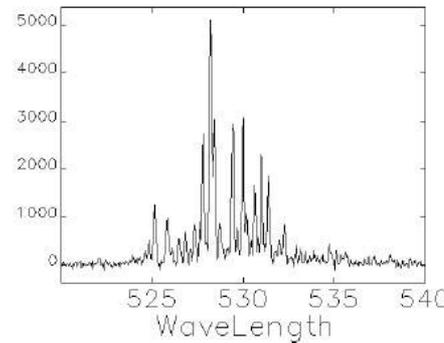
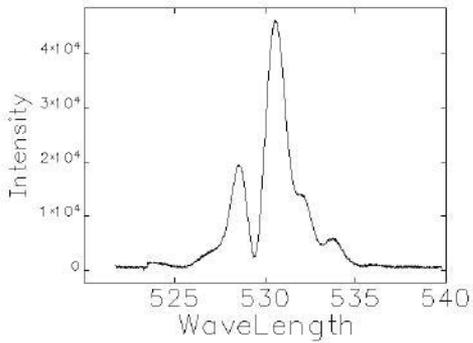
Use a spectrometer to observe spikes in single-shot spectrum
Sajaev, Argonne Nat'l Labs





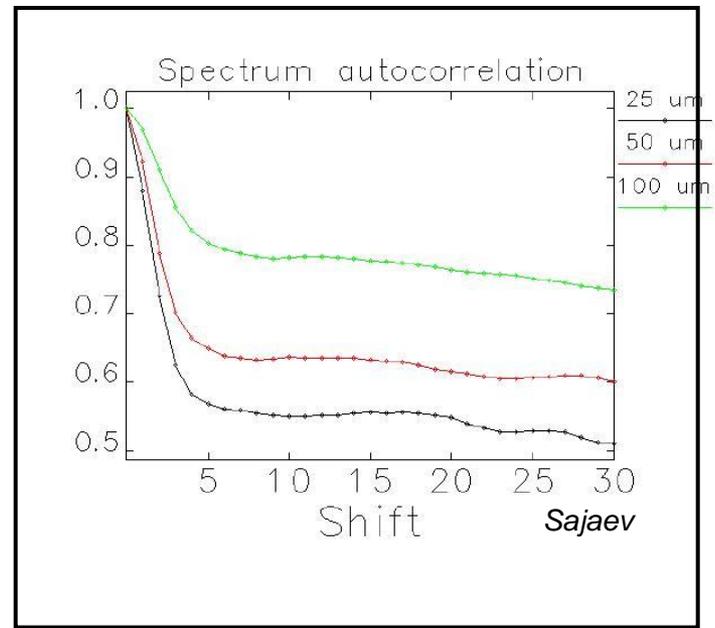
Frequency Domain Experiments (cont'd)

Bunch length proportional to Fourier transform of spectrum autocorrelation



Large frequency correlation
Short bunch length

Small frequency correlation
long bunch length





Fluctuations in Interference Visibility Pattern

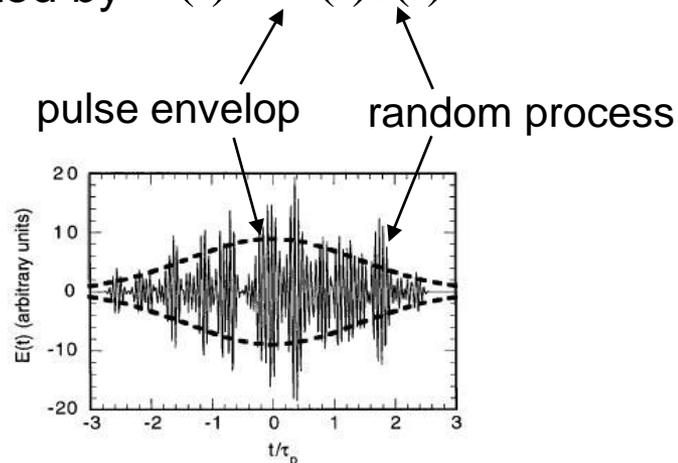
Landmark paper : Zolotarev and Stupakov (1996)

Measure fluctuations in the coherence function of the incoherent electric field

$$\Gamma(\tau) = \int E(t)E^*(t - \tau)dt$$

Utilizes a two-beam interferometer to measure $\Gamma(\tau)$

In simulation, the electric field is represented by $E(t) = A(t)e(t)$





USPAS Simulator - Pulse Energy Fluctuations

Each pulse of light is a superposition of randomly-phased 'wave packets'

Simulator generates wave packets at random times t_k

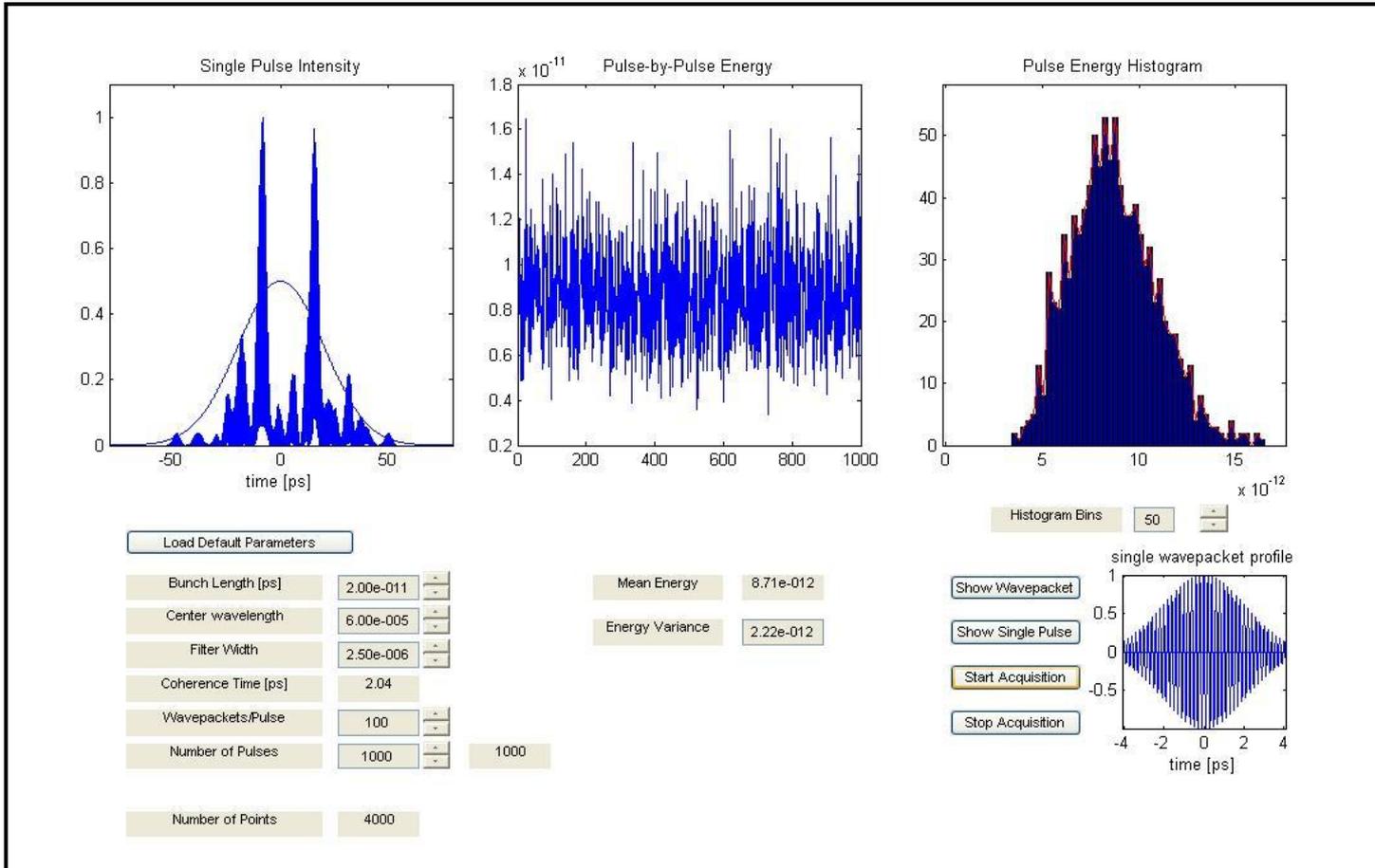
Computes wavepacket superposition and resulting intensity E^*E

Records statistics of shot-to-shot photon beam energy $U = \int E^* E dt$
to deduce pulse length

Very much like Sinnabale experiment and USPAS laboratory
but you 'see' effects not physically observable



Simulator for Pulse-Energy Fluctuations





USPAS Simulator (cont'd)

Part I: Photon beam properties

Calculate wavelength, energy, photon flux, etc.

Part II: Coherence properties

Coherence length with BP filter, etc

Part III: Time-base calculations for simulator code

Need simulate with 1 μ m radiation

Part IV: The simulator interface

Part V: Wavepackets

Study as a function of wavelength, bandwidth, etc

Part VI: Study pulse-to-pulse statistics as a function of bunch length, filter width, etc

Independent study



Summary Fluctuation Techniques

- Wavepacket emission is a statistically random process
- In the time domain
 - use a filter to make coherence length~bunch length
 - look for fluctuations in shot-to-shot intensity
- Fluctuations in interferometer visibility pattern
- In the frequency domain
 - use a spectrometer to observe fluctuations in spectra
- Simulator for this afternoon