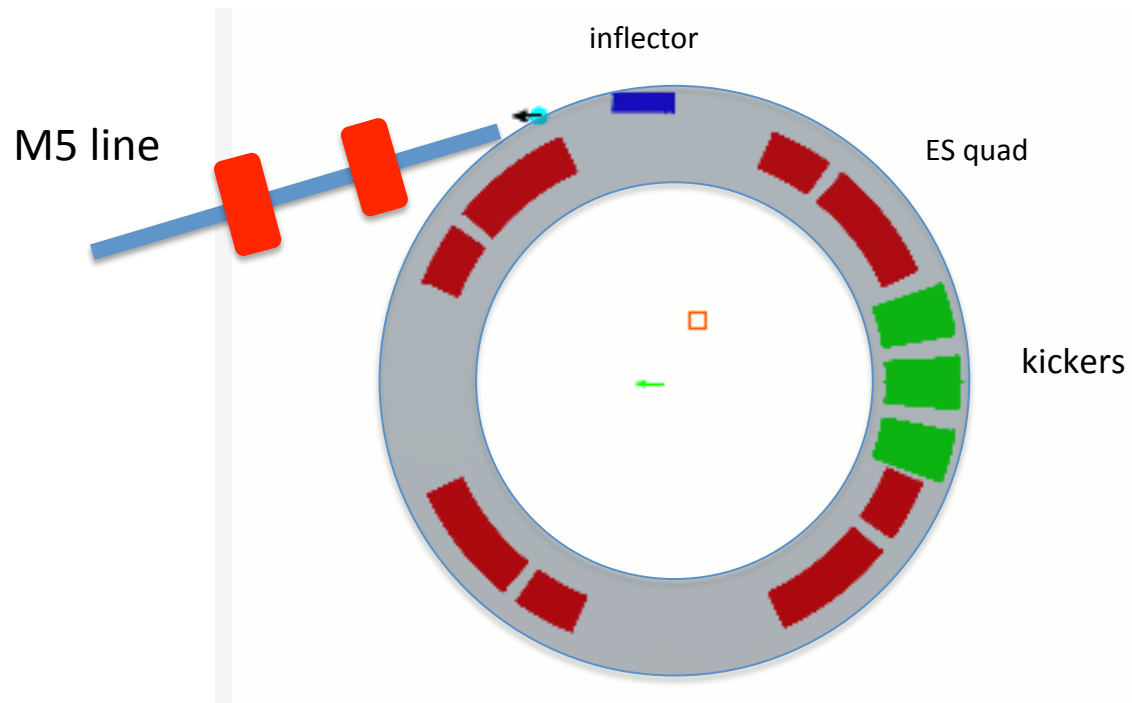
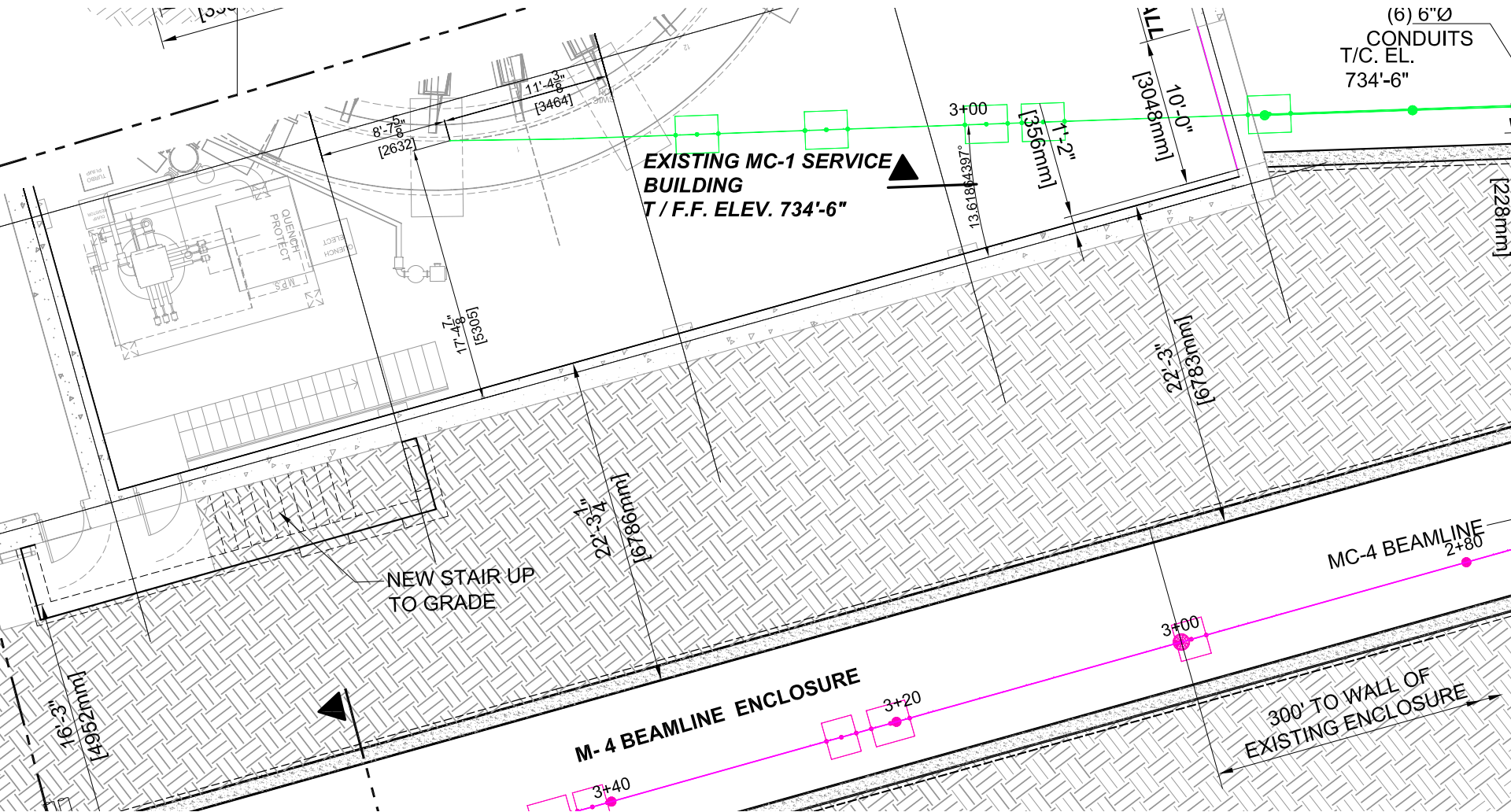


Injection

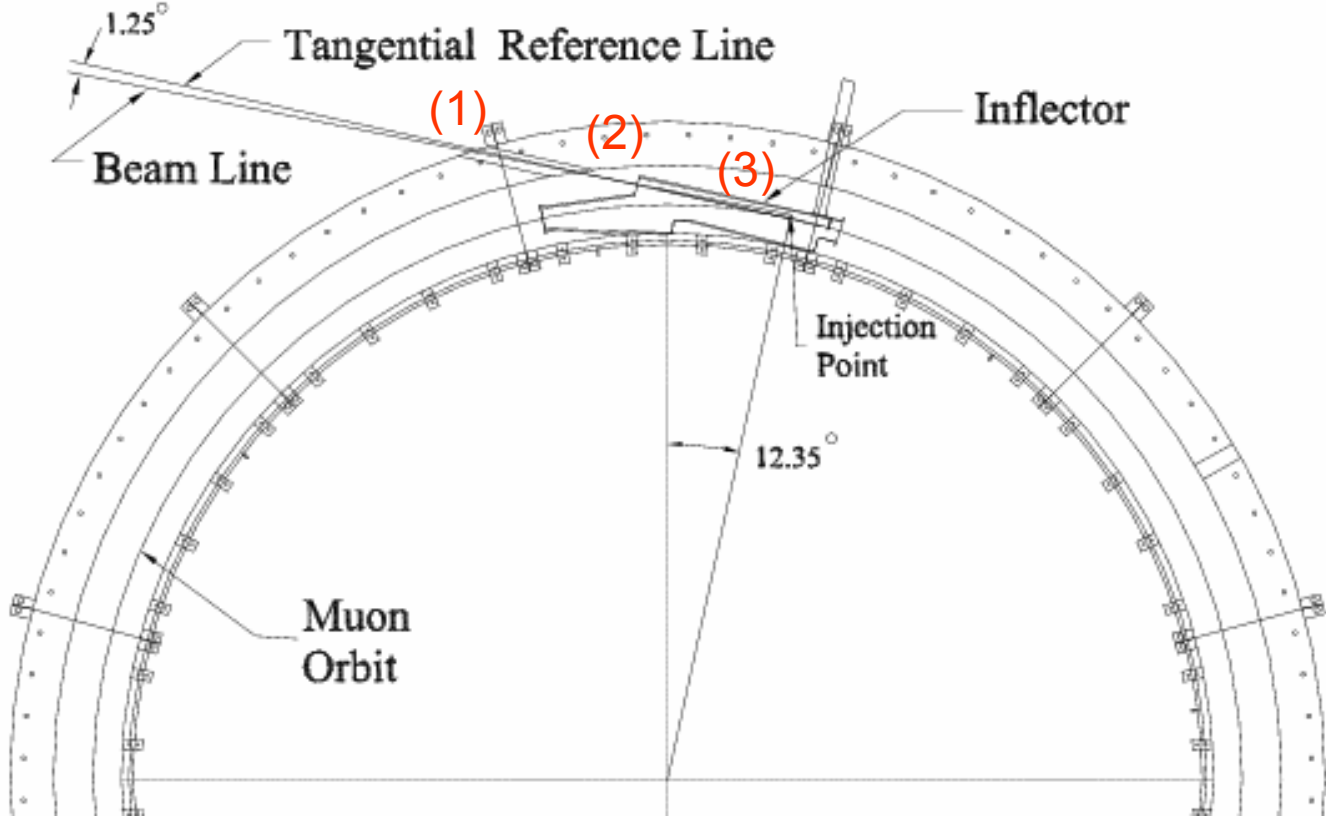
D. Rubin
USPAS
2019

Muon Storage Ring

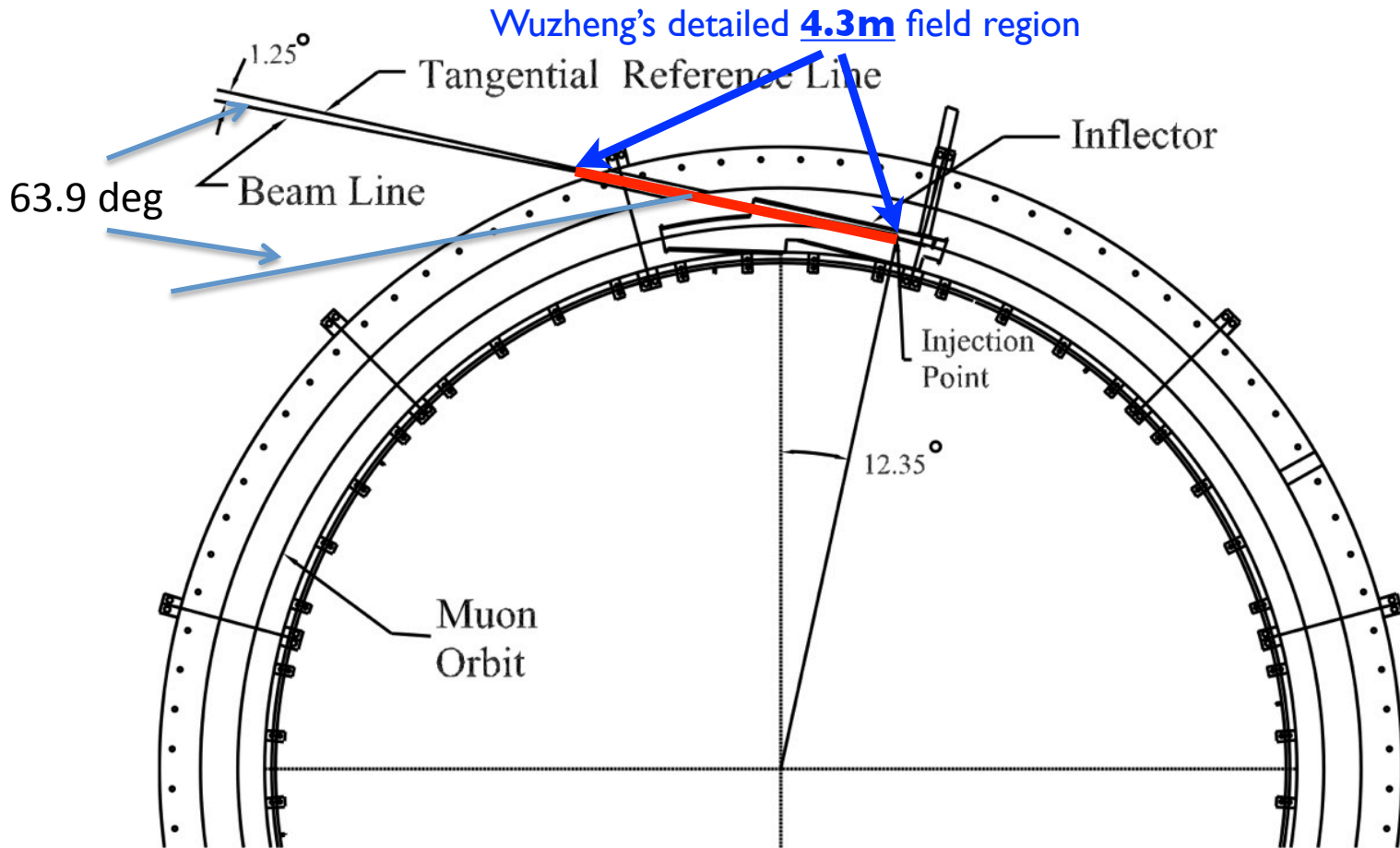




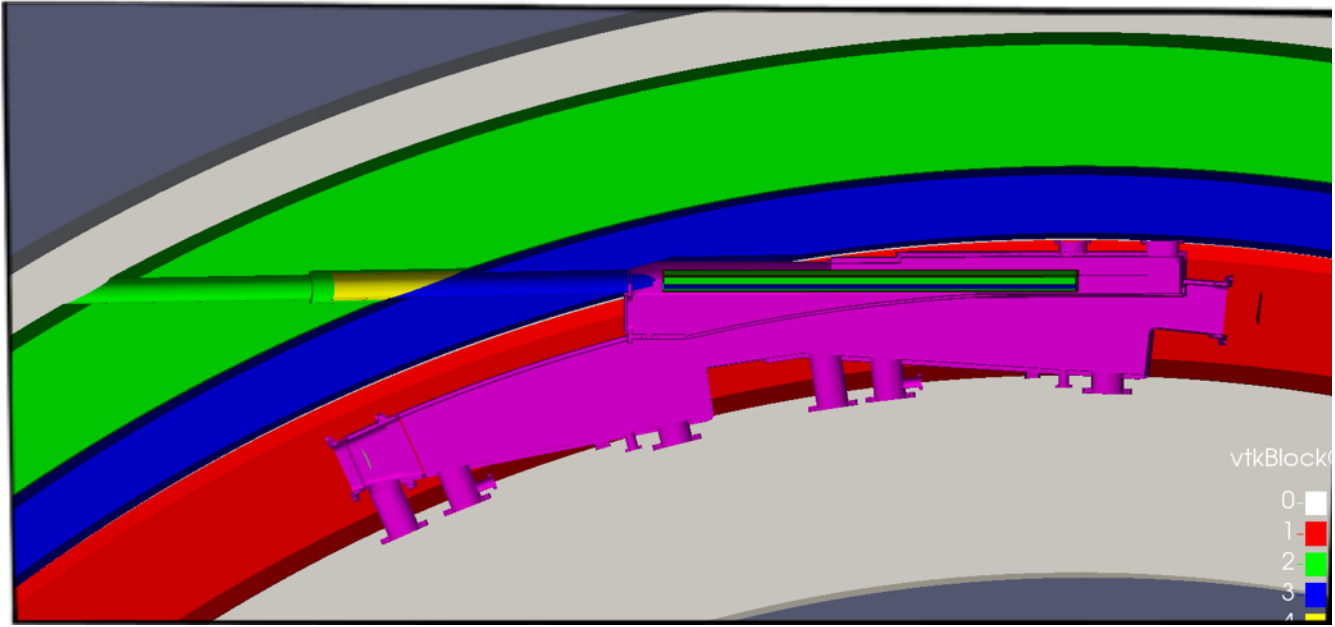
How do twiss parameters propagate through iron, cryostat, inflector into ring



(1) hole in back leg, (2) storage ring fringe field, (3) inflector channel

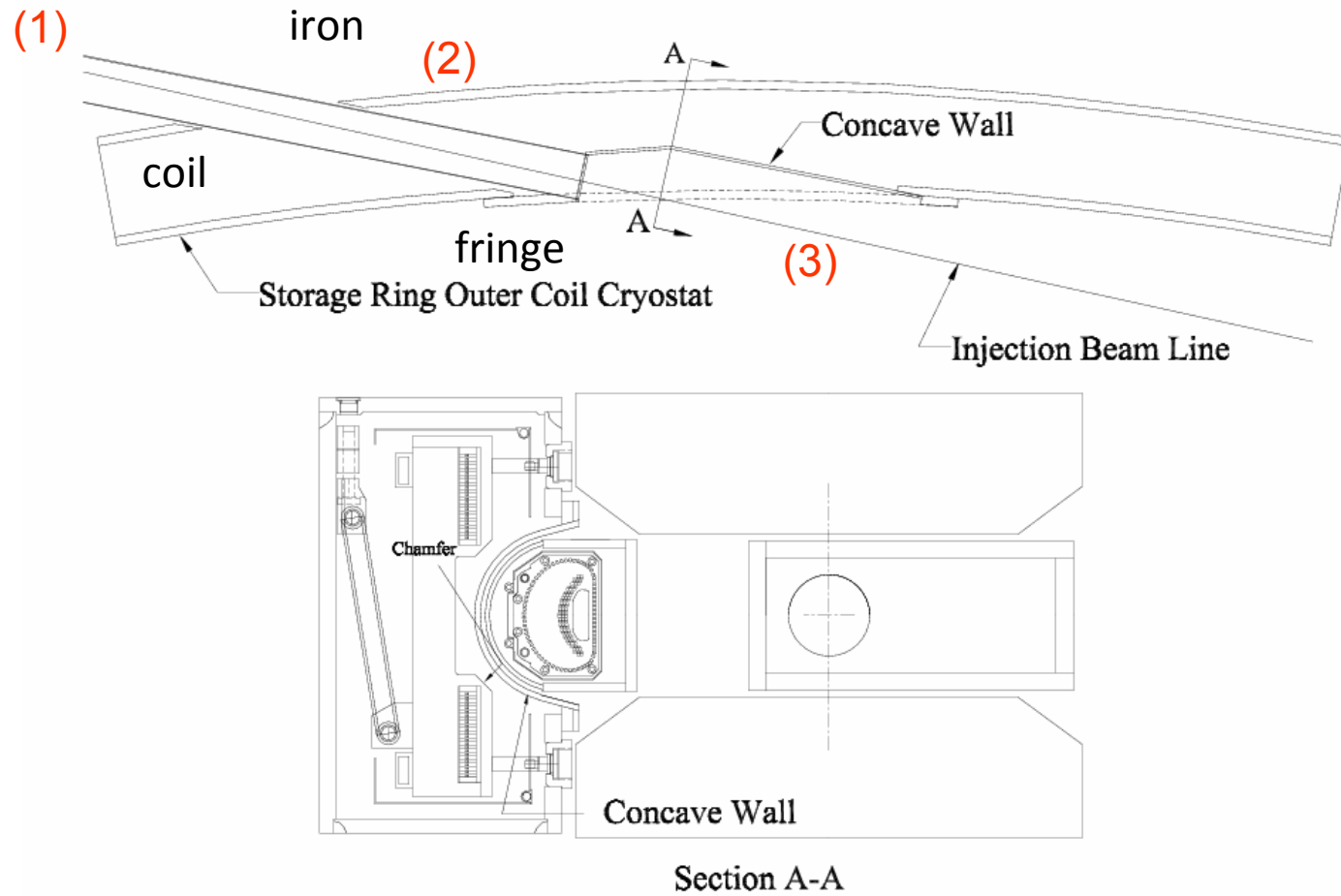


Effective gradient $G \approx (1 + \cos(63.9))/2 (\Delta B_y / \Delta R)$

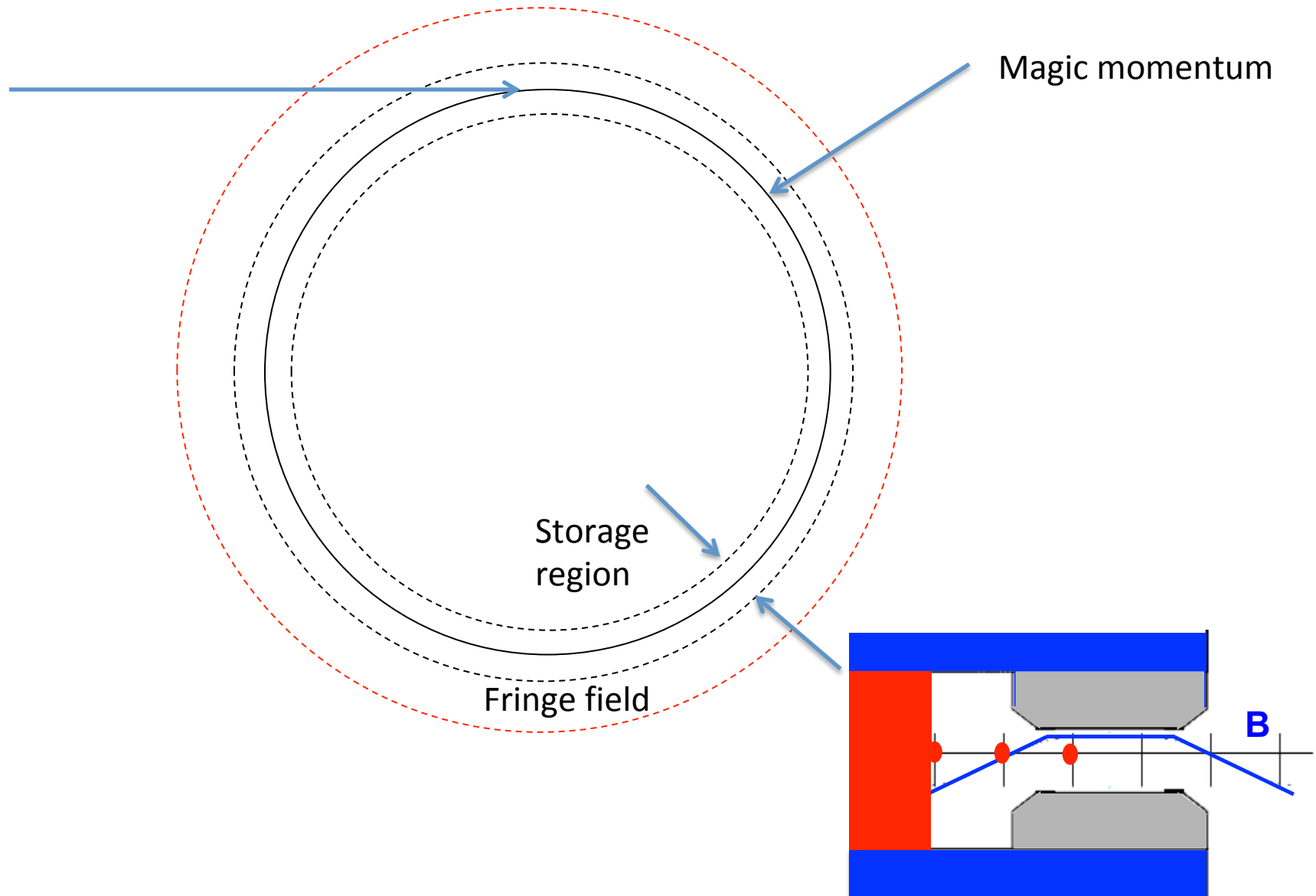


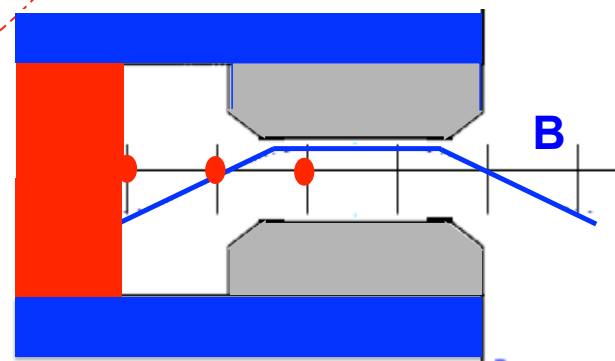
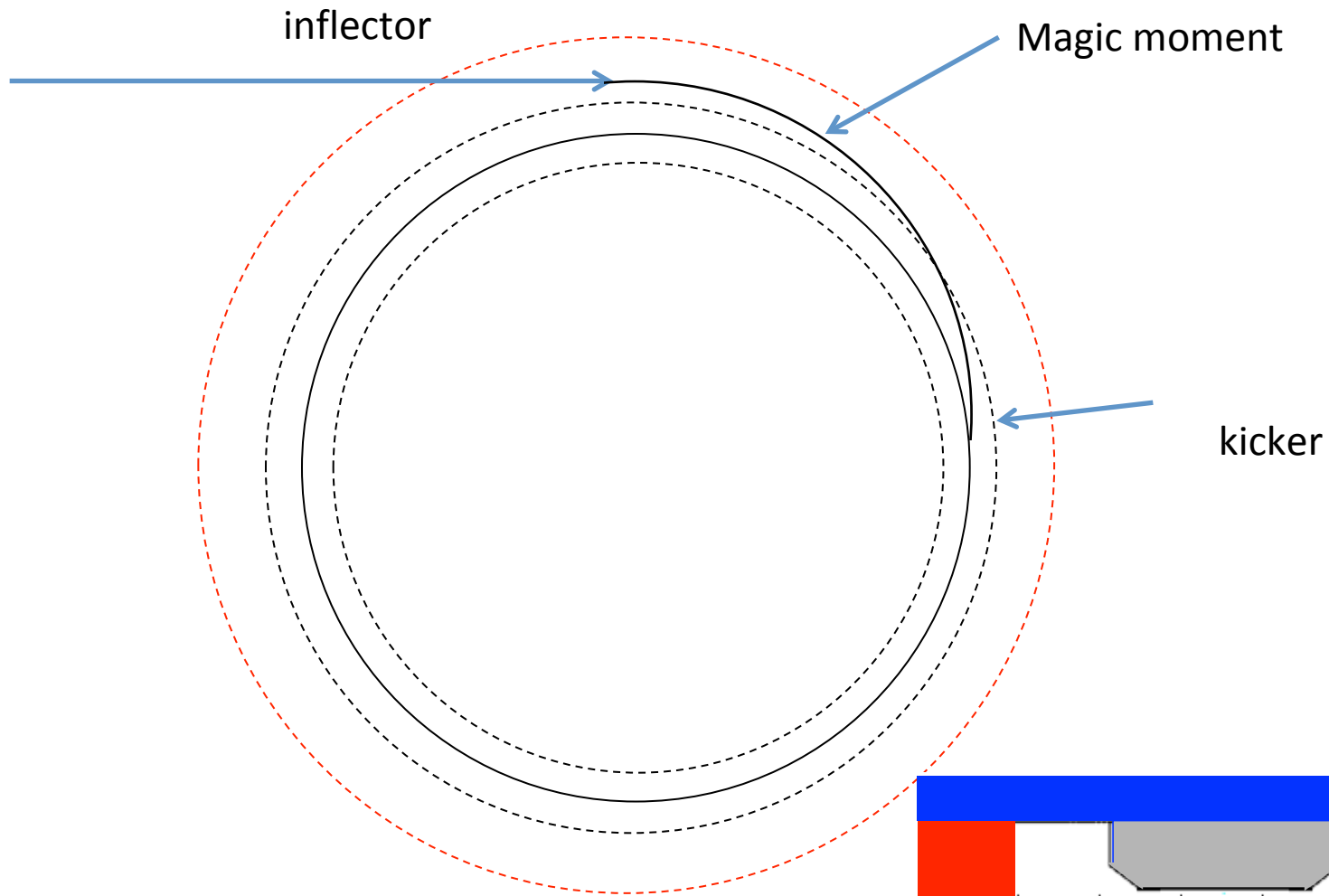
What are the fields along the route of injected beam ?

Another View

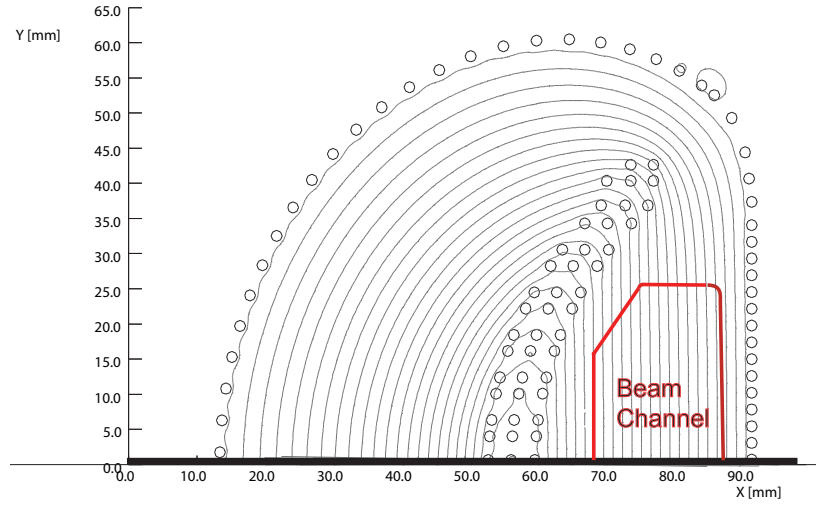
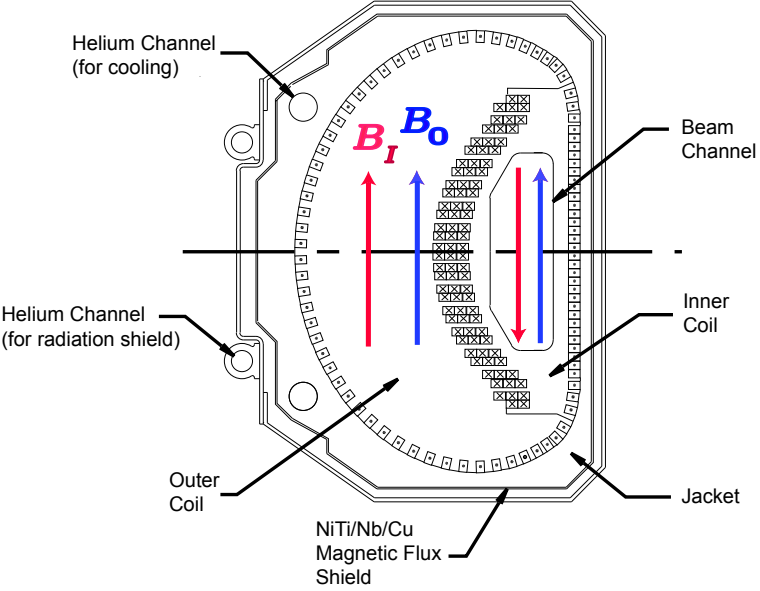


From W. Meng, June 2008 collaboration meeting





Inflector

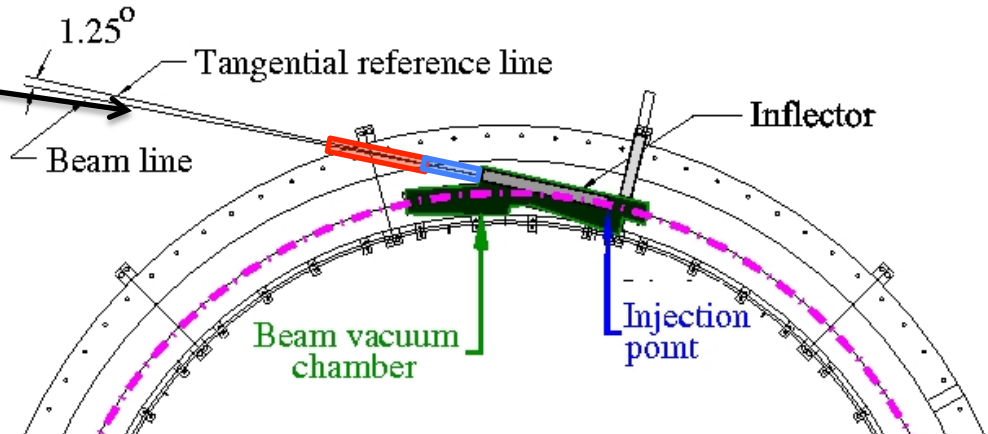


Injection Channel

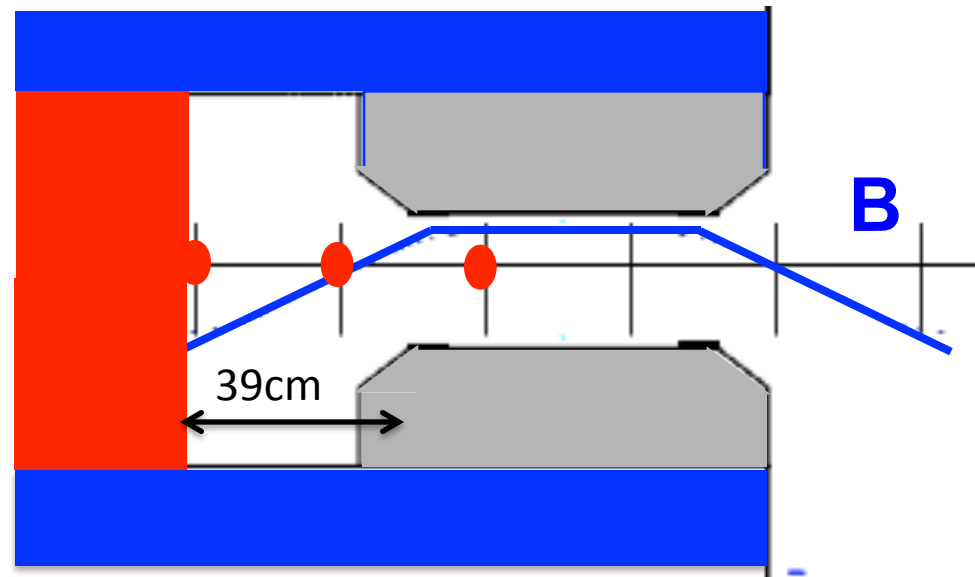
Muons come to the end of the M5 line and then propagate through:

- Hole in magnet yoke
- Dipole fringe field
- Inflector

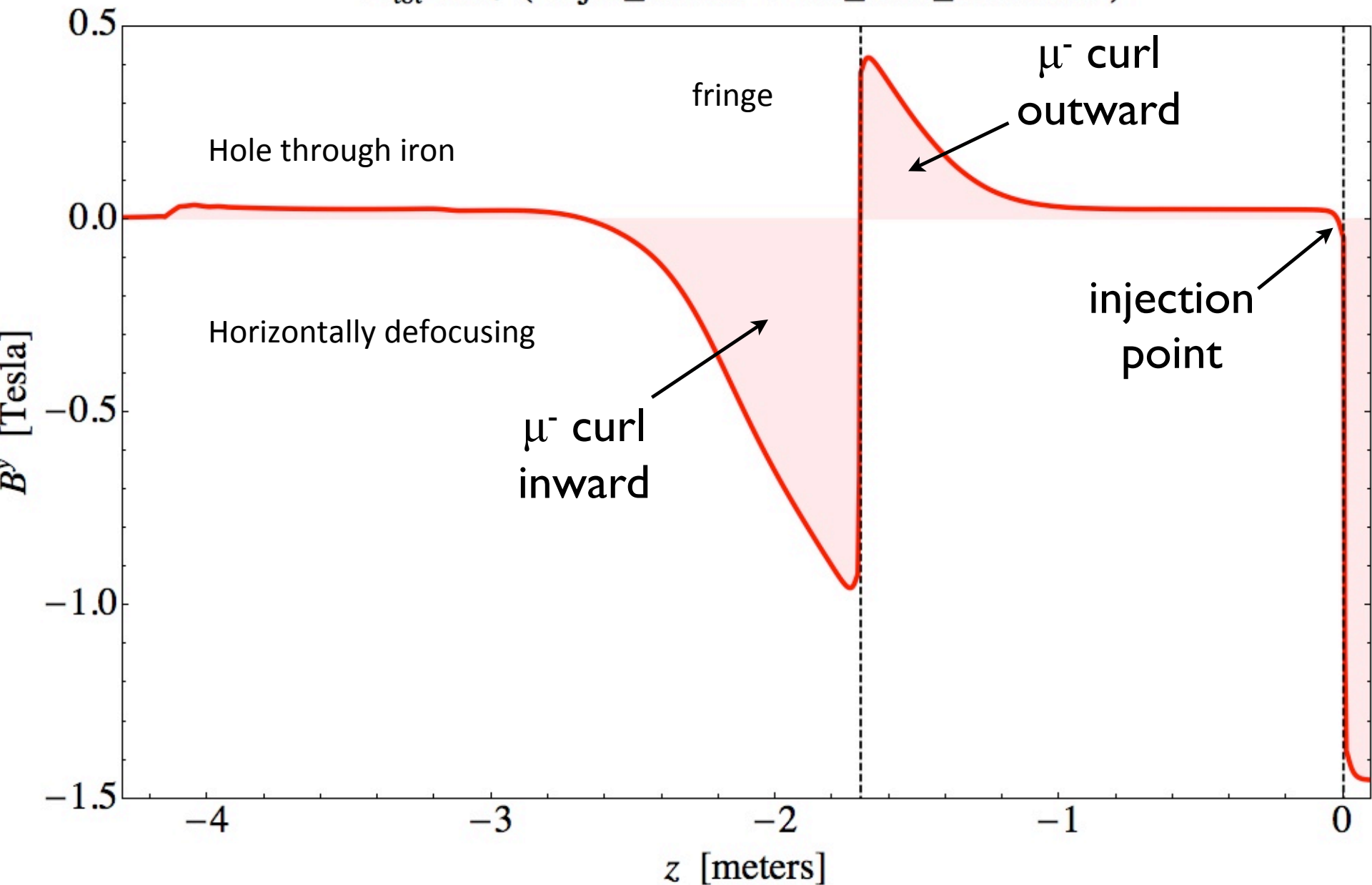
And exit the inflector 77 mm from the center of the dipole aperture



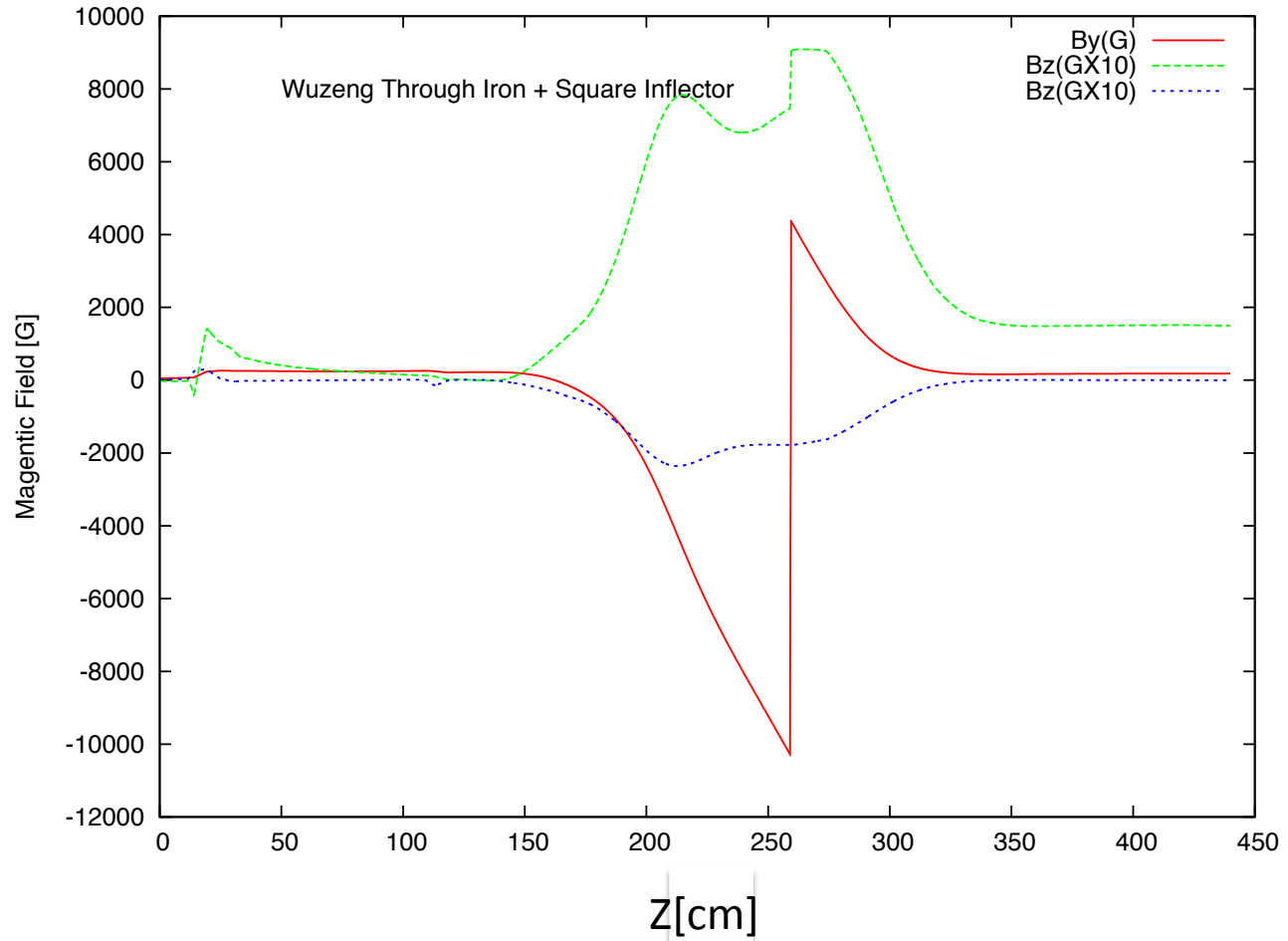
The magnetic field is near zero at the inner surface of the yoke, and rises to 1.45T between the magnet poles, over a distance of $\sim 39\text{cm}$



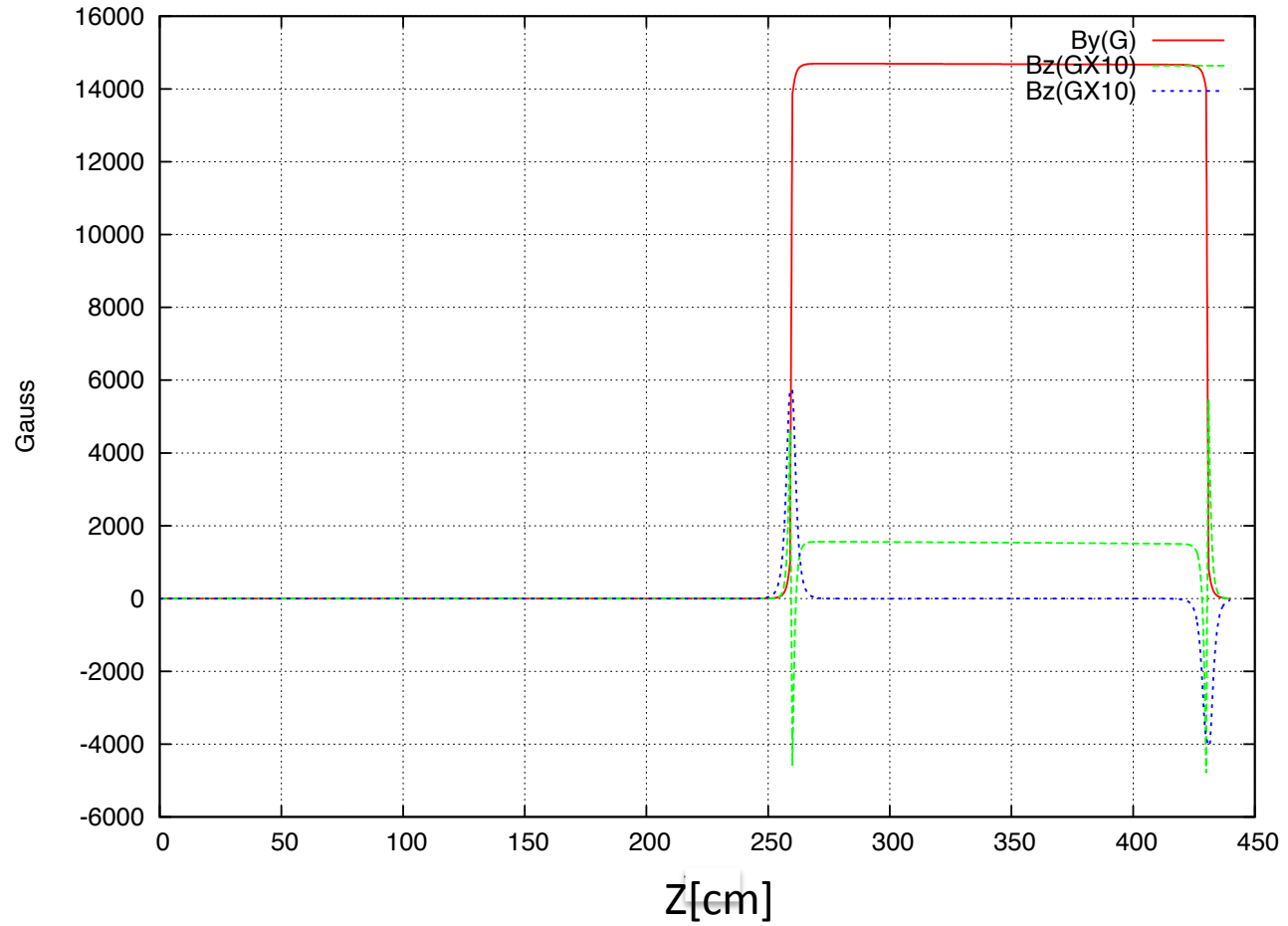
B_{tot}^y vs. z (“injec_fld.dat”+“inf_field_alone.dat”)



Inflector and main dipole

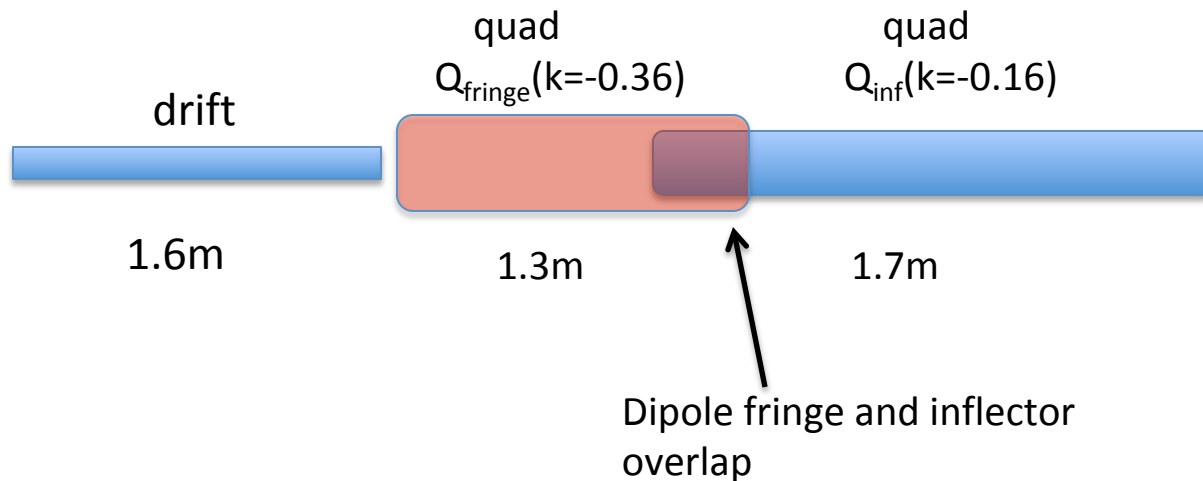


Inflector



Injection channel defocusing

The 4.3m injection channel can be modeled as follows

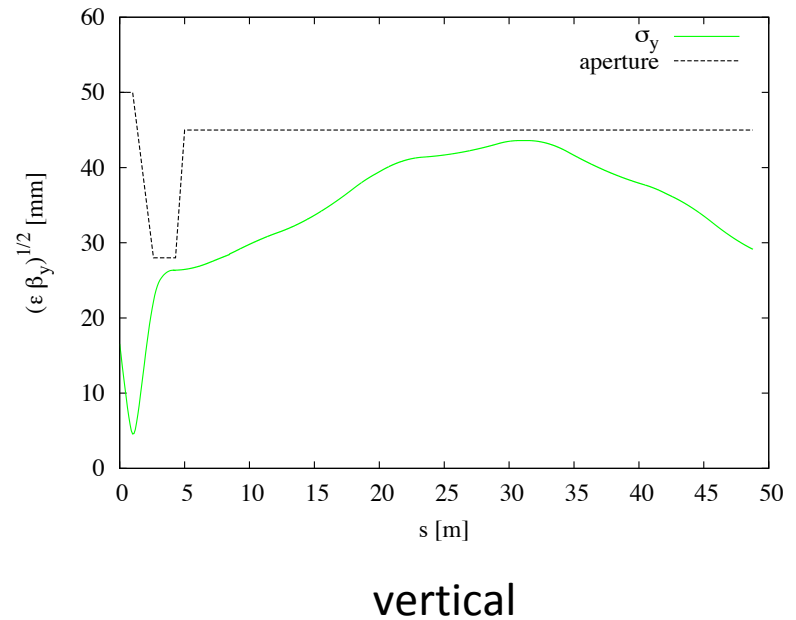
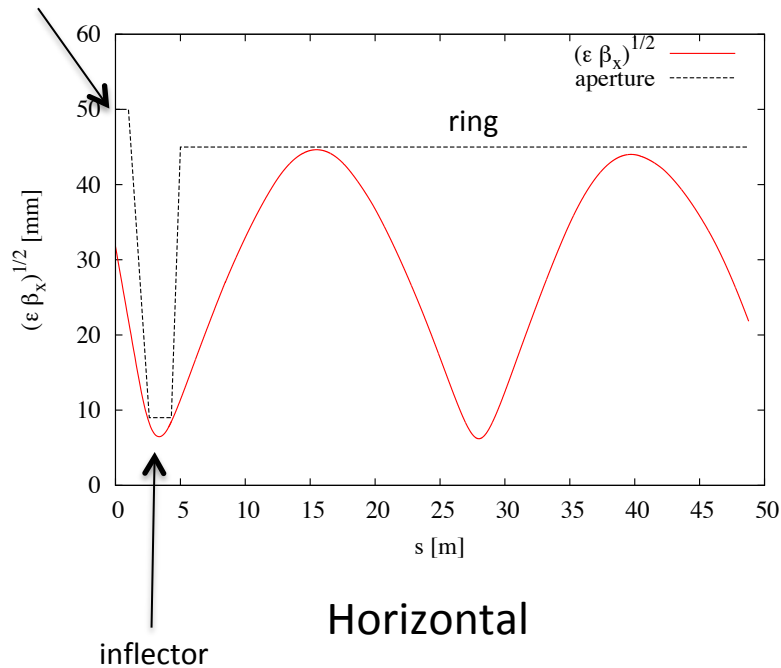


Or we can compute the transfer maps numerically by tracking

Acceptance

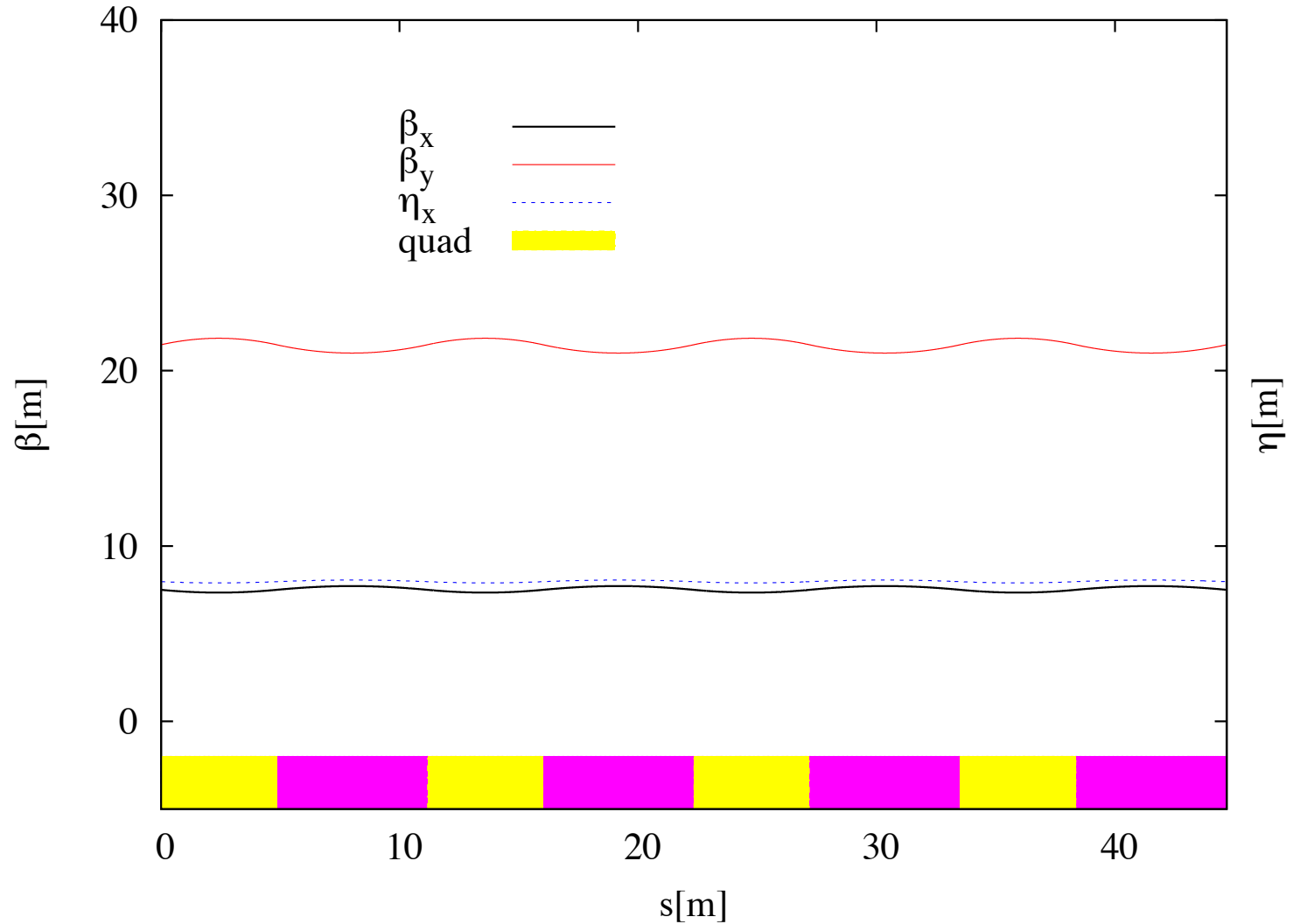
β, α, η at end of M5 line (entrance to iron) are chosen so that inflector and ring have equivalent acceptance

Entrance
to iron



Aperture and beam size

Twiss parameters (A-matrix) is chosen to be single valued in ring



Acceptance

We maximize transmission and capture by focusing the beam so that inflector and ring have equivalent acceptance

$$\sqrt{\frac{\beta_{\text{inf}}}{\beta_{\text{max-ring}}}} = \frac{A_{\text{inf}}}{A_{\text{ring}}} \rightarrow \beta_{\text{inf}} = \left(\frac{A_{\text{inf}}}{A_{\text{ring}}} \right)^2 \beta_{\text{max-ring}}$$

The maximum β in the ring depends on the matching according to

$$\beta_{\text{max-ring}} \sim \frac{\beta_{\text{ring}}^2}{\beta_{\text{inf}}} \quad \left(\begin{array}{l} \beta_{\text{ring}}^x = 8\text{m} \\ \beta_{\text{ring}}^y = 16\text{m} \end{array} \right)$$

Then
$$\beta_{\text{inf}} = \frac{A_{\text{inf}}}{A_{\text{ring}}} \beta_{\text{ring}}$$

We find that inflector and ring have equivalent acceptance if

$$\beta_{\text{inf}}^x = 1.6\text{m}, \quad \beta_{\text{inf}}^y = 10.0\text{m}$$

Dispersion

Energy acceptance of inflector and ring are equivalent if

$$\eta_{\text{inf}} = \frac{A_{\text{inf}}}{A_{\text{ring}}} \eta_{\text{max-ring}} \quad \text{and since} \quad \eta_{\text{max-ring}} = 2\eta_{\text{ring}} - \eta_{\text{inf}}$$

$$\Rightarrow \eta_{\text{inf}} \sim 2.5\text{m}$$

(Tracking simulations indicate optimal capture for $\eta \sim 0$)

The energy acceptance follows from

$$\sigma = \sqrt{\beta_x \epsilon + \eta_{\text{max-ring}}^2 \delta^2}$$

$$\Rightarrow \delta \sim 0.1\%$$

$$\begin{pmatrix} \beta_{\text{max-ring}}^x = & 40\text{m} \\ \eta_{\text{max-ring}} = & 17\text{m} \\ \epsilon_x = & 40\mu\text{m} \end{pmatrix}$$

Dispersion

Energy acceptance of inflector and ring are equivalent if

$$\eta_{\text{inf}} = \frac{A_{\text{inf}}}{A_{\text{ring}}} \eta_{\text{max-ring}} \quad \text{and since} \quad \eta_{\text{max-ring}} = 2\eta_{\text{ring}} - \eta_{\text{inf}}$$

$$\Rightarrow \eta_{\text{inf}} \sim 2.5\text{m}$$

Tracking simulations indicate optimal capture for

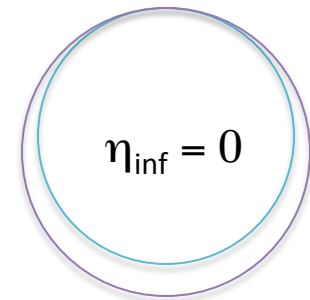
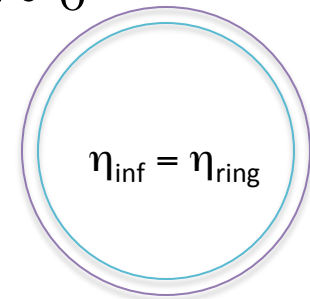
$$\eta \sim 0$$

The energy acceptance follows from

$$\sigma = \sqrt{\beta_x \epsilon + \eta_{\text{max-ring}}^2 \delta^2}$$

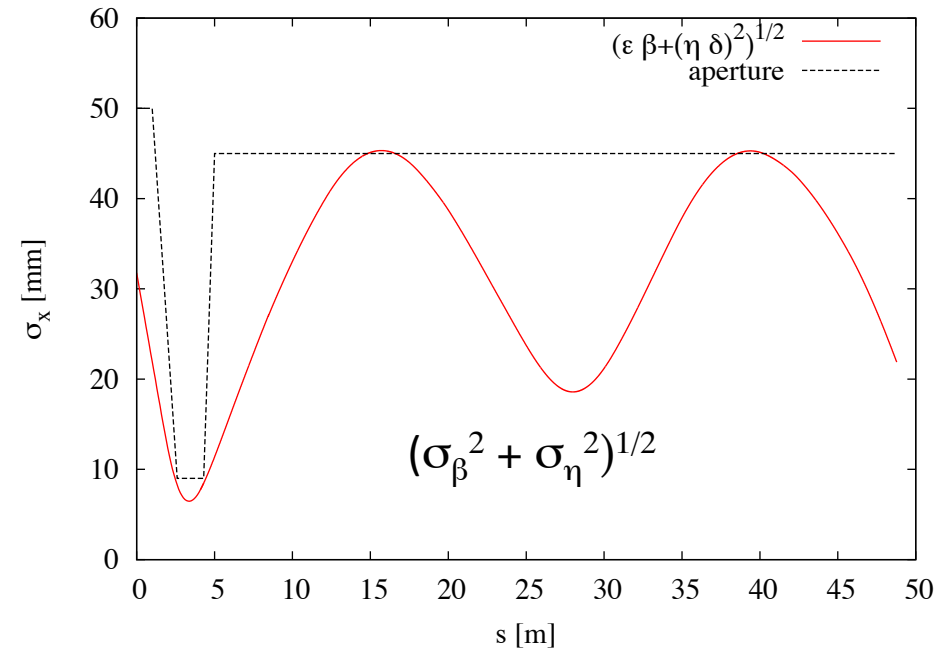
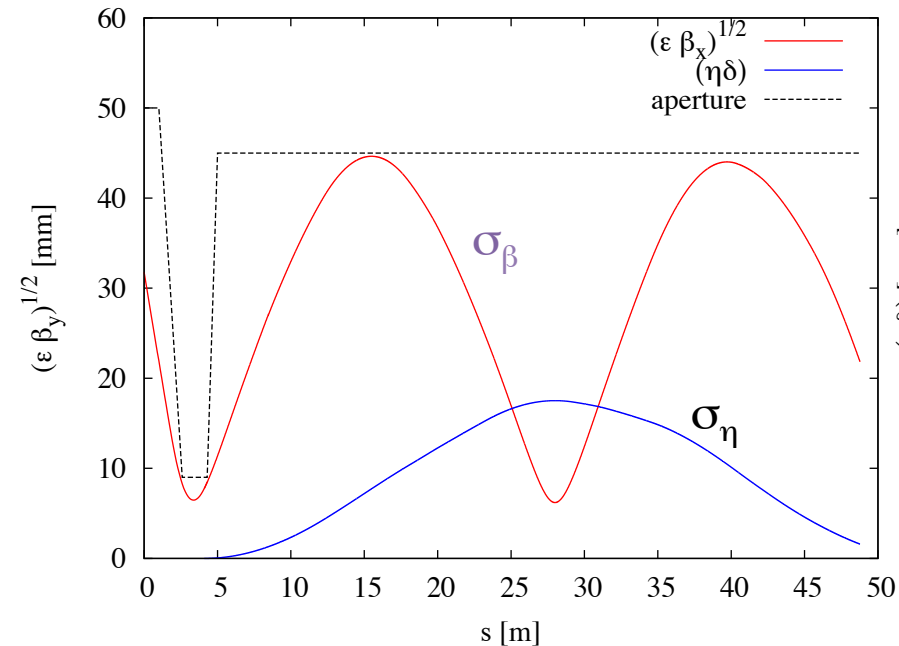
$$\Rightarrow \delta \sim 0.1\%$$

$$\begin{pmatrix} \beta_{\text{max-ring}}^x = & 40\text{m} \\ \eta_{\text{max-ring}} = & 17\text{m} \\ \epsilon_x = & 40\mu\text{m} \end{pmatrix}$$



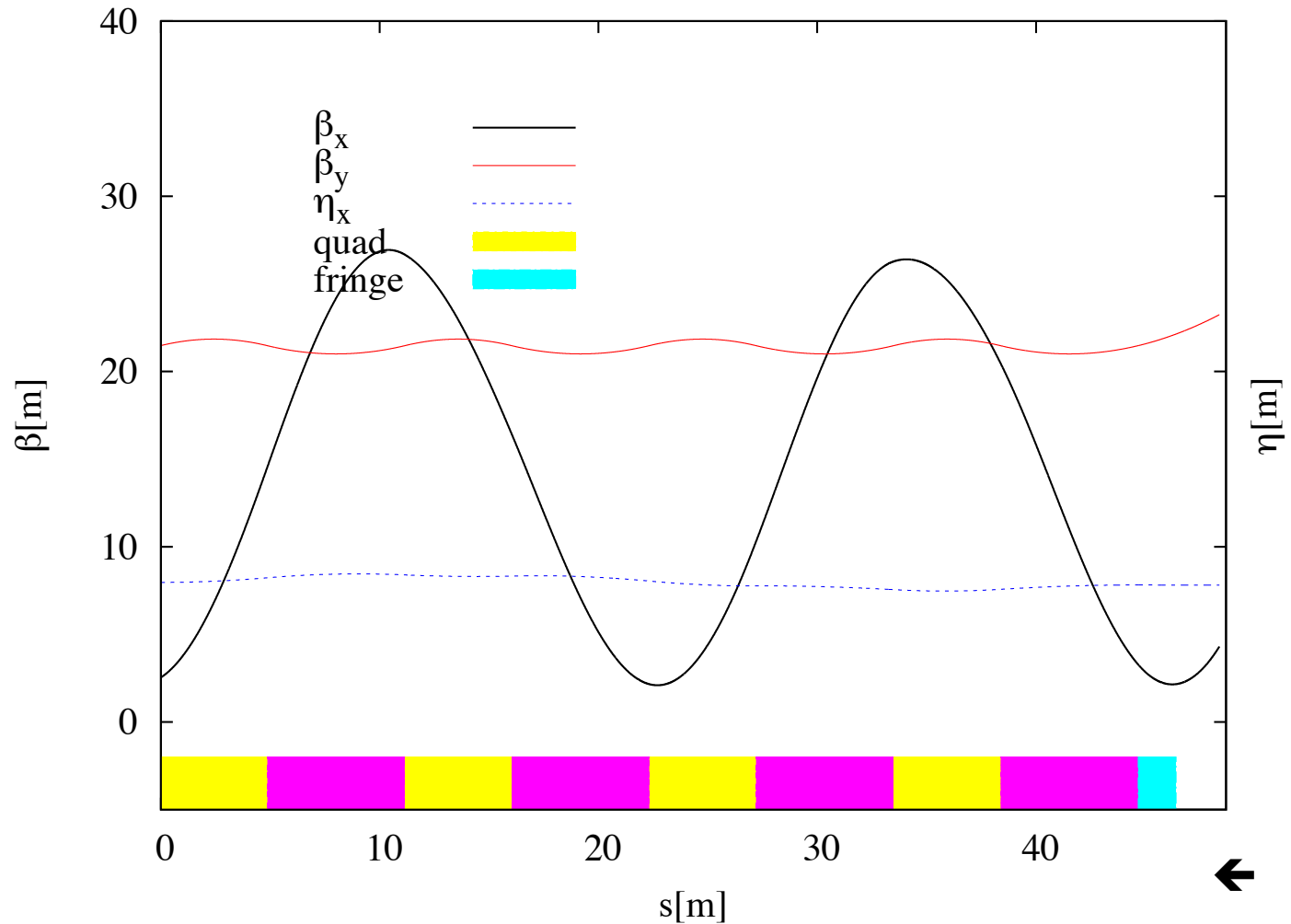
Dispersion

Contribution to beam size from betatron motion and energy spread



Horizontal aperture and beam size
Including 0.1% energy spread

Suppose we choose β_x β_y at upstream end of inflector so that 40 mm-mrad so that most of the beam fits through the inflector aperture
 (Assuming *ideal* inflector: zero field, opened ends)

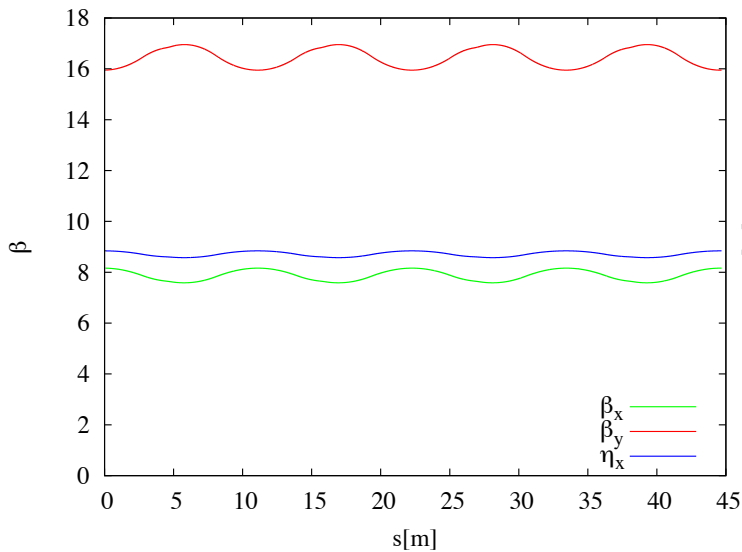


Mismatch

- Emittance of the muon beam is large and the number of captured muons is limited by apertures
- Inflector aperture \ll ring aperture
- Mismatch of twiss parameters at inflector-ring interface \Rightarrow peak β & η in

the ring are always greater than closed ring values:

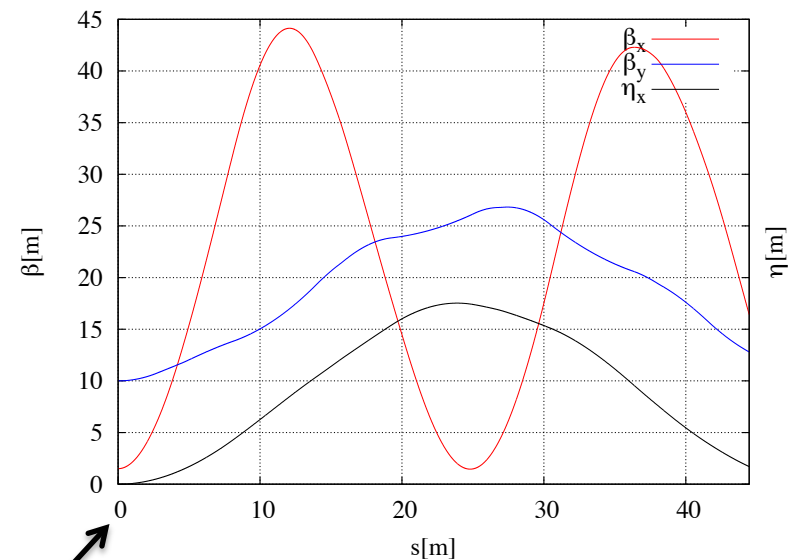
$$\beta_{\text{max-ring}} \sim \frac{\beta_{\text{ring}}^2}{\beta_{\text{inf}}}$$



Closed ring

$$\begin{aligned} \beta_{\text{inf}}^x &= 1.5\text{m} \\ \beta_{\text{inf}}^y &= 10.0\text{m} \Rightarrow \\ \eta_{\text{inf}} &= 0 \end{aligned}$$

Inflector exit



Mismatch

How do we propagate

- twiss parameters
- Beam size
- Phase space distribution

How do we determine M?

Invariant

$$a = \gamma x^2 + 2\alpha x x' + \beta x'^2$$

Define

$$\vec{\gamma} = \begin{pmatrix} \gamma \\ 2\alpha \\ \beta \end{pmatrix} \quad \vec{X} = \begin{pmatrix} x^2 \\ x x' \\ x'^2 \end{pmatrix} \quad a = \vec{\gamma} \cdot \vec{X}$$

$$x_f = M x_i$$

$$\vec{X}_f = \mathcal{M} \vec{X}_i = \begin{pmatrix} (M_{11}x_i + M_{12}x'_i)^2 \\ (M_{11}x_i + M_{12}x'_i)(M_{21}x_i + M_{22}x'_i) \\ (M_{21}x_i + M_{22}x'_i)^2 \end{pmatrix}$$

$$\mathcal{M} = \begin{pmatrix} M_{11}^2 & 2M_{11}M_{12} & M_{12}^2 \\ M_{11}M_{21} & (M_{11}M_{22} + M_{12}M_{21}) & M_{12}M_{22} \\ M_{21}^2 & 2M_{21}M_{22} & M_{22}^2 \end{pmatrix}$$

$$\begin{aligned} \vec{\gamma} \cdot \vec{X} &= \gamma^T X = \gamma^T \mathcal{M}^{-1} \mathcal{M} X \\ &= (\mathcal{M}^{-1})^T \vec{\gamma} \cdot \mathcal{M} \vec{X} \\ \rightarrow (\mathcal{M}^{-1})^T \vec{\gamma}_i &= \vec{\gamma}_f \end{aligned}$$

If we know M we can construct \mathcal{M} and \mathcal{N} to propagate twiss parameters

To construct \mathcal{N} replace the elements of M with M^{-1} in \mathcal{M} and transpose

$$\mathcal{N}^T = \begin{pmatrix} M_{22}^2 & -2M_{22}M_{12} & M_{12}^2 \\ -M_{22}M_{21} & (M_{11}M_{22} + M_{12}M_{21}) & -M_{12}M_{11} \\ M_{21}^2 & -2M_{21}M_{11} & M_{11}^2 \end{pmatrix}$$

$$\mathcal{N} = \begin{pmatrix} M_{22}^2 & -M_{22}M_{21} & M_{21}^2 \\ -2M_{22}M_{12} & (M_{11}M_{22} + M_{12}M_{21}) & -2M_{21}M_{11} \\ M_{12}^2 & -M_{12}M_{11} & M_{11}^2 \end{pmatrix}$$

In a drift $M = \begin{pmatrix} 1 & L \\ 0 & 1 \end{pmatrix}$

$$\mathcal{N} = \begin{pmatrix} 1 & 0 & 0 \\ -2L & 1 & 0 \\ L^2 & -L & 1 \end{pmatrix}$$

If at $L=0$, β is at a minimum then

$$\vec{\gamma}(L) = \mathcal{N} \begin{pmatrix} 1/\beta_0 \\ 0 \\ \beta_0 \end{pmatrix} = \begin{pmatrix} 1/\beta_0 \\ -\frac{2L}{\beta_0} \\ \frac{L^2}{\beta_0} + \beta_0 \end{pmatrix}$$

$$\beta(L) = \beta_0 + \frac{L^2}{\beta_0}$$

β increases rapidly away from a minimum

Now we know how to propagate twiss parameters along a beam line

Extract twiss parameters from particle distribution ?

$$x = \sqrt{2\beta\epsilon_x} \cos \phi_x + \eta_x \delta$$

$$\beta_x \epsilon_x = \langle xx \rangle - \frac{\langle x\delta \rangle^2}{\langle \delta\delta \rangle}$$

Notation

$$\beta_x \epsilon_x = \sigma(1, 1) - \frac{\sigma^2(1, 6)}{\sigma(6, 6)}$$

Betatron dist

$$\sigma_\beta(i, j) \equiv \sigma(i, j) - \frac{\sigma(i, 6) \sigma(j, 6)}{\sigma(6, 6)}$$

$$\epsilon_x^2 = \sigma_\beta(1, 1) \sigma_\beta(2, 2) - \sigma_\beta^2(1, 2)$$

We can compute $\sigma(i, j)$ for the distribution and extract twiss

$$\epsilon_x \begin{pmatrix} \beta_x & -\alpha_x \\ -\alpha_x & \gamma_x \end{pmatrix} = \begin{pmatrix} \sigma_\beta(1, 1) & \sigma_\beta(1, 2) \\ \sigma_\beta(1, 2) & \sigma_\beta(2, 2) \end{pmatrix}$$

If we know M we can propagate twiss

And if we have a distribution we can compute twiss

What if we know twiss, how do we determine M?

M transports particles from point 1 with $\vec{\gamma}_1$ to point 2 with $\vec{\gamma}_2$

Write M in terms of $\vec{\gamma}_1, \vec{\gamma}_2$

$$x_1 = a\sqrt{\beta_1} \cos \phi_1, \quad x'_1 = -a \left(\frac{\alpha_1}{\sqrt{\beta_1}} \cos \phi_1 + \frac{1}{\sqrt{\beta_1}} \sin \phi_1 \right)$$

$$\begin{pmatrix} a\sqrt{\beta_2} \cos \phi_2 \\ -a \left(\frac{\alpha_2}{\sqrt{\beta_2}} \cos \phi_2 + \frac{1}{\sqrt{\beta_2}} \sin \phi_2 \right) \end{pmatrix} = M \begin{pmatrix} a\sqrt{\beta_1} \cos \phi_1 \\ -a \left(\frac{\alpha_1}{\sqrt{\beta_1}} \cos \phi_1 + \frac{1}{\sqrt{\beta_1}} \sin \phi_1 \right) \end{pmatrix}$$

Define
$$G = \begin{pmatrix} \sqrt{\beta_i} & 0 \\ -\frac{\alpha_i}{\sqrt{\beta_i}} & -\frac{1}{\sqrt{\beta_i}} \end{pmatrix}$$

So that
$$\begin{pmatrix} x \\ x' \end{pmatrix} = G \begin{pmatrix} \cos \phi \\ \sin \phi \end{pmatrix}$$

$$\begin{pmatrix} x \\ x' \end{pmatrix} = G \begin{pmatrix} \cos \phi \\ \sin \phi \end{pmatrix}$$

$$G_2 a \begin{pmatrix} \cos \phi_2 \\ \sin \phi_2 \end{pmatrix} = M G_1 a \begin{pmatrix} \cos \phi_1 \\ \sin \phi_1 \end{pmatrix}$$

$$G_2^{-1} M G_1 = \begin{pmatrix} \cos \Delta\phi & -\sin \Delta\phi \\ \sin \Delta\phi & \cos \Delta\phi \end{pmatrix}$$

$$M = G_2 \begin{pmatrix} \cos \Delta\phi & -\sin \Delta\phi \\ \sin \Delta\phi & \cos \Delta\phi \end{pmatrix} G_1^{-1}$$

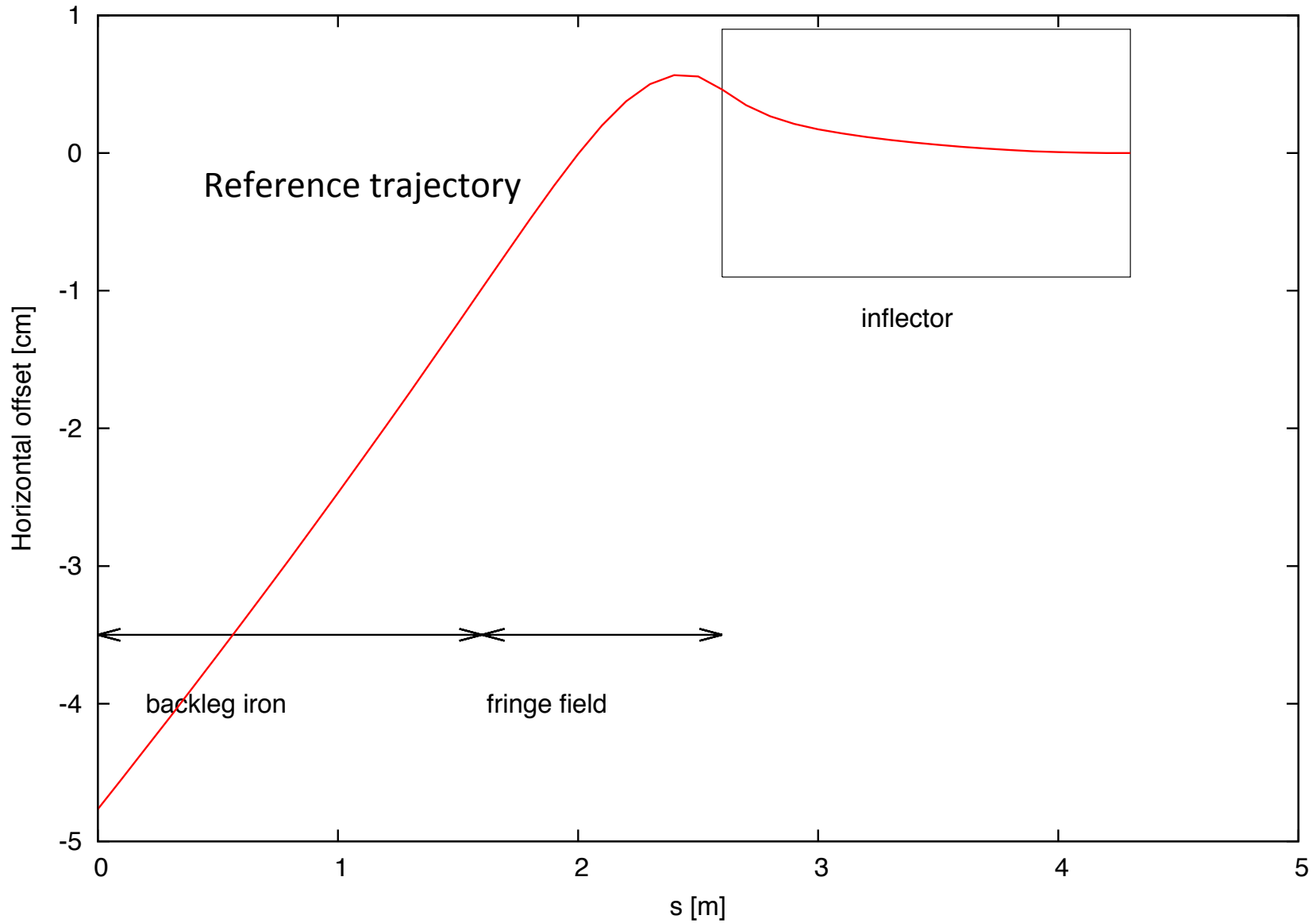
$$M = \begin{pmatrix} \sqrt{\frac{\beta_2}{\beta_1}} [\cos \Delta\phi + \alpha_1 \sin \Delta\phi] & \sqrt{\beta_2 \beta_1} \sin \Delta\phi \\ -\frac{1}{\sqrt{\beta_1 \beta_2}} [(\alpha_1 - \alpha_2) \cos \Delta\phi - (1 + \alpha_2 \alpha_1) \sin \Delta\phi] & \sqrt{\frac{\beta_1}{\beta_2}} [\cos \Delta\phi - \alpha_2 \sin \Delta\phi] \end{pmatrix}$$

If we know M we can propagate twiss
And if we have a distribution we can compute twiss
From twiss we can determine M (to a phase advance)

What if all we have is a field map?
How do we compute M

M is the Jacobian of the mapping from 1 to 2
We need to compute derivatives by tracking
Derivatives are a change with respect to a reference

$$M = \begin{pmatrix} \frac{\partial x_1^f}{\partial x_1^i} & \frac{\partial x_1^f}{\partial x_2^i} \\ \frac{\partial x_2^f}{\partial x_1^i} & \frac{\partial x_2^f}{\partial x_2^i} \end{pmatrix}$$



To compute the transfer matrix with respect to a reference trajectory track from 2 sets of initial coordinates near the reference

$$\vec{x}_i \rightarrow \vec{x}_f \quad \text{Offset from reference by } \vec{\Delta}^j$$

$$M \vec{x}_i^1 = \vec{x}_f^1$$

Or for two non degenerate trajectories

$$M \vec{x}_i^1 \vec{x}_i^2 = \vec{x}_f^1 \vec{x}_f^2$$

$$M X_1 = X_2$$

$$M = X_2 X_1^{-1}$$

Closed orbit reference

$$M(\vec{x}_i - \vec{x}_c) = \vec{x}_f - \vec{x}_c$$

$$M\vec{x}_i - (M\vec{x}_c - \vec{x}_c) = \vec{x}_f$$

$$M\vec{x}_i - \vec{v} = \vec{x}_f$$

Define

$$P = \begin{pmatrix} M & \begin{pmatrix} -v \\ -v' \end{pmatrix} \\ 0 & 1 \end{pmatrix} \quad PX_i = X_f \quad X_i = \begin{pmatrix} x_i \\ x'_i \\ 1 \end{pmatrix}$$

$$Y_i = X_i^1 X_i^2 X_i^3 \quad PY_i = Y_f \quad P = Y_f Y_i^{-1}$$

$$\rightarrow M, \vec{v} \quad \vec{x}_c = (M - I)^{-1} \vec{v}$$

Conclusion: track 3 non degenerate orbits through 1 full turn to find closed orbit and M
Track 7 non degenerate orbits to find M for full 6-d phase space

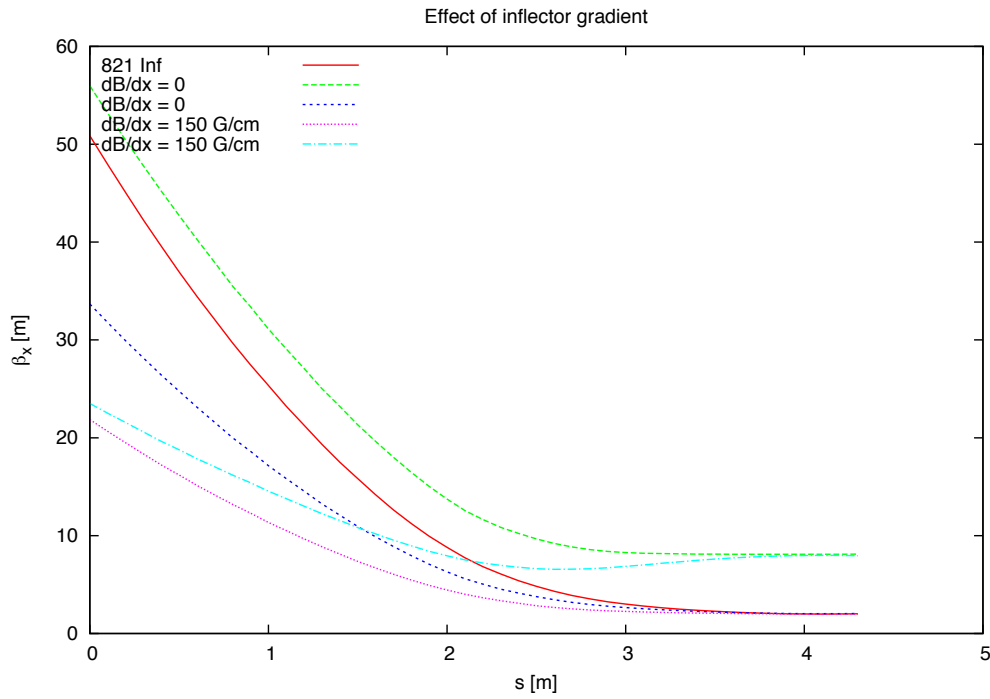
Suppose we want to study the dependence of capture on beta in the inflector

- Beta in the inflector is determined by quads in M5 line.
- Changing Beta in the inflector => change beta in M5

The distribution to be tracked is defined at the end of M5 with unknown twiss

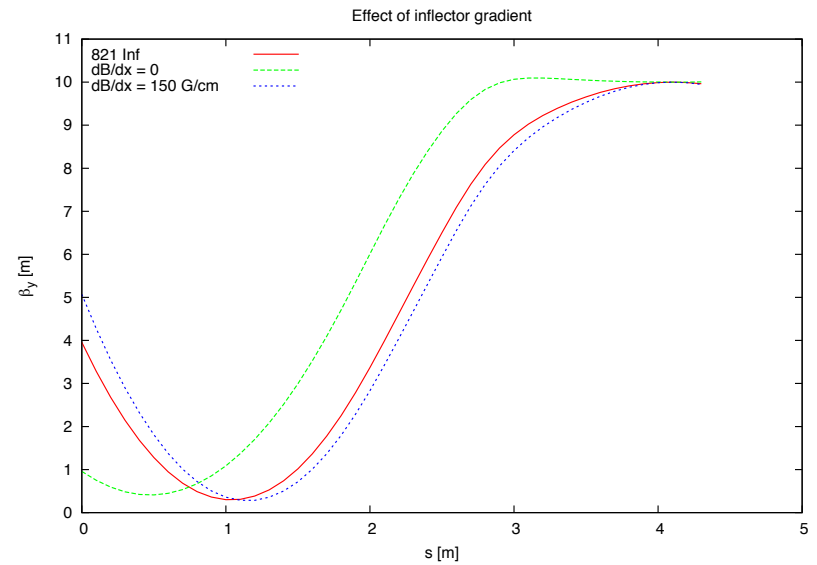
1. Construct M through the injection channel by tracking through the field maps
2. Use M for injection channel to propagate twiss in inflector $\vec{\gamma}_I$ back to end of M5 $\vec{\gamma}_{M5}$
3. Compute twiss $\vec{\gamma}_D$ for Diktys's distribution, (which is defined at end of M5)
4. Construct $M(\vec{\gamma}_D, \vec{\gamma}_{M5})$ and use it to transform the distribution to the target twiss
5. The transformed distribution will be characterized by $\vec{\gamma}_{M5}$ at the start of tracking and by $\vec{\gamma}_I$ when tracked (using M) to the inflector

$dB_y/dx = 150\text{G/cm}$



Horizontal beta

Vertical beta

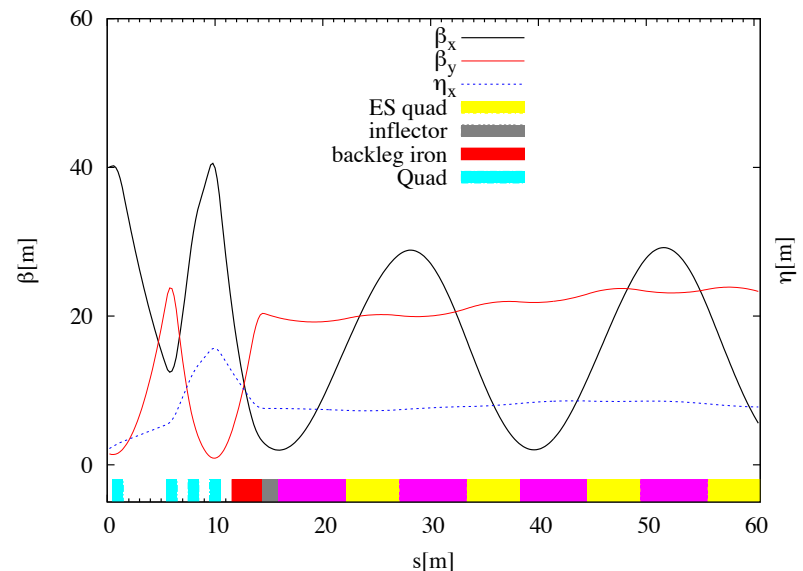


Example

Phase space matching

- 40 mm-mrad clears inflector if at inflector midpoint $\beta_x=2\text{m}$, $\beta_y=19\text{m}$
- Transfer map for propagation of twiss parameters determined by tracking through Wuzeng field map

$$T = \begin{pmatrix} 1.61951 & 5.68287 & -0.10265 & -0.12099 & 0.00032 & 0.02890 \\ 0.29750 & 1.65985 & -0.03583 & -0.05500 & 0.00025 & 0.00047 \\ -0.08510 & -0.05582 & 0.44649 & 3.19115 & -0.00014 & 0.00884 \\ -0.02370 & -0.04136 & -0.25378 & 0.40416 & -0.00010 & -0.00031 \\ 0.00776 & 0.04529 & -0.00296 & 0.00323 & 1.00000 & 0.00046 \\ 0.00000 & 0.00000 & 0.00000 & 0.00000 & 0.00000 & 1.00000 \end{pmatrix}$$



Large aperture inflector - Kickers

Kicker fields scale with displacement d

With larger aperture inflector

$$d = 77\text{mm} \Rightarrow 91\text{mm}$$

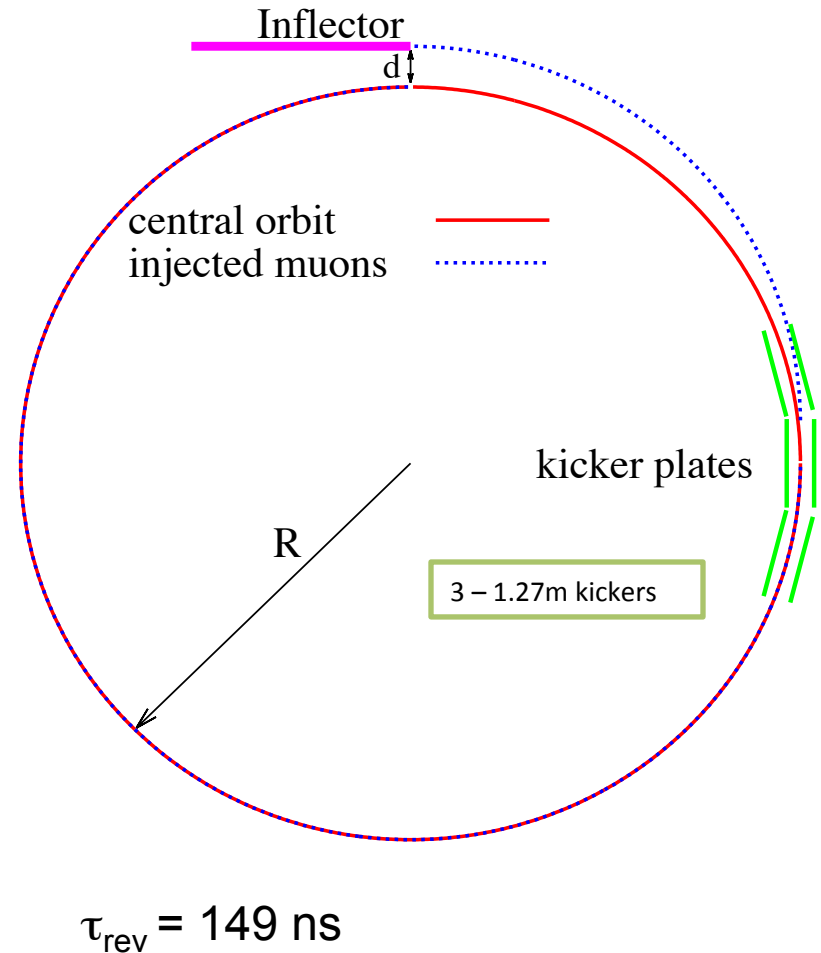
Kick angle ~ 10.8 to 12.8 mrad

Integrated B-field ~ 1.11 to 1.31 kG-m

$$\text{Kick Angle} = d/R$$

$$R = R_0(1+\delta)$$

$$\theta = \theta_0(1-\delta)$$



Orientation

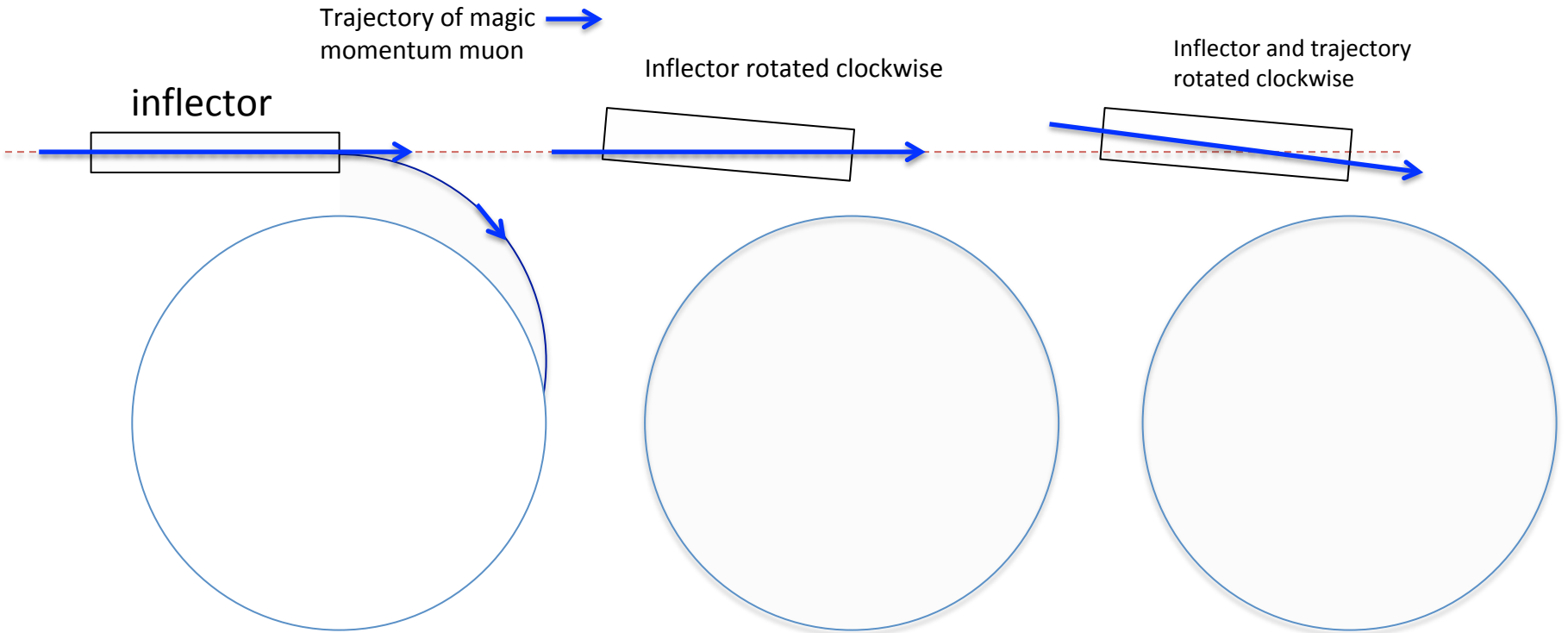
Injection channel

Trajectory of magic
momentum muon →

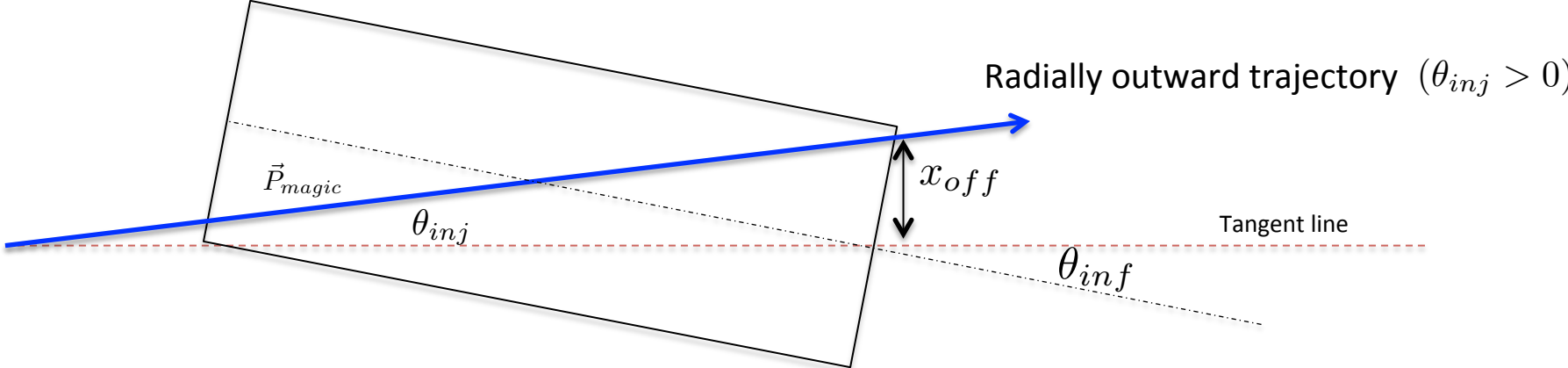
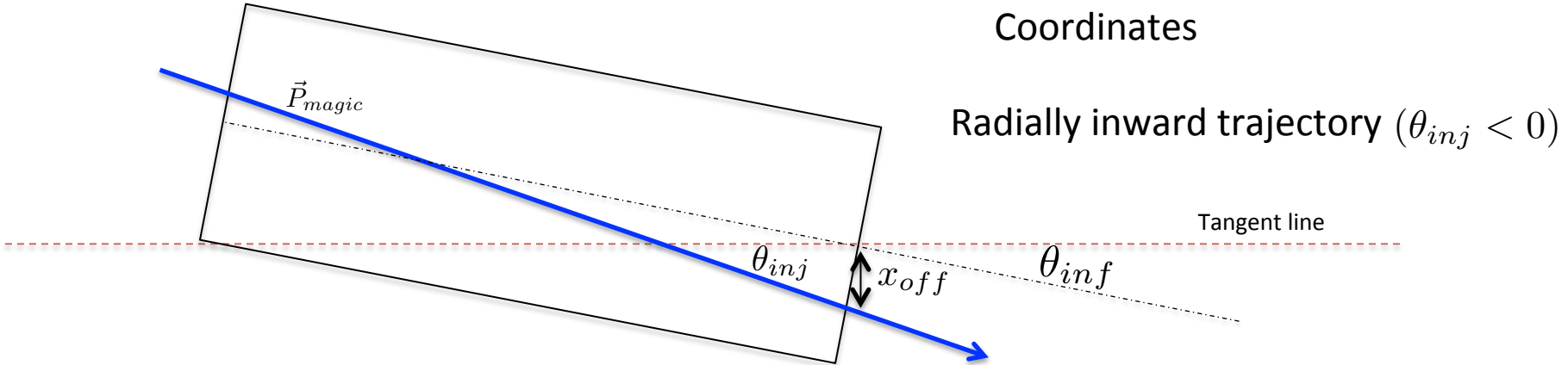
Inflector rotated clockwise

Inflector and trajectory
rotated clockwise

inflector



Coordinates



Angles and offset of magic momentum trajectory

Trajectory of injected particle

$$\begin{aligned}x_{inf} &= x_{\beta 0} + x_{\delta 0} \\ &= x_{\beta 0} + \eta \delta\end{aligned}$$

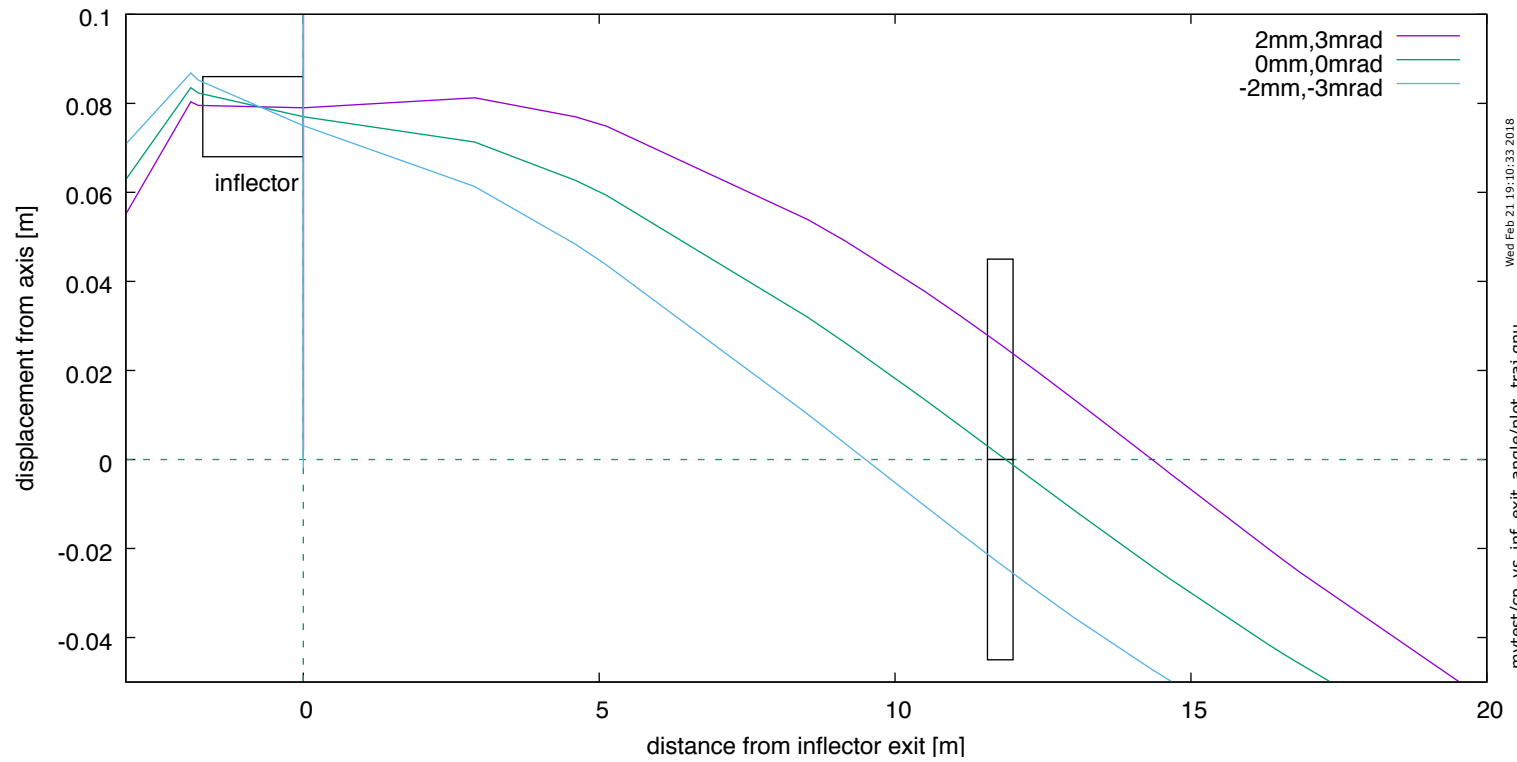
$$x_{\beta 0} = x_{inf} - \eta(0)\delta$$

Trajectory of injected particle

$$\begin{aligned}x_{inf} &= x_{\beta 0} + x_{\delta 0} \\ &= x_{\beta 0} + \eta \delta\end{aligned}$$

$$x_{\beta 0} = x_{inf} - \eta(0)\delta$$

$$x(s) = x_{\beta 0} \sqrt{\frac{\beta(s)}{\beta(0)}} \cos(\phi(s)) + \eta(s)\delta$$



Wed Feb 21 19:10:33 2018

mytest/cp_vs_inf_exit_angle/plot_traj.gnu

Trajectory of injected particle

$$\begin{aligned}x_{inf} &= x_{\beta 0} + x_{\delta 0} \\ &= x_{\beta 0} + \eta \delta\end{aligned}$$

$$x_{\beta 0} = x_{inf} - \eta(0)\delta$$

$$x(s) = x_{\beta 0} \sqrt{\frac{\beta(s)}{\beta(0)}} \cos(\phi(s)) + \eta(s)\delta$$

$$\Delta x(s) = k \sqrt{\beta_{kick} \beta(s)} \sin(\phi(s) - \pi/2) \quad \text{Change due to kicker}$$

Trajectory of injected particle

$$\begin{aligned}x_{inf} &= x_{\beta 0} + x_{\delta 0} \\ &= x_{\beta 0} + \eta \delta\end{aligned}$$

$$x_{\beta 0} = x_{inf} - \eta(0)\delta$$

$$x(s) = x_{\beta 0} \sqrt{\frac{\beta(s)}{\beta(0)}} \cos(\phi(s)) + \eta(s)\delta$$

$$\Delta x(s) = k \sqrt{\beta_{kick} \beta(s)} \sin(\phi(s) - \pi/2) \quad \text{Change due to kicker}$$

$$x(s) = (x_{inf} - \delta\eta_0 - k\beta_0) \cos(\phi(s)) + \eta_0\delta \quad \text{All together}$$

Trajectory of injected particle

$$\begin{aligned}x_{inf} &= x_{\beta 0} + x_{\delta 0} \\ &= x_{\beta 0} + \eta\delta\end{aligned}$$

$$x_{\beta 0} = x_{inf} - \eta(0)\delta$$

$$x(s) = x_{\beta 0} \sqrt{\frac{\beta(s)}{\beta(0)}} \cos(\phi(s)) + \eta(s)\delta$$

$$\Delta x(s) = k \sqrt{\beta_{kick}\beta(s)} \sin(\phi(s) - \pi/2) \quad \text{Change due to kicker}$$

$$x(s) = (x_{inf} - \delta\eta_0 - k\beta_0) \cos(\phi(s)) + \eta_0\delta \quad \text{All together}$$

$$x_{min} = x_{inf} - k\beta_0$$

$$x_{max} = -x_{inf} + k\beta_0 + 2\eta_0\delta$$

Perfect kick

$$k\beta_0 = x_{inf}$$

If the particle exits the inflector with some angle x'_{inj} (with respect to the tangent line)

Then before it reaches the kicker

$$x(s) = x_{\beta 0} \cos \phi(s) + x'_{xinf} \beta_0 \sin \phi(s) + \eta_0 \delta$$

If the particle exits the inflector with some angle x'_{inj} (with respect to the tangent line)

Then before it reaches the kicker

$$x(s) = x_{\beta 0} \cos \phi(s) + x'_{xinf} \beta_0 \sin \phi(s) + \eta_0 \delta$$

And beyond the kicker

$$\begin{aligned} x(s) &= (x_{inf} - \delta \eta_0) \cos(\phi(s)) + \eta_0 \delta - k \beta_0 \cos(\phi(s)) + x'_{inf} \beta_0 \sin \phi(s) \\ &= A \cos(\phi(s) + \phi_0) + \eta_0 \delta \end{aligned}$$

where

$$A = \pm \sqrt{(x_{inf} - \delta \eta_0 - k \beta_0)^2 + (x'_{inf} \beta_0)^2} \quad (\text{CBO amplitude})$$

and

$$\tan \phi_0 = \frac{x'_{inf} \beta_0}{x_{inf} - \delta \eta_0 - k \beta_0} \quad (\text{CBO phase})$$

The extremes of displacement are

$$x_{ext} = \pm |A| + \eta_0 \delta \quad (\text{CBO envelope})$$

$$A = \pm \sqrt{(x_{inf} - \delta\eta_0 - k\beta_0)^2 + (x'_{inf}\beta_0)^2}$$
$$\rightarrow |A| \geq x'_{inf}\beta_0 \sim (3 \text{ mrad})(7 \text{ m}) \sim 21 \text{ mm}$$

The kicker can not compensate a nonzero injection angle

$$A = \pm \sqrt{(x_{inj} - \delta\eta_0 - k\beta_0)^2 + (x'_{inj}\beta_0)^2}$$

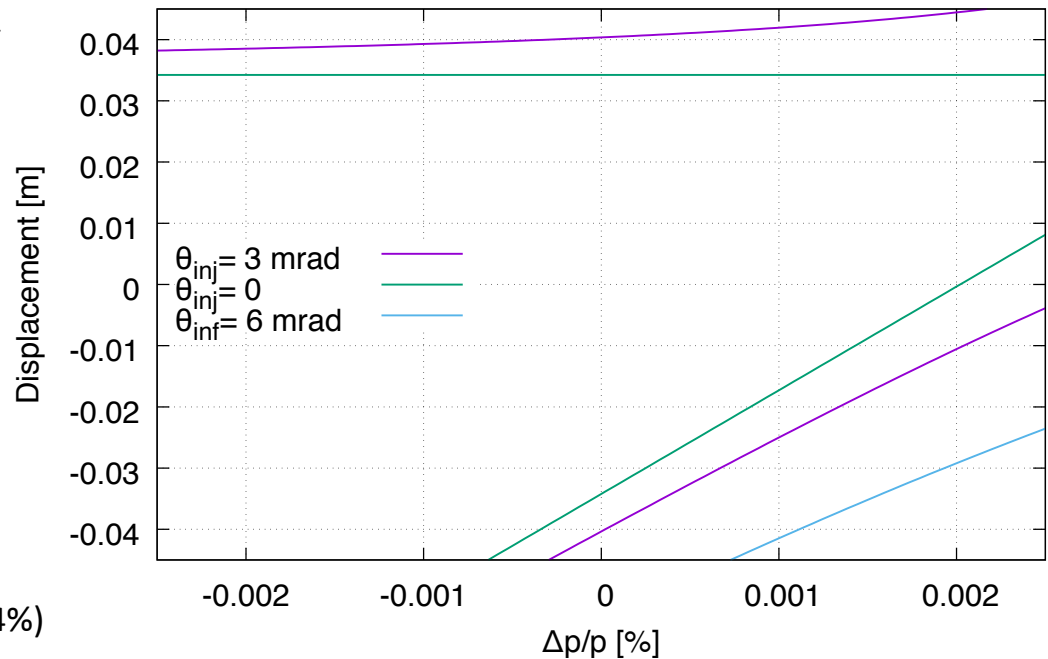
$$\rightarrow |A| \geq x'_{inj}\beta_0 \sim (3 \text{ mrad})(7 \text{ m}) \sim 21 \text{ mm}$$

The kicker can not compensate a nonzero injection angle

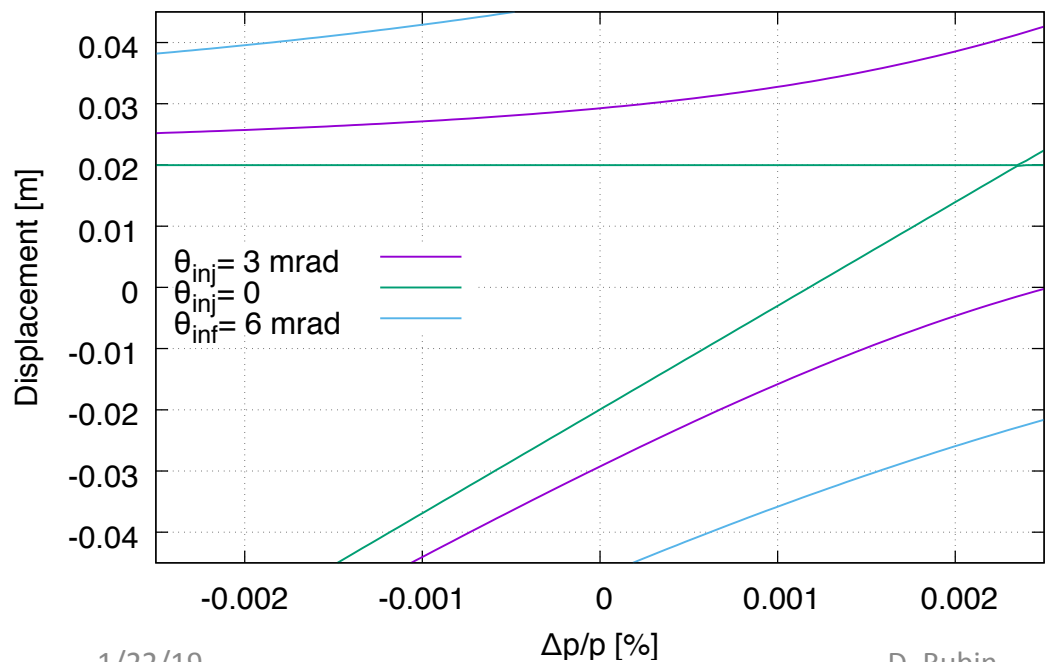
Set $\left\{ \begin{array}{l} x_{inj} = 77 \text{ mm} \\ \eta_0 \sim 8 \text{ m} \\ \beta_0 \sim 7 \text{ m} \end{array} \right.$ And vary kick k , momentum δ , θ_{inj}

$$A = \pm \sqrt{(x_{inf} - \delta\eta_0 - k\beta_0)^2 + (x'_{inf}\beta_0)^2}$$

kick angle = 6 mrad (56%)



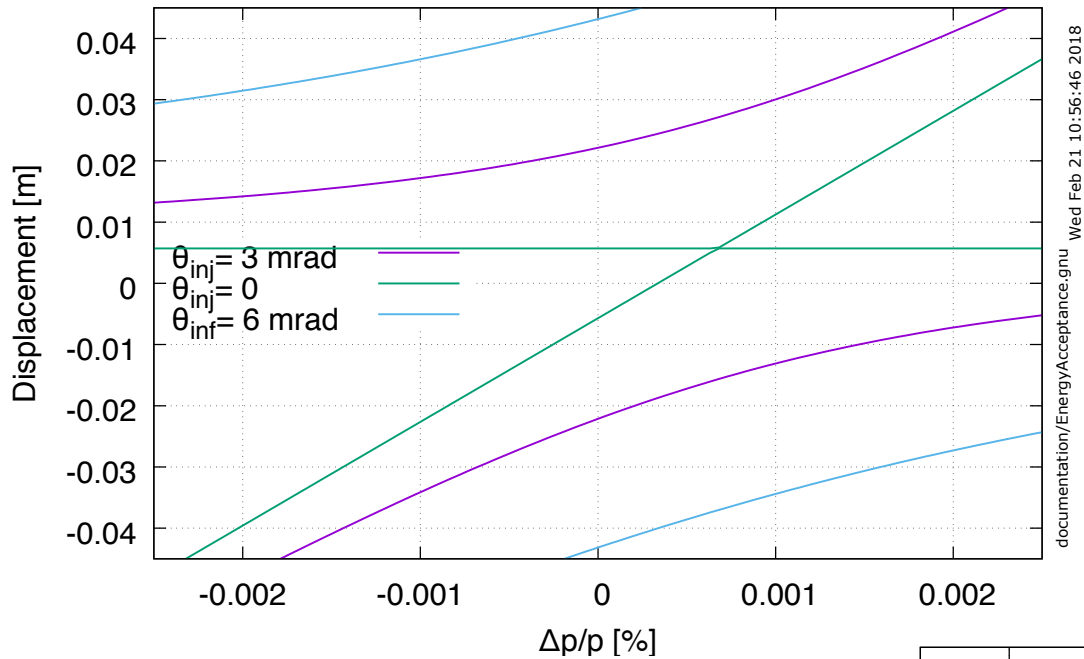
kick angle = 8 mrad (74%)



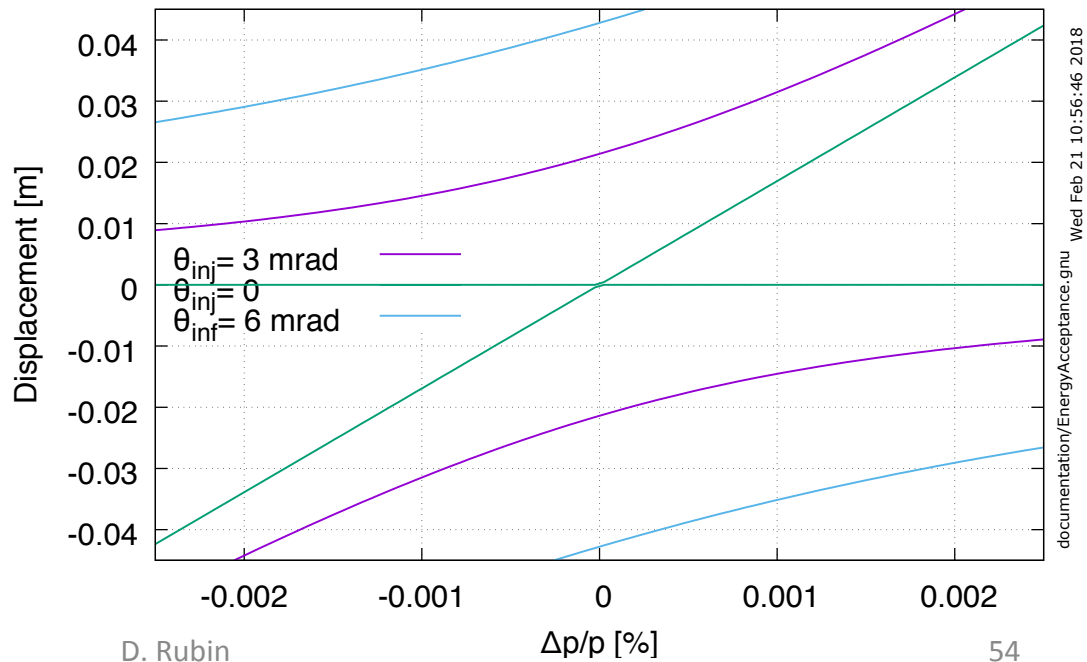
documentation/EnergyAcceptance.gnu Wed Feb 21 10:56:46 2018

documentation/EnergyAcceptance.gnu Wed Feb 21 10:56:46 2018

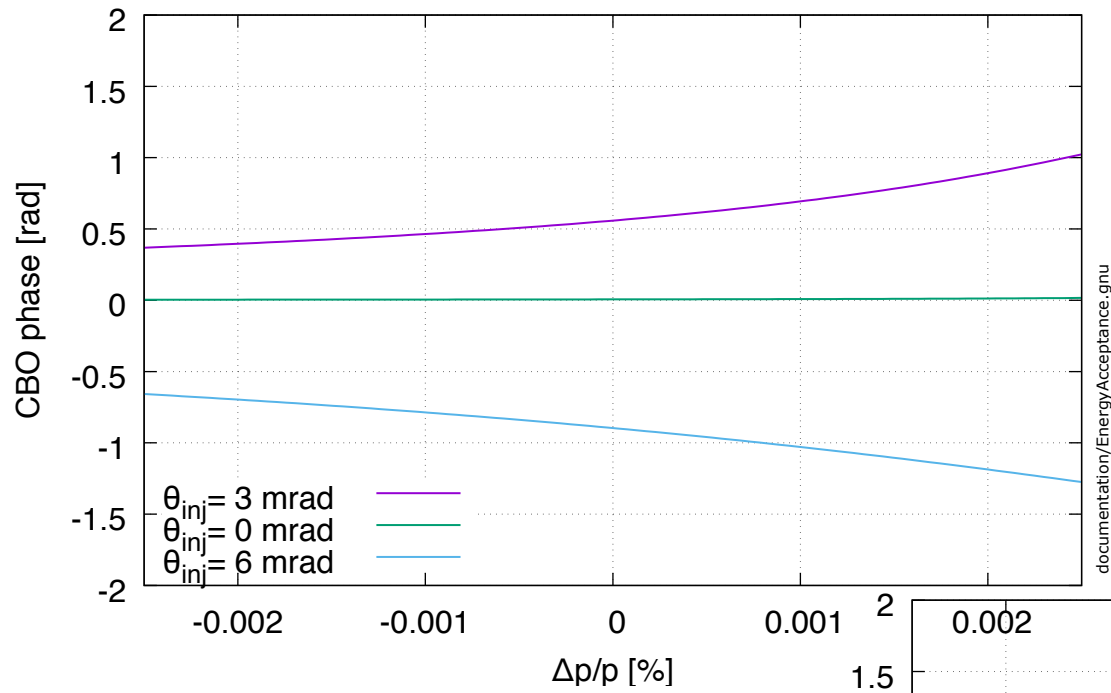
kick angle = 10 mrad (93%)



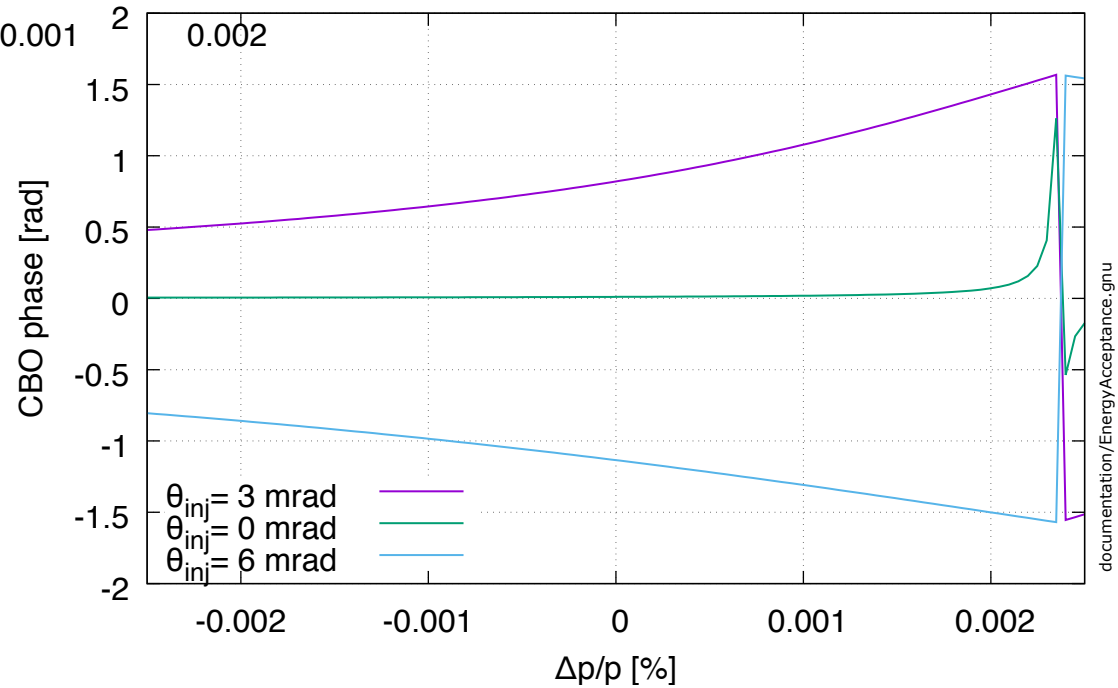
kick angle = 10.8 mrad (100%)



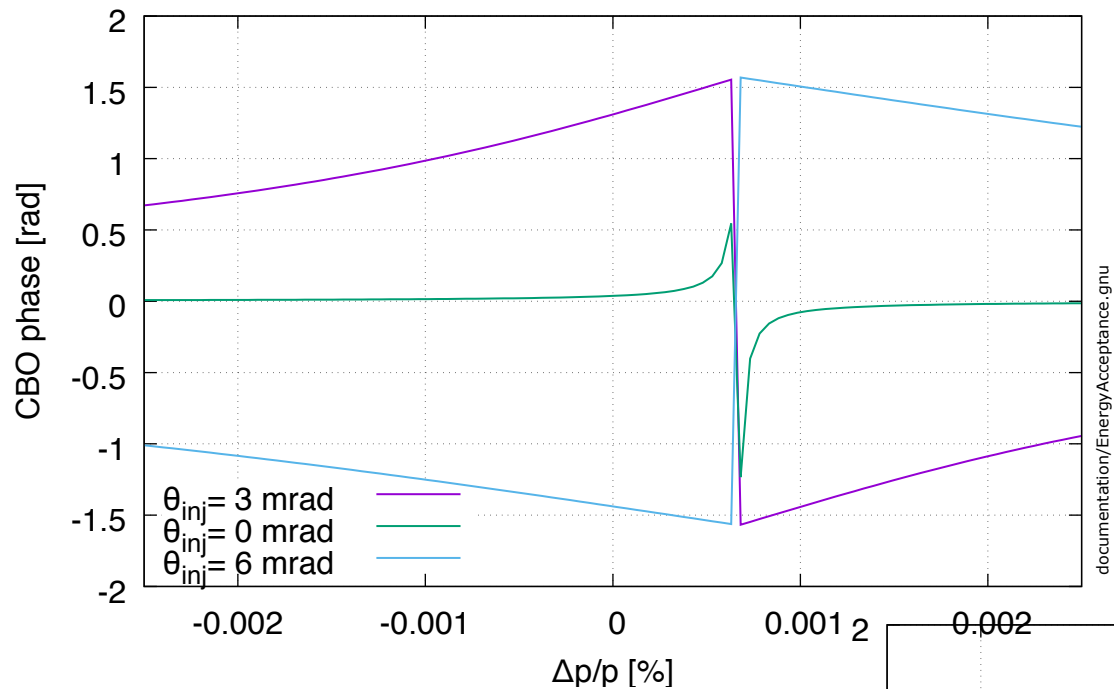
kick angle = 6 mrad



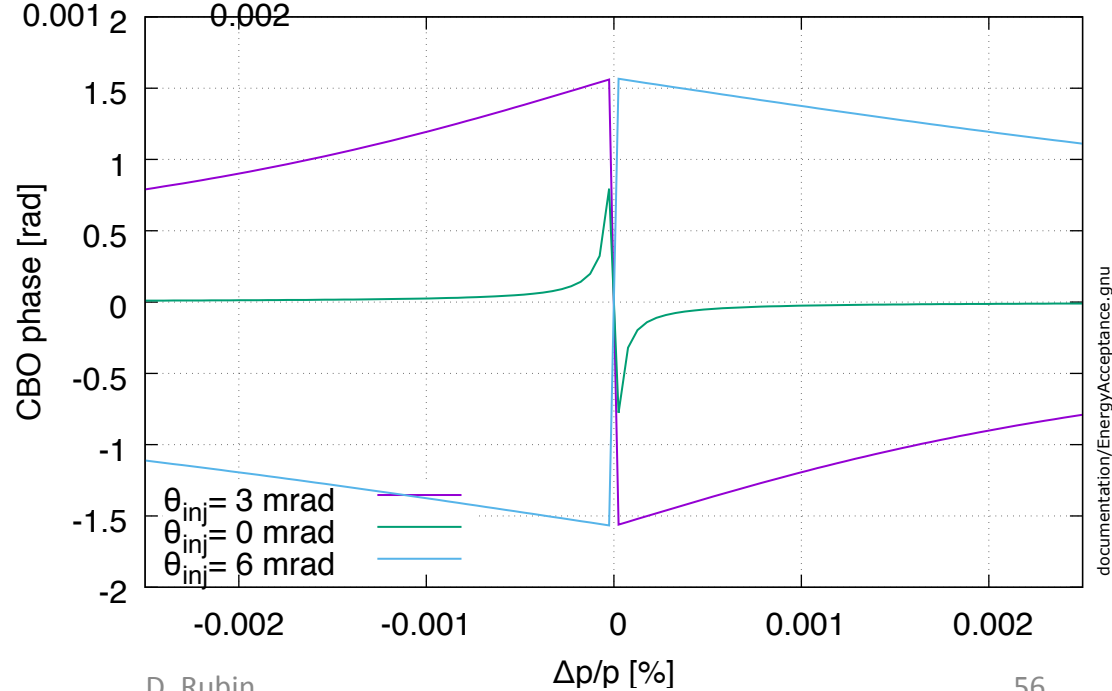
kick angle = 8 mrad



kick angle = 10 mrad



kick angle = 10.8 mrad



Suppose there is finite dispersion at the inflector exit

$$A = \pm \sqrt{(x_{inf} - \delta(\eta_0 - \eta_{inf}) - k\beta_0)^2 + ((x'_{inf} + \eta'_{inf}\delta)\beta_0)^2}$$

and

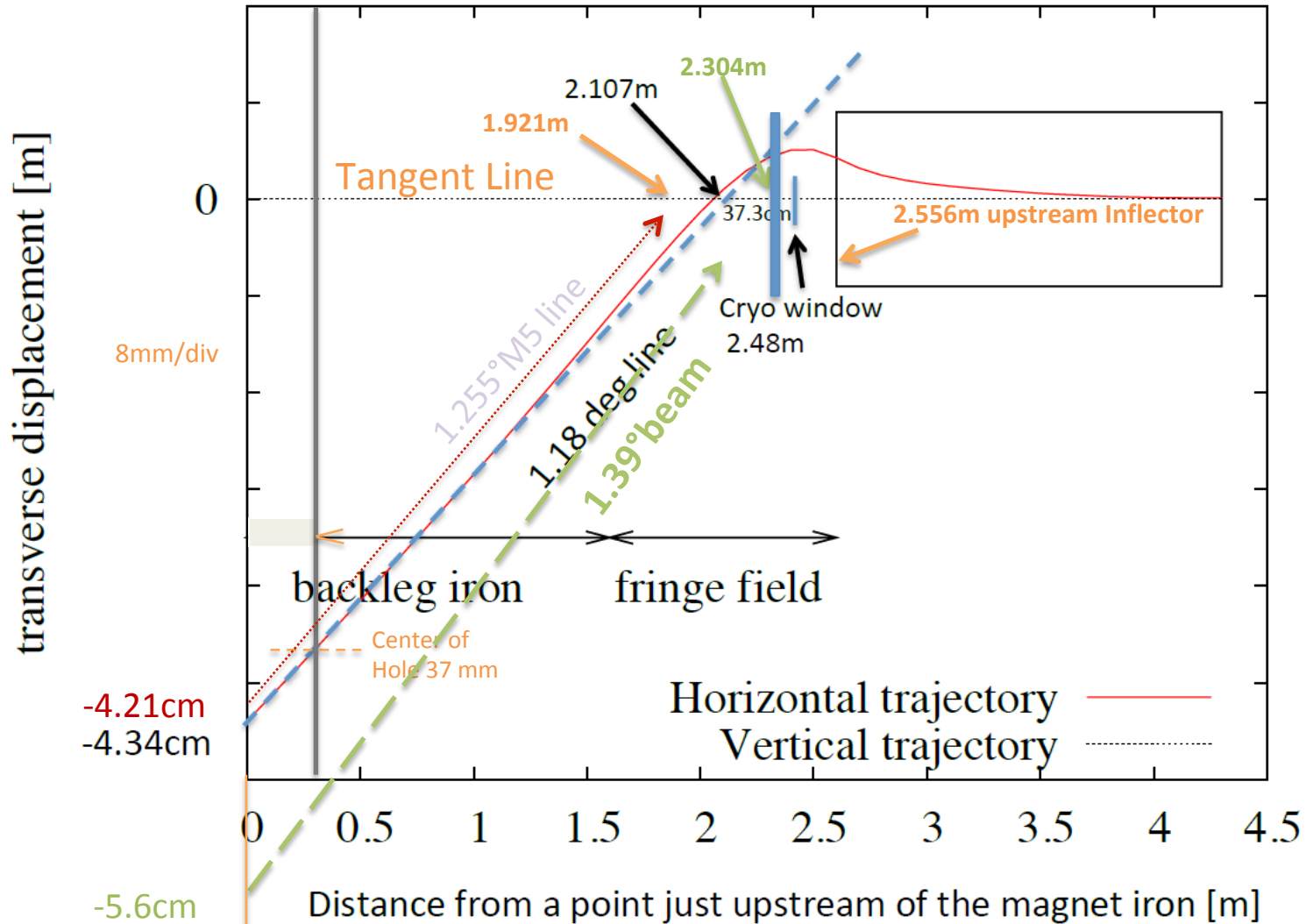
$$\tan \phi_0 = \frac{x_{inf} - \delta(\eta_0 - \eta_{inf}) - k\beta_0}{(x'_{inf} + \eta'_{inf}\delta)\beta_0}$$

The extremes of displacement become

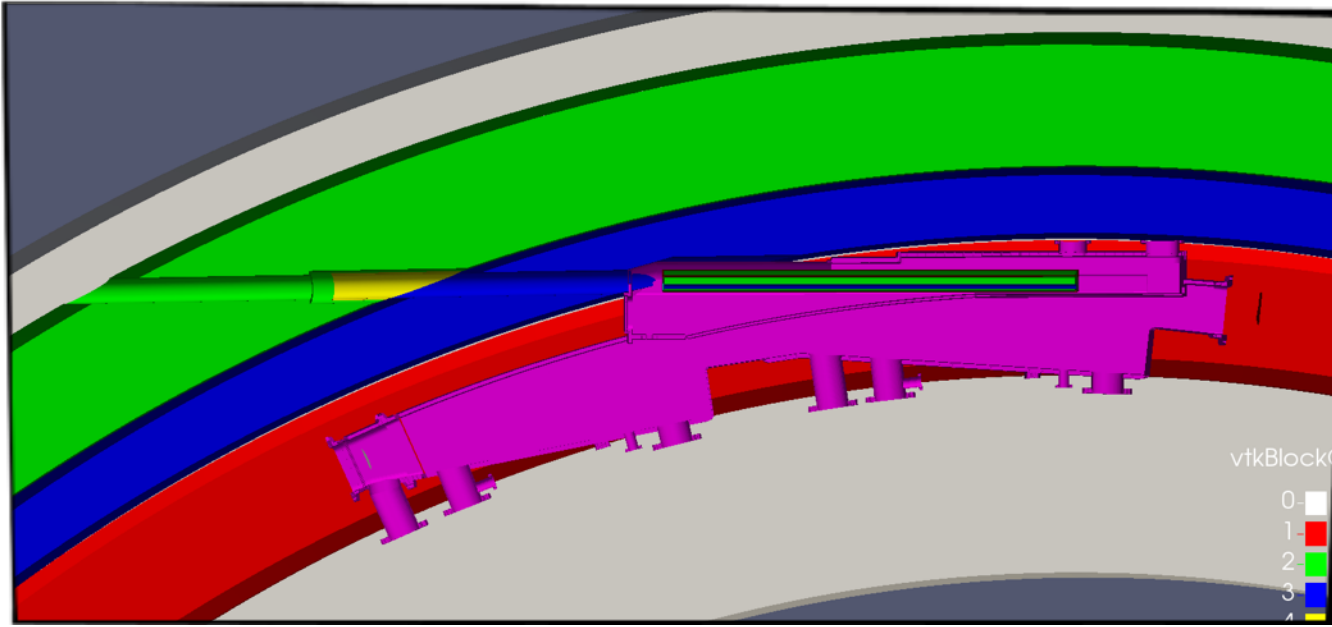
$$x_{ext} = \pm |A| + (\eta_0 - \eta_{inf})\delta$$

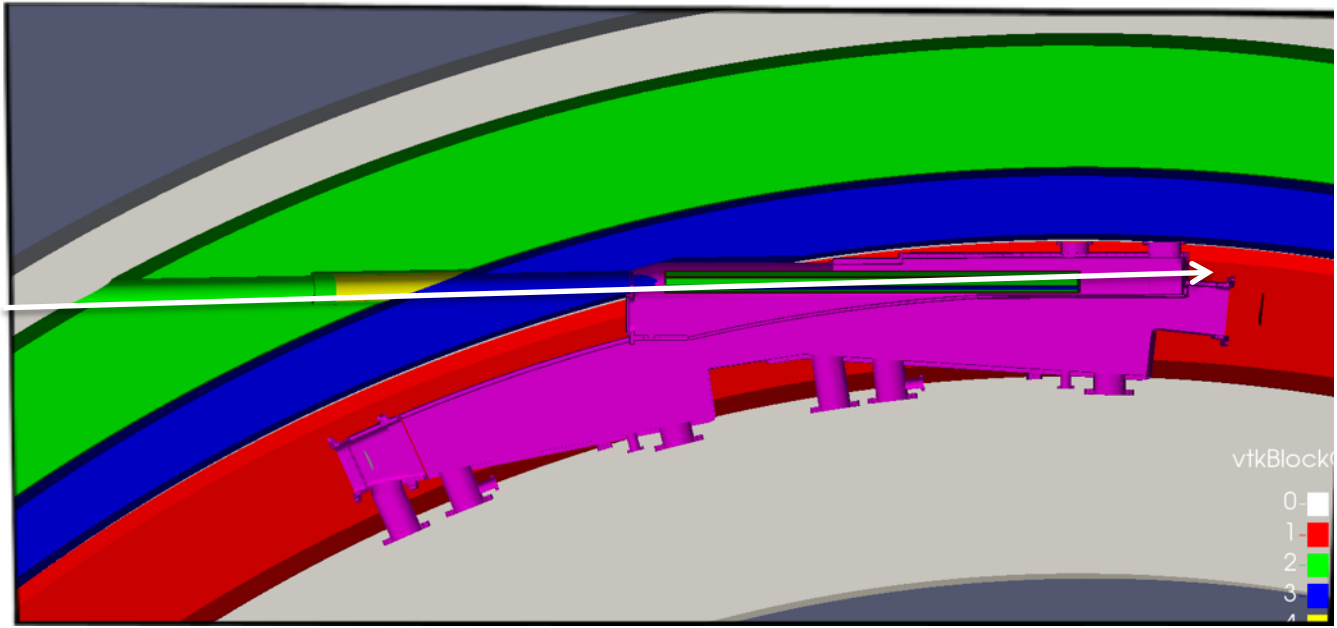
Some Facts

- Storage fraction is optimal if beam enters inflector with outward going angle (+3.5 mrad)
- Inflector is rotated about downstream end by -2.3 mrad (pointing radially inward)
- Storage optimized with $B_{\text{inf}} \sim 1.5\%$ high
- Storage fraction increases rapidly with kicker voltage (with no plateau in sight)
- CBO amplitude ($\sim 28\text{mm}$ peak-peak), depends weakly on kicker voltage and inflector field
- Closed orbit displaced 6-10 mm outward.
 - Displacement depends weakly on kicker voltage and inflector field
- De-coherence time is long; $\sim 300\text{ us}$ => narrow momentum acceptance
- De-coherence time increases with kicker voltage
 - More muons are stored, but in a narrower momentum band?

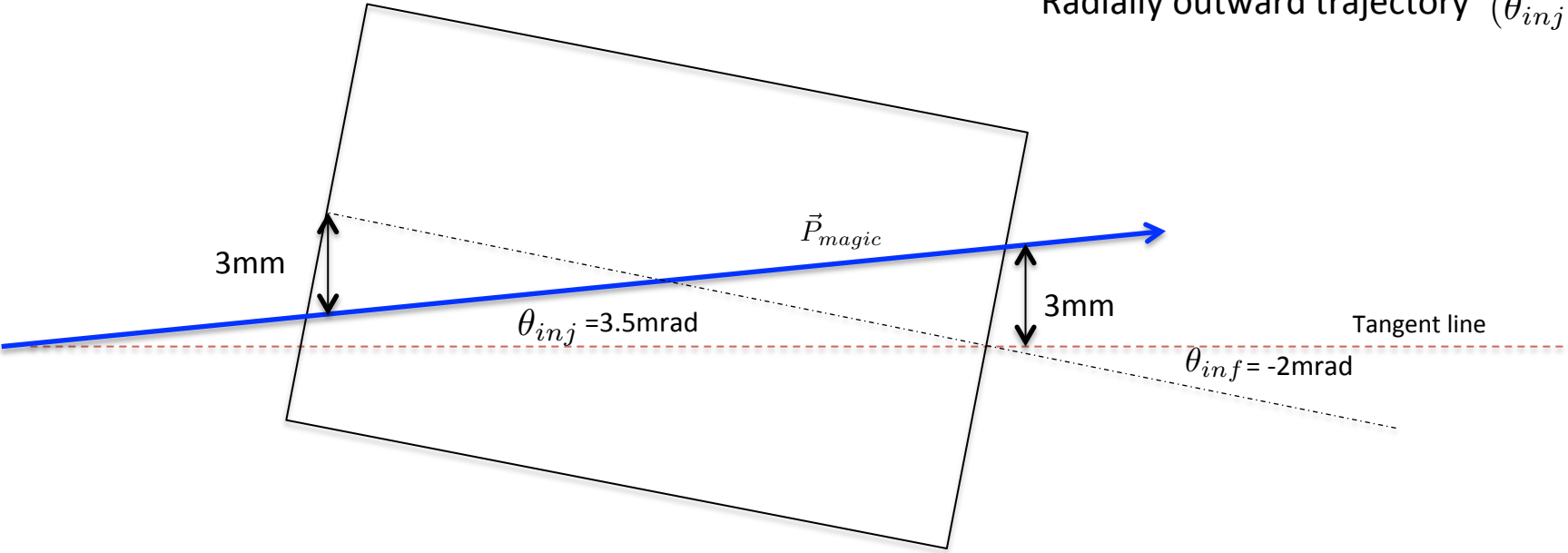


Jim Morgan

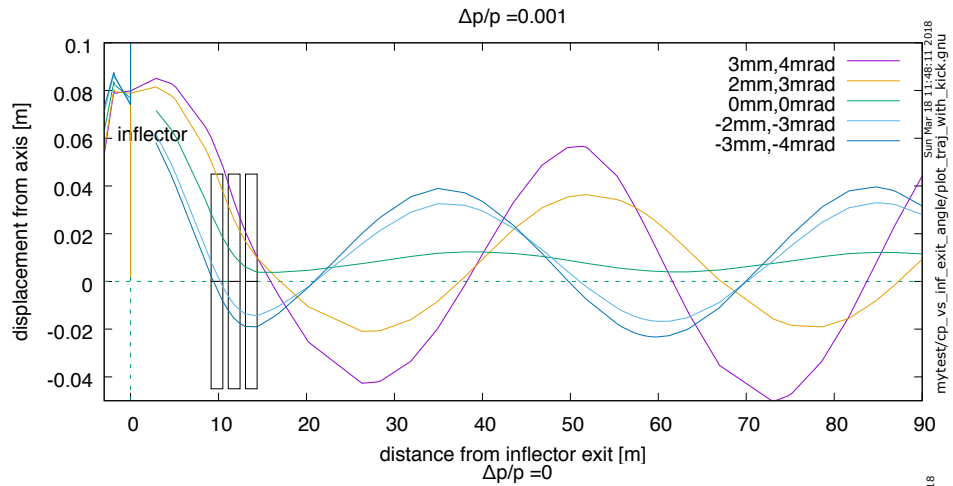




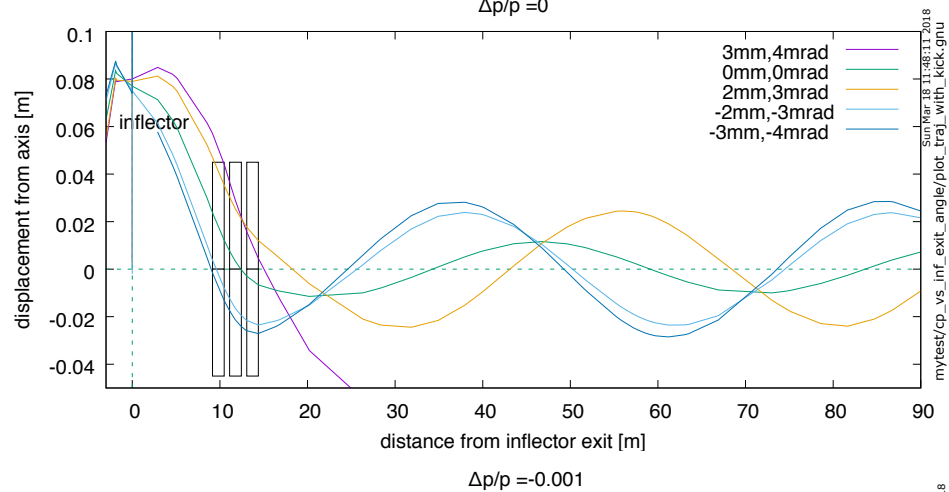
Radially outward trajectory ($\theta_{inj} > 0$)



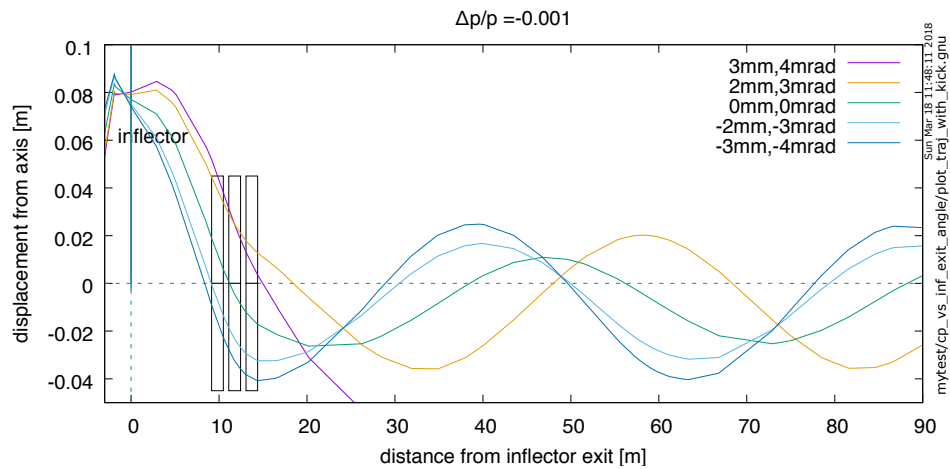
$$\Delta p/p = +0.001$$



Magic momentum



$$\Delta p/p = -0.001$$



Summary

What we know (or suspect)

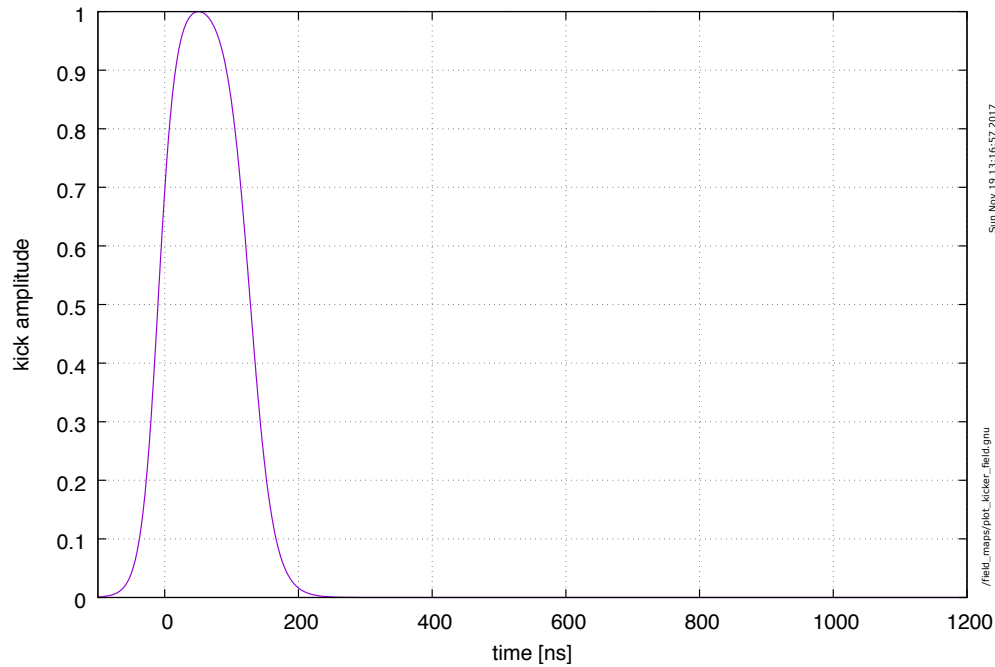
- Inflector is tilted clockwise about exit (-2.3 mrad)
- Beam trajectory from end of M5 and into backleg and inflector is directed radially outwards by a few mrad
 - Beam exiting inflector is likewise directed radially outwards
- 'Large' angle at inflector exit with respect to tangent line
 - CBO amplitude depends weakly on kick
 - Momentum acceptance skewed high
 - Optimal kick is 'high'
- Optimal inflector current (inflector exit angle) is high
- Kicker field is ?

There are alternative scenarios (parameter sets) that are equally compelling;

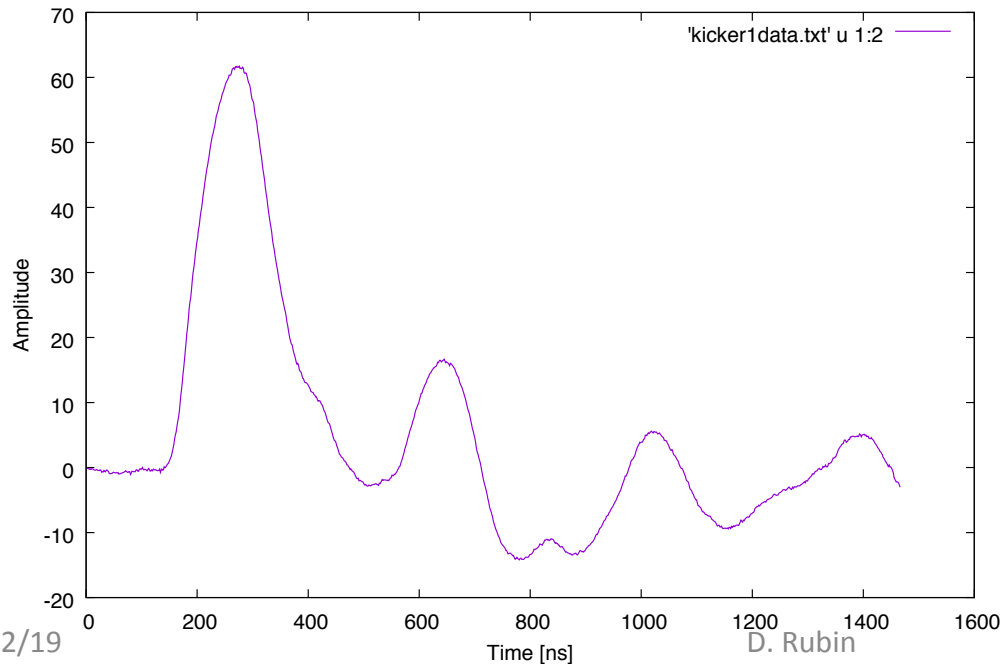
- *trajectory exiting inflector is radially inward and kicker field is low*
- *Trajectory is straight ahead and kicker field is very low*

With our now stable conditions

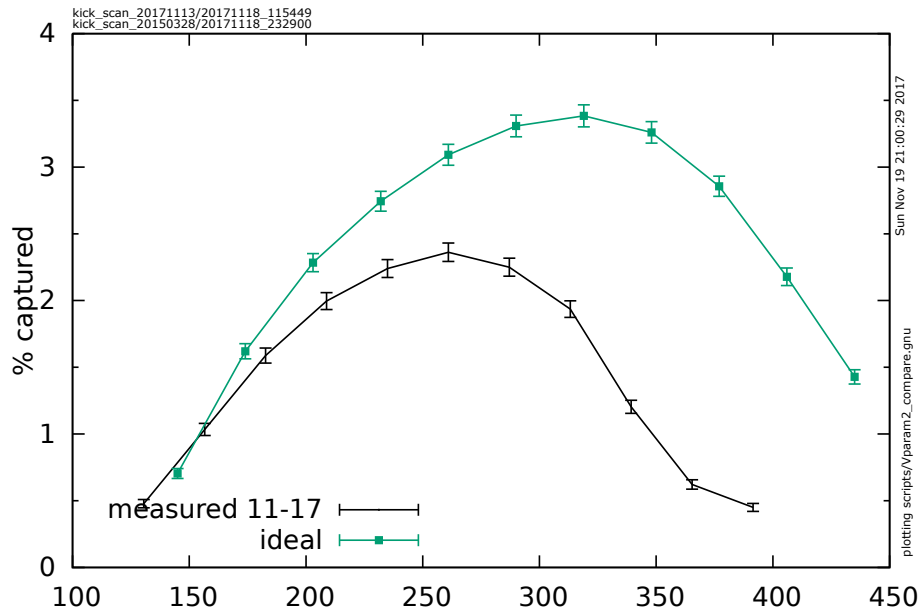
Data from dedicated scans of inflector field (exit angle) and kicker voltage, with fiber harps would help to constrain the phase space of solutions



'Ideal' kicker pulse

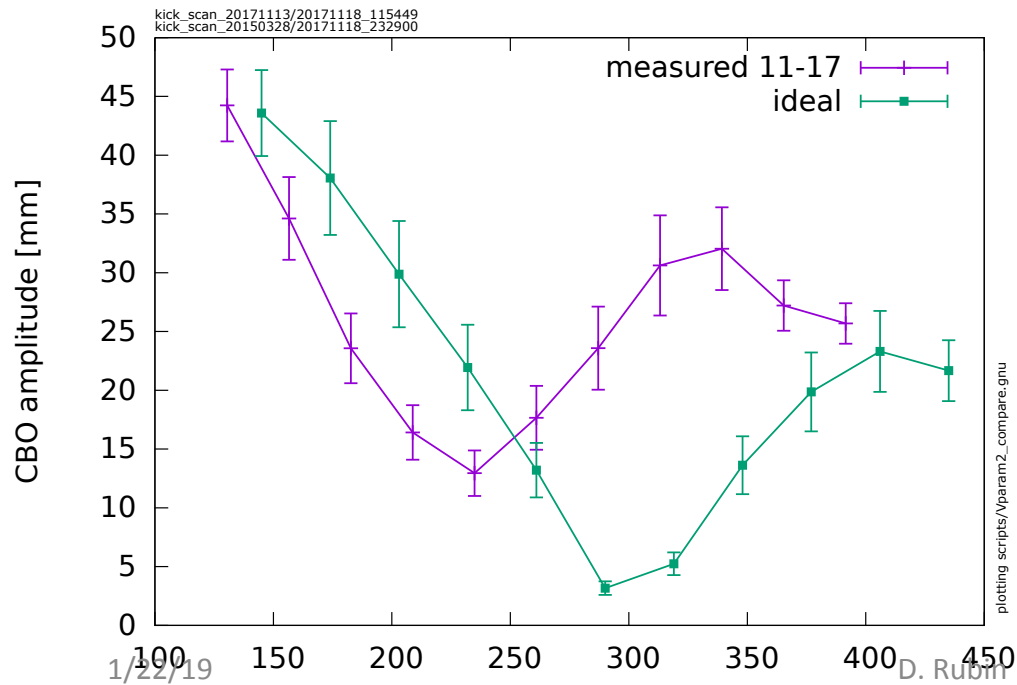


Magnetometer
measurement of kicker 1

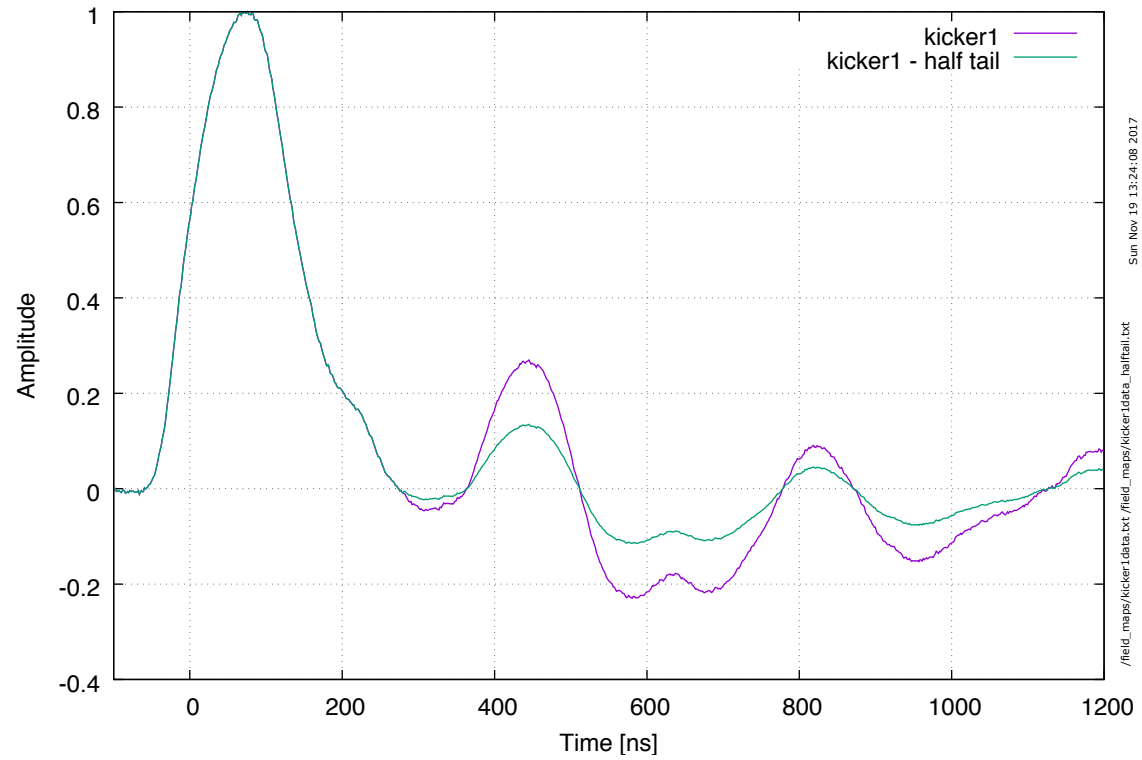


Measured/ideal = 70%

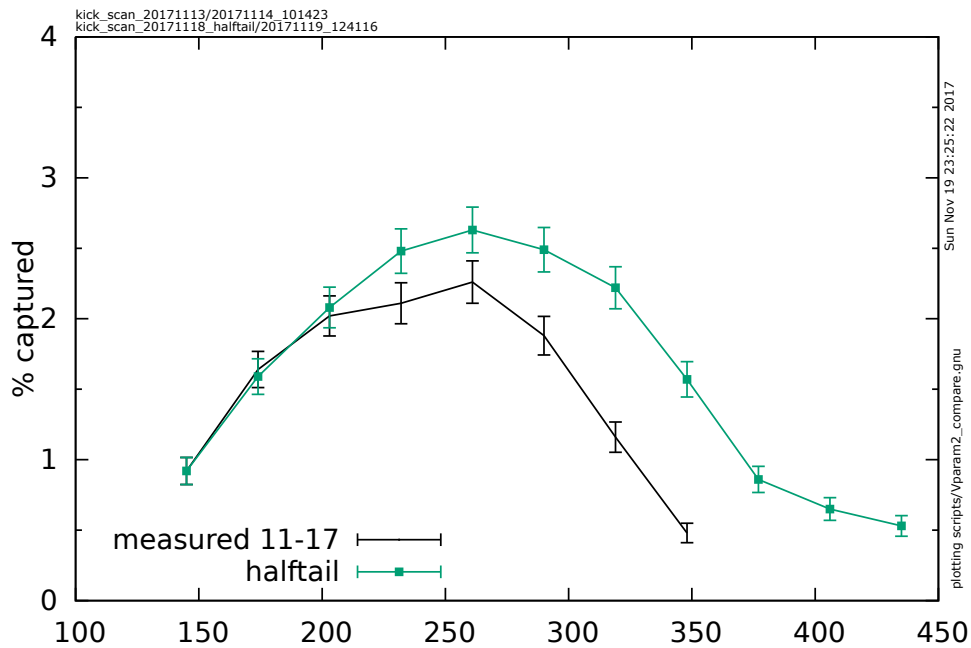
Capture vs pulse shape



Residual CBO vs pulse shape



Suppose we are able to reduce the secondary pulses by 1/2



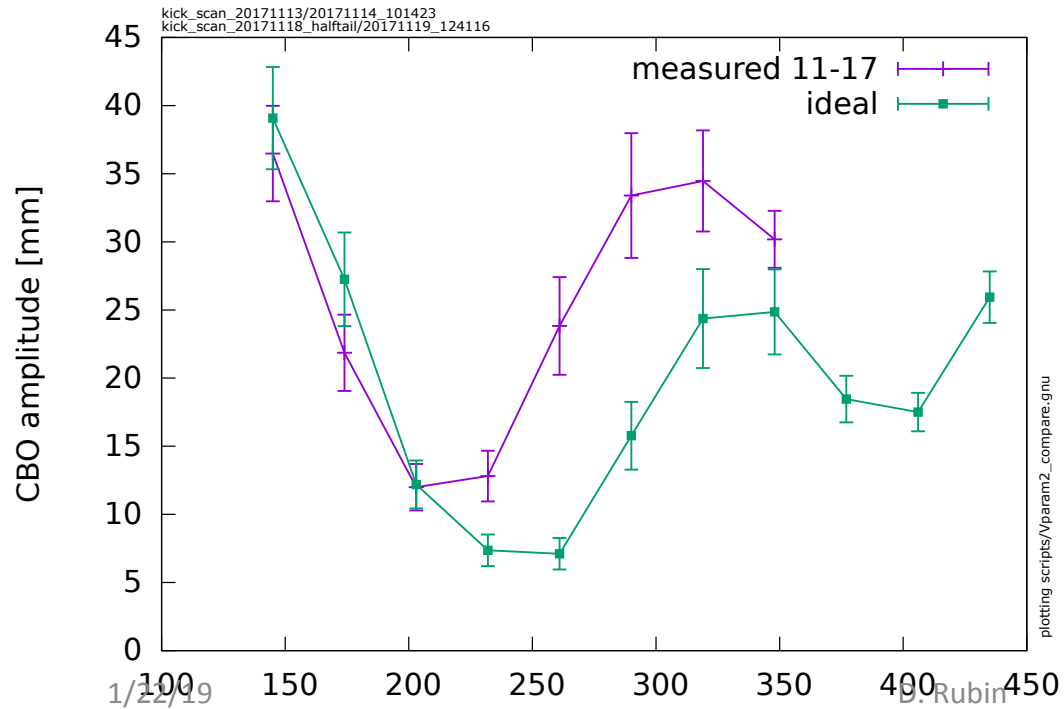
Compare

- Measured pulse shape
- Measured with 1/2 tail

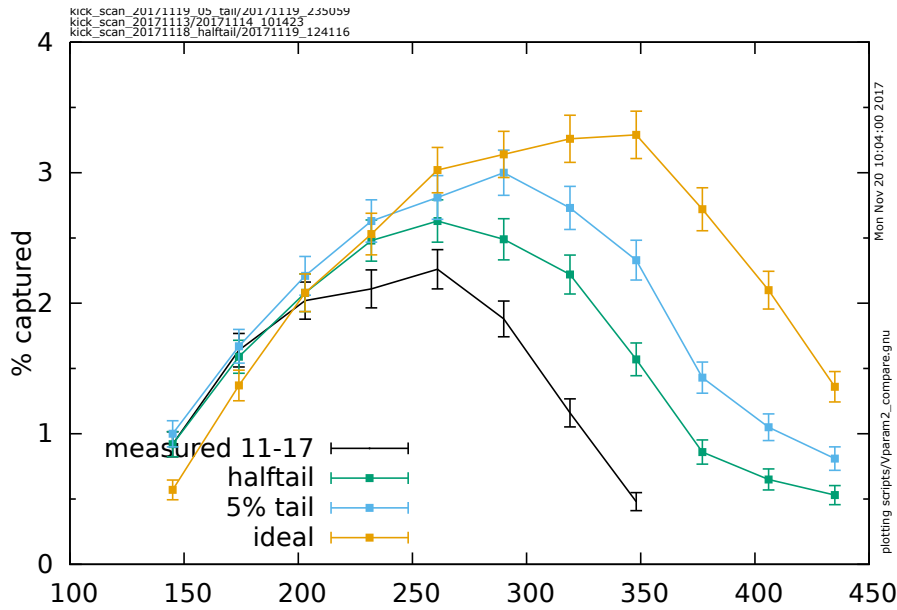
Capture efficiency

Reduced tail/ideal = 78%

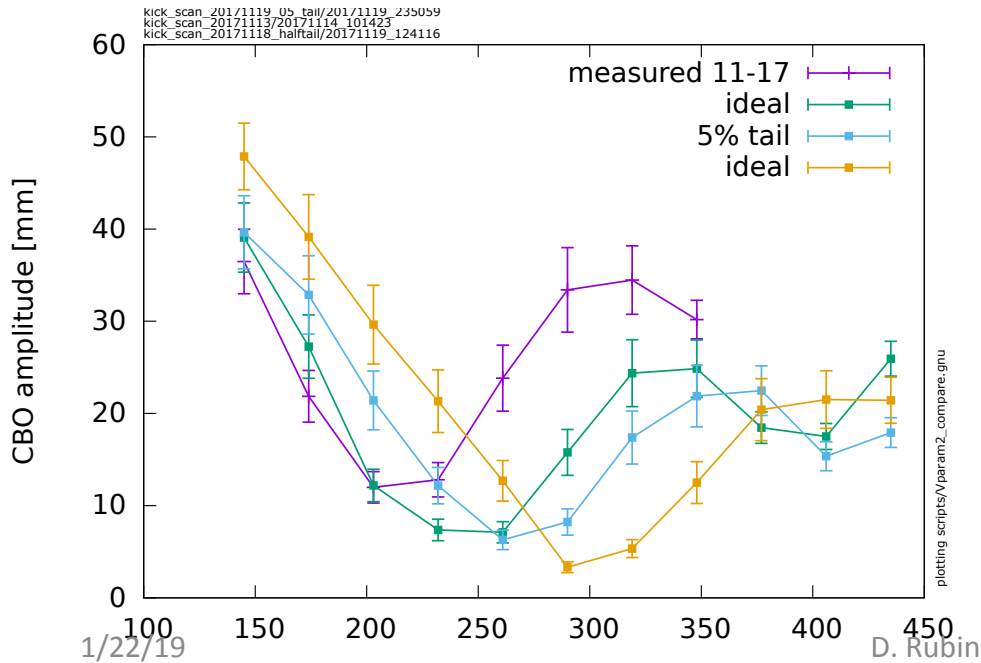
Measured/reduced tail = 86%



Residual CBO

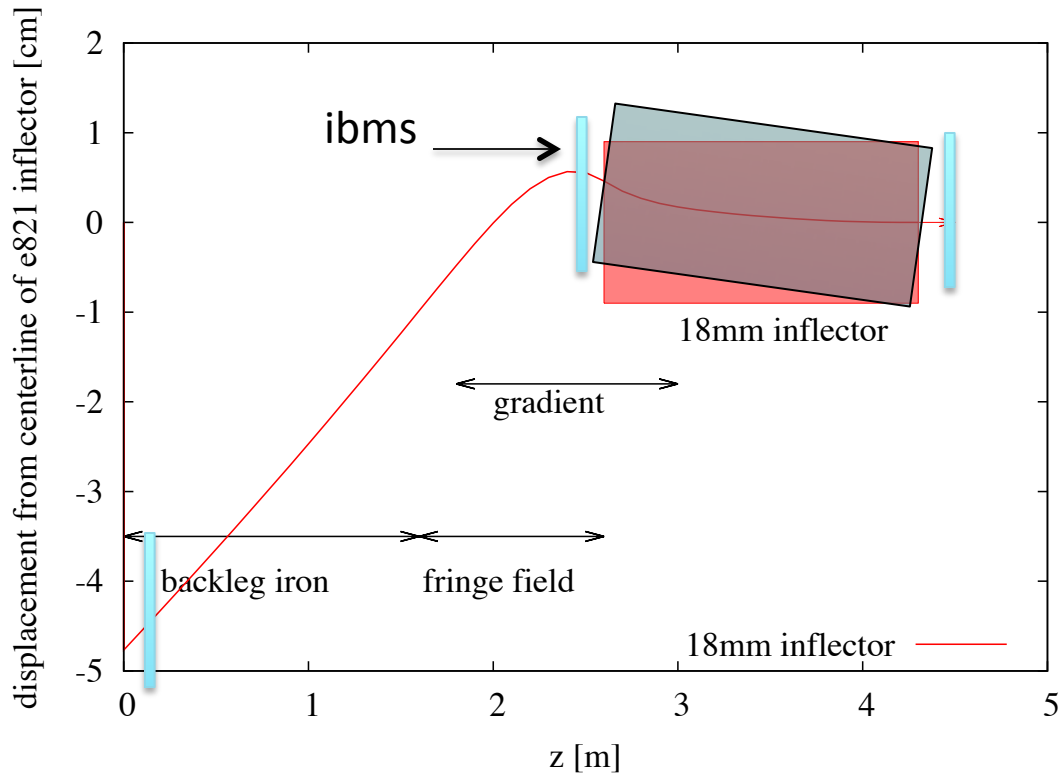


Capture vs pulse shape

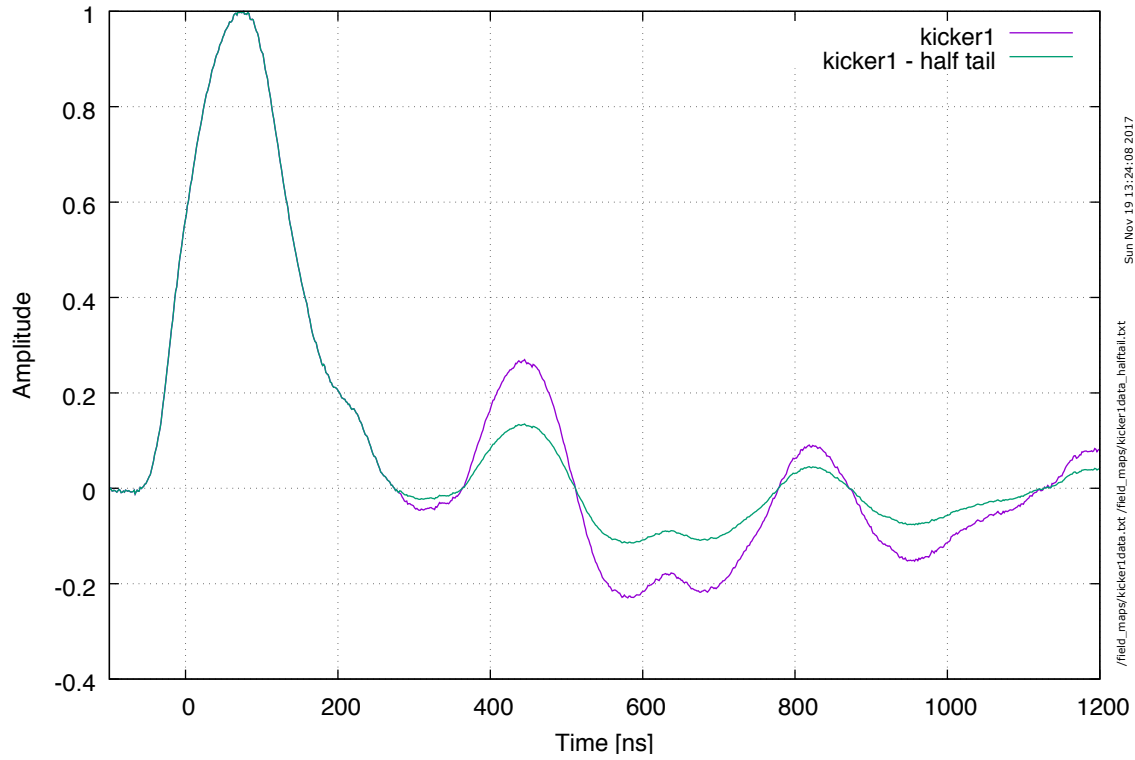


Residual CBO vs pulse shape

Trajectory through injection channel



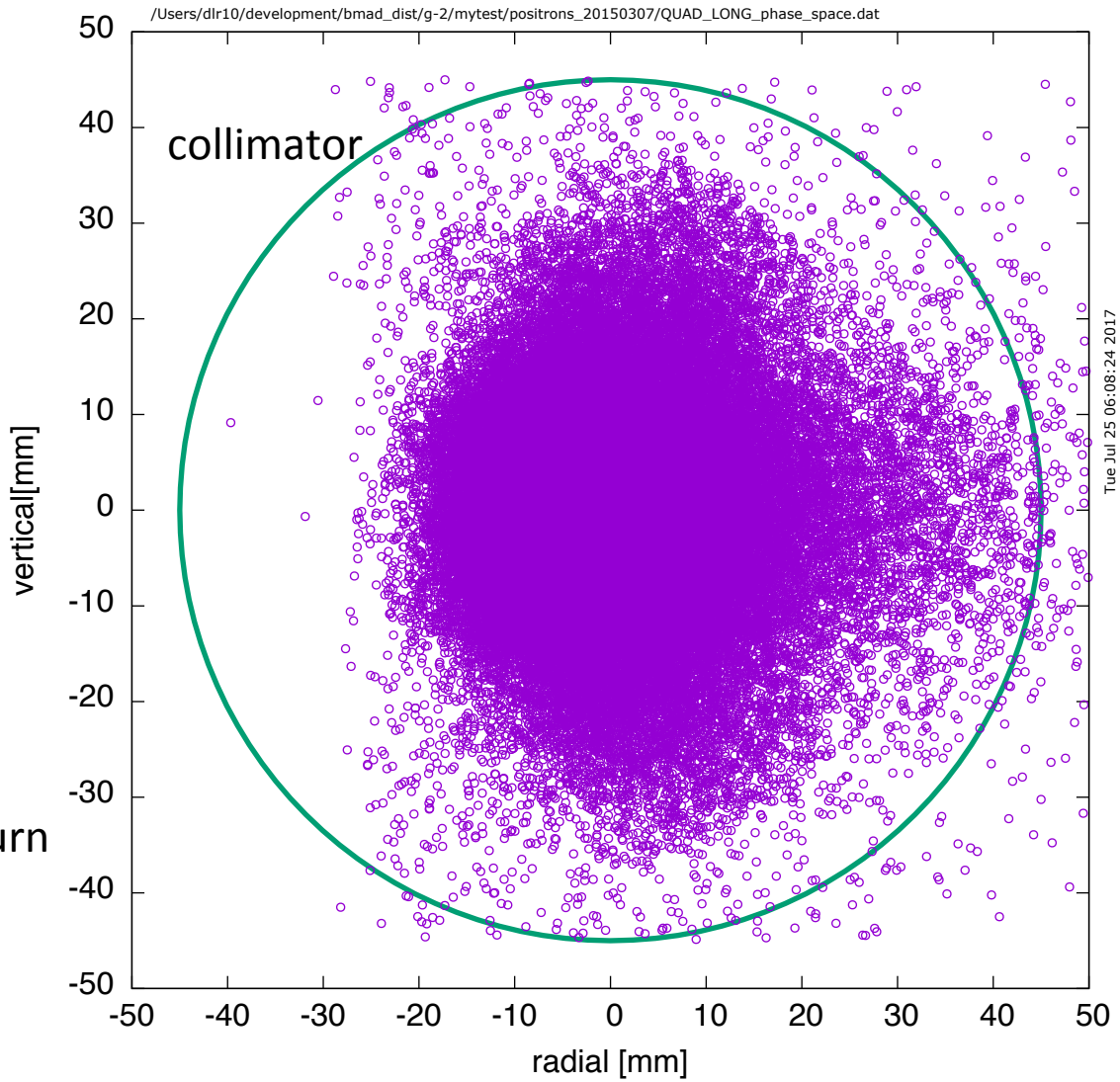
Beam exits centered in inflector
tangent to and displaced 77mm
from ring central orbit



Sun Nov 19 13:24:08 2017

/field_maps/kicker1data.txt /field_maps/kicker1data_half tail.txt

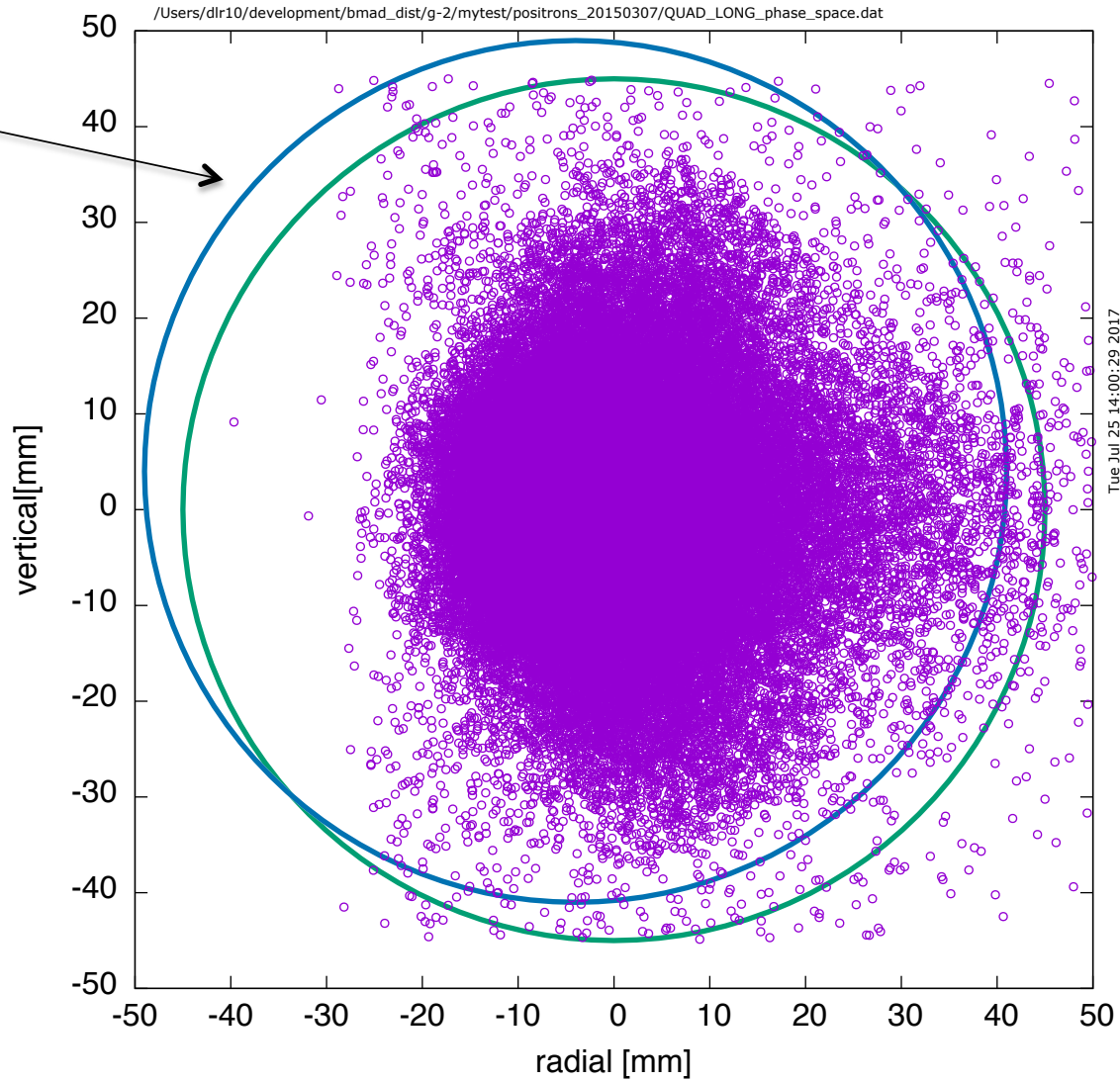
Scraping

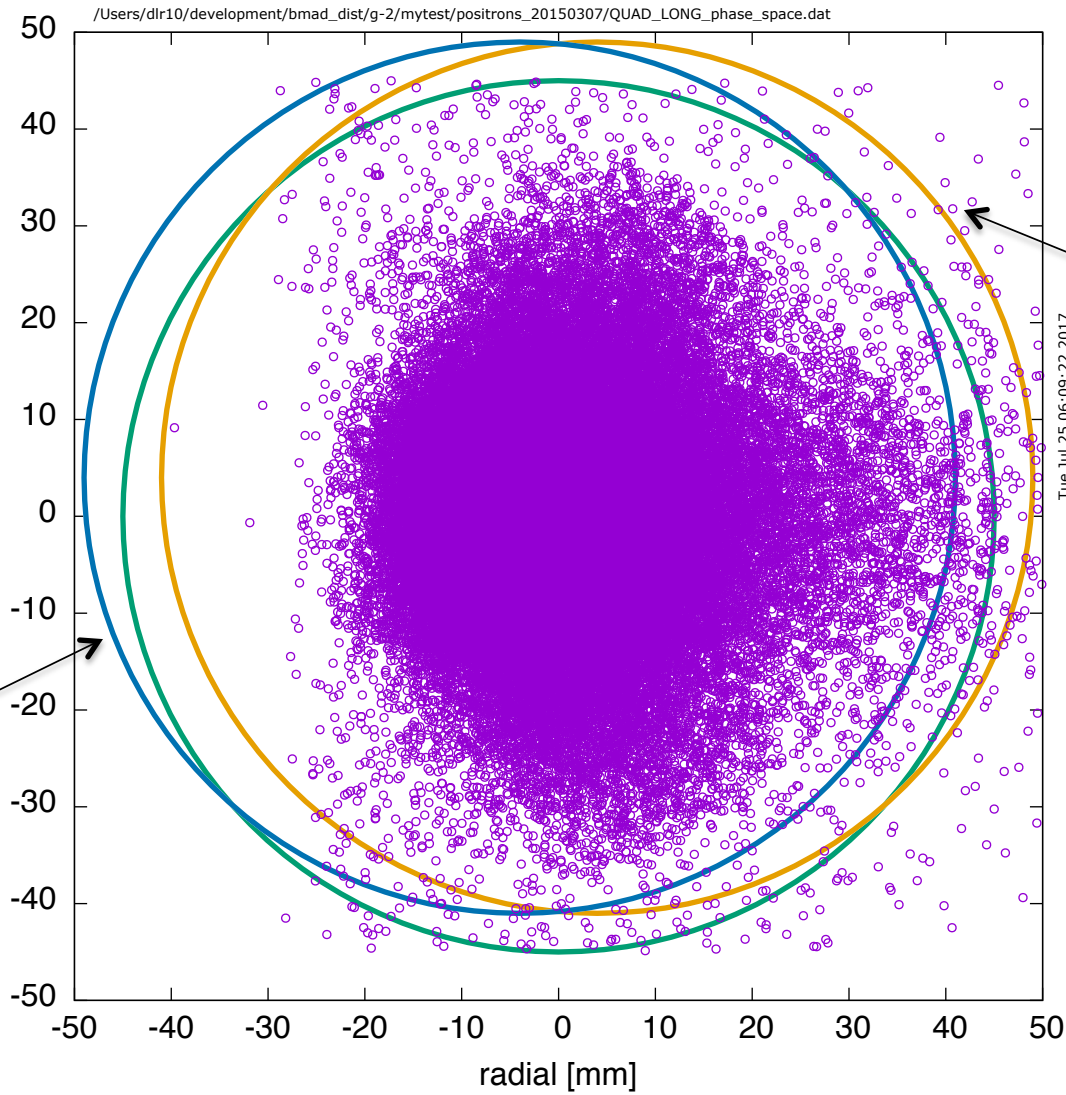


Distribution on first turn

Displace collimator to scrape off tails

Displaced
collimator





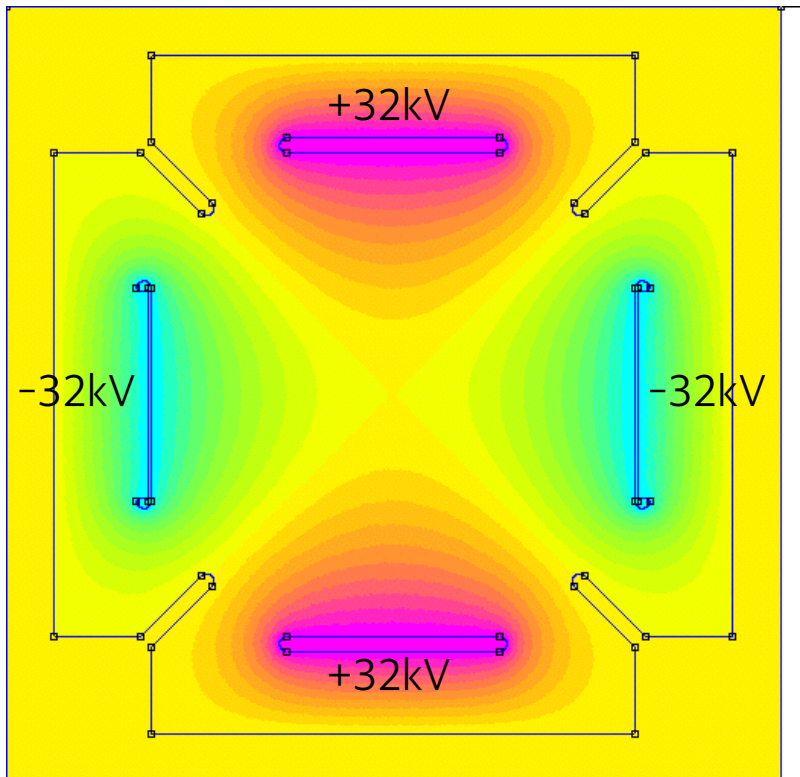
Displaced
radially out
and up

Displaced
radially in
and up

It is difficult to move the scrapers

Instead we displace the muon closed orbit

Nominal field configuration



Reduce voltage on bottom plate
and orbit shifts downward

Reduce voltage on inner(outer) plate
and orbit shifts out(in)

Table 10

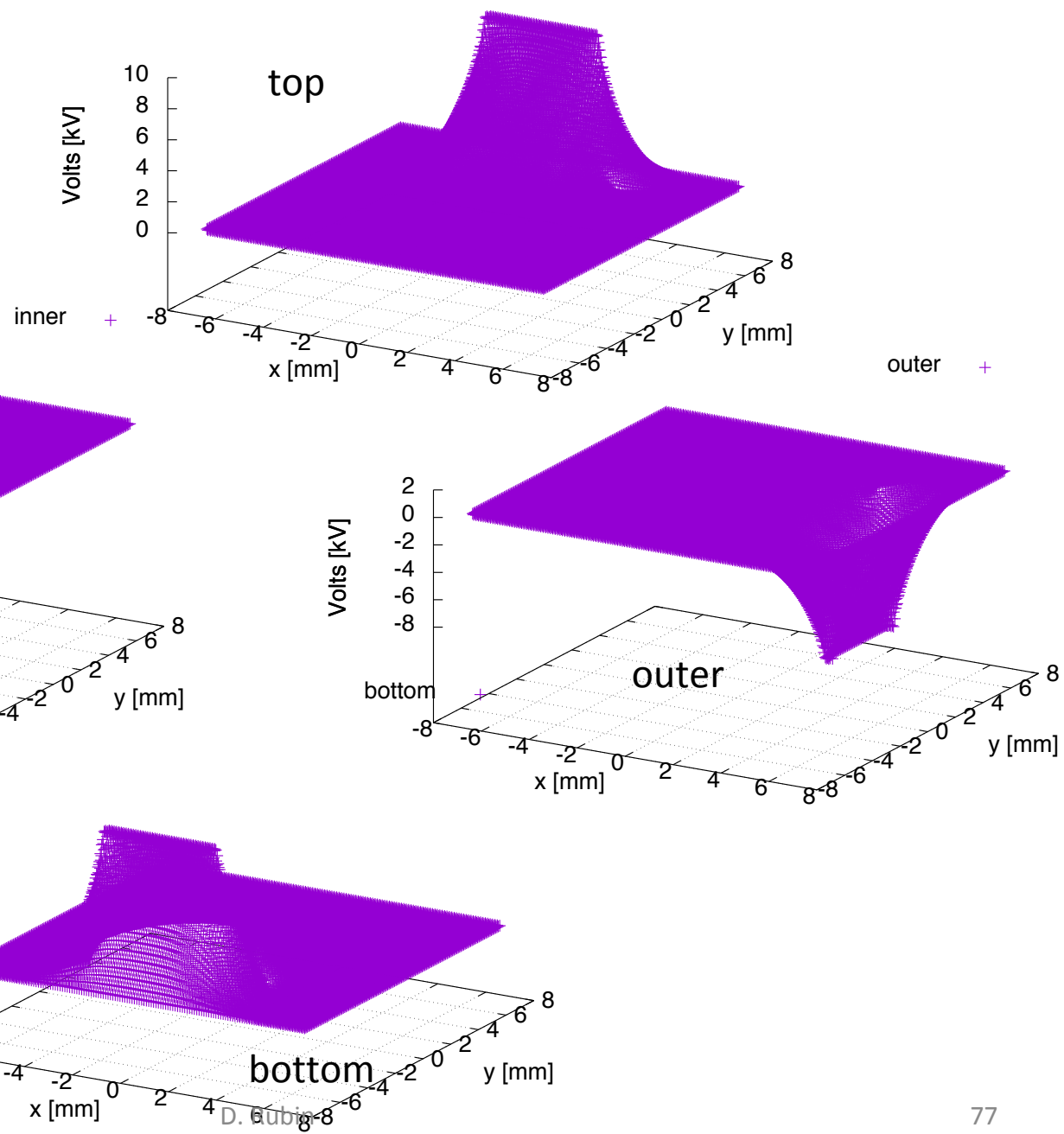
Quadrupole plate voltage [kV] at injection time for the negative muon storage polarity

Quad section	Top plates	Bottom plates	Inner plates	Outer plates
Q ₁	-24	-17	+24	+24
Q ₂	-24	-17	+17	+24
Q ₃	-24	-17	+24	+24
Q ₄	-24	-17	+24	+17

Temporary values (kV) for E989

	top	bottom	inner	outer
Q1	+32	+22.7	-32	-32
Q2	+32	+22.7	-22.7	-32
Q3	+32	+22.7	-32	-32
Q4	+32	+22.7	-32	-22.7

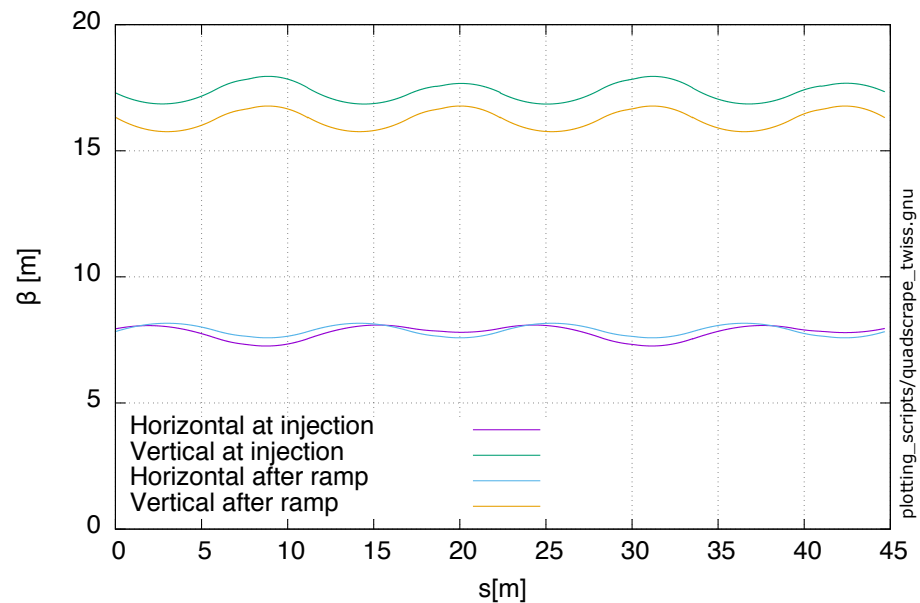
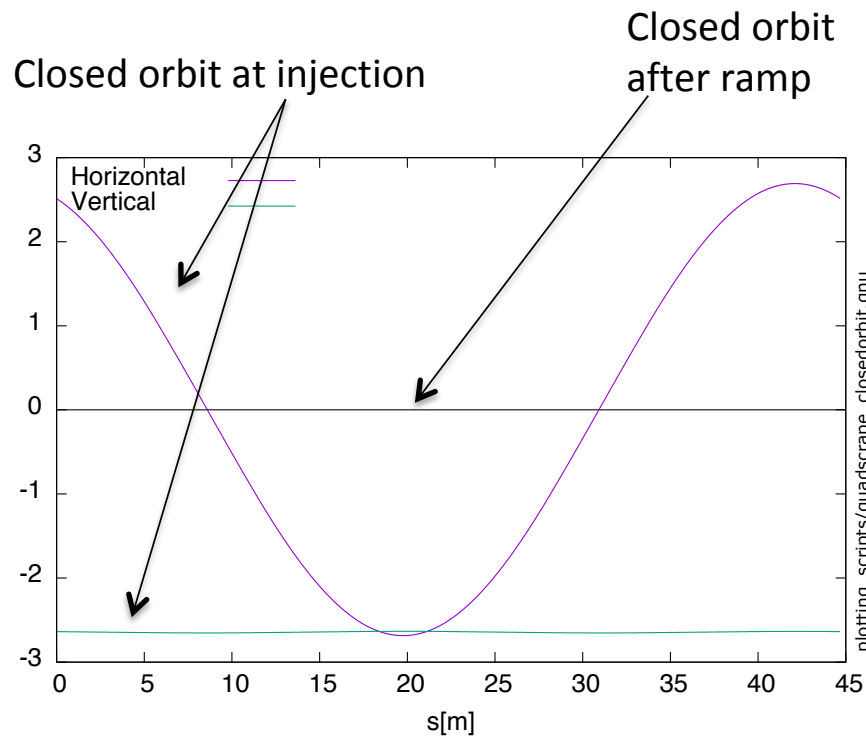
Potential map for each of four quad plates



Closed orbit, tunes, and β at injection and after ramp

Quad voltage at injection

	Top	Bottom	Inner	Outer
Q1	32	22.7	-32	-32
Q2	32	22.7	-22.7	-32
Q3	32	22.7	-32	-32
Q4	32	22.7	-32	-22.7

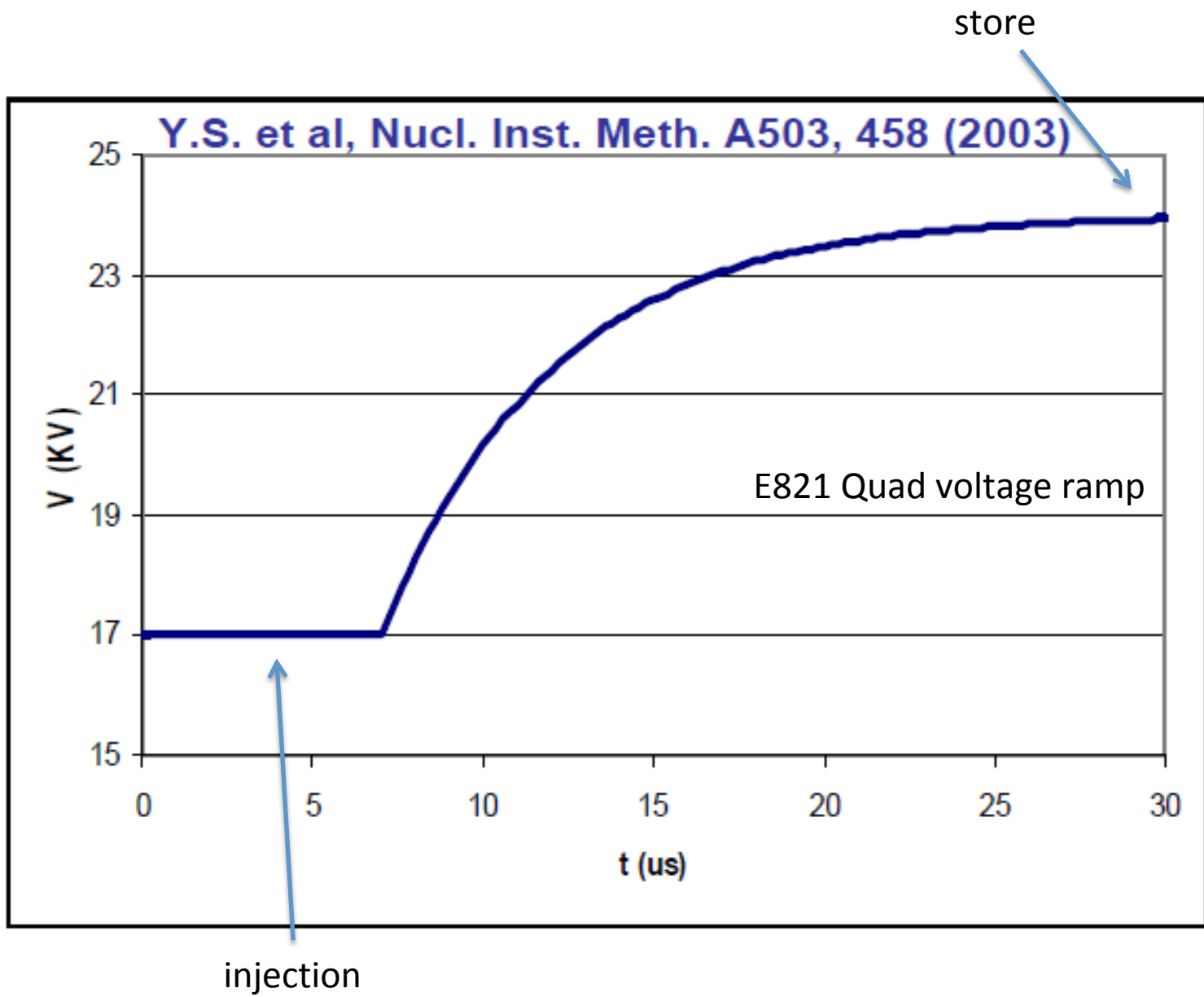


Tunes at Injection

Q_x 0.9137
 Q_y 0.4090

Tunes after ramp

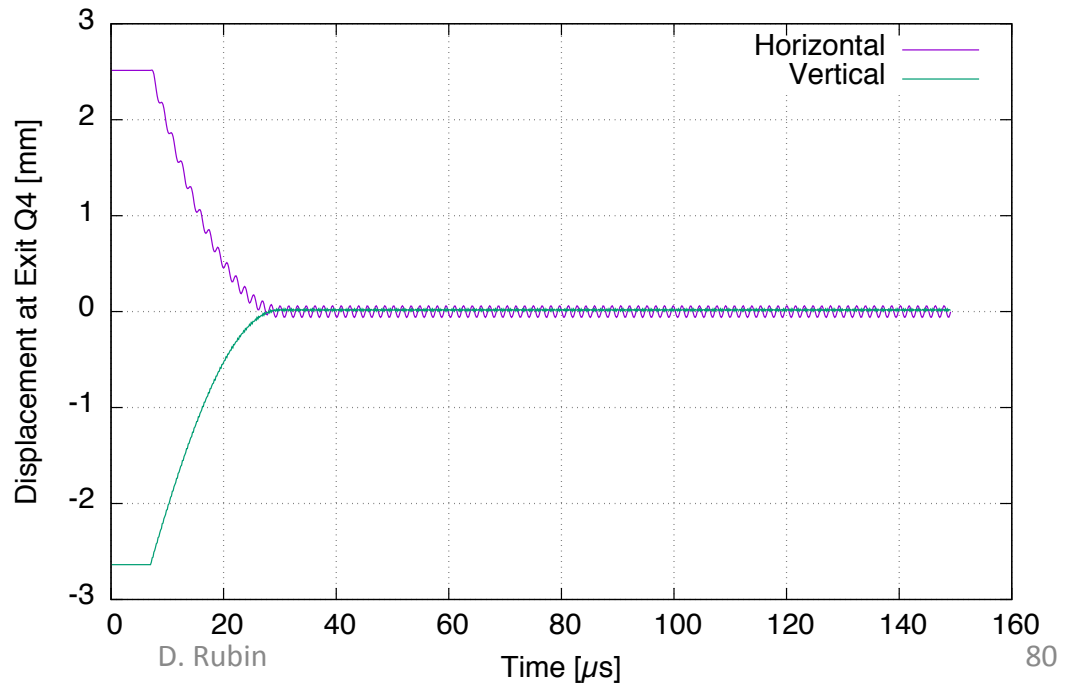
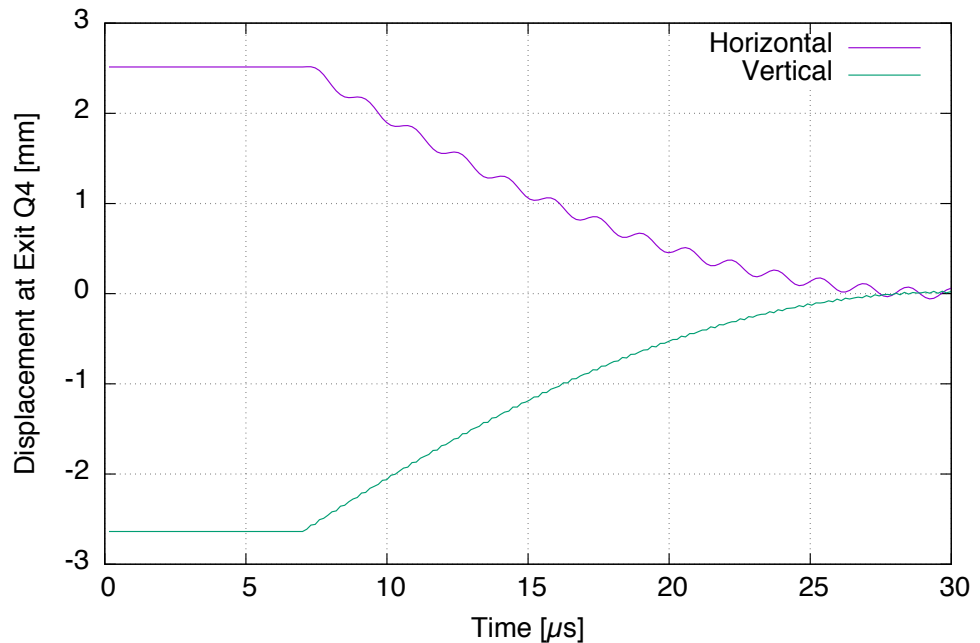
0.9039
 0.4403



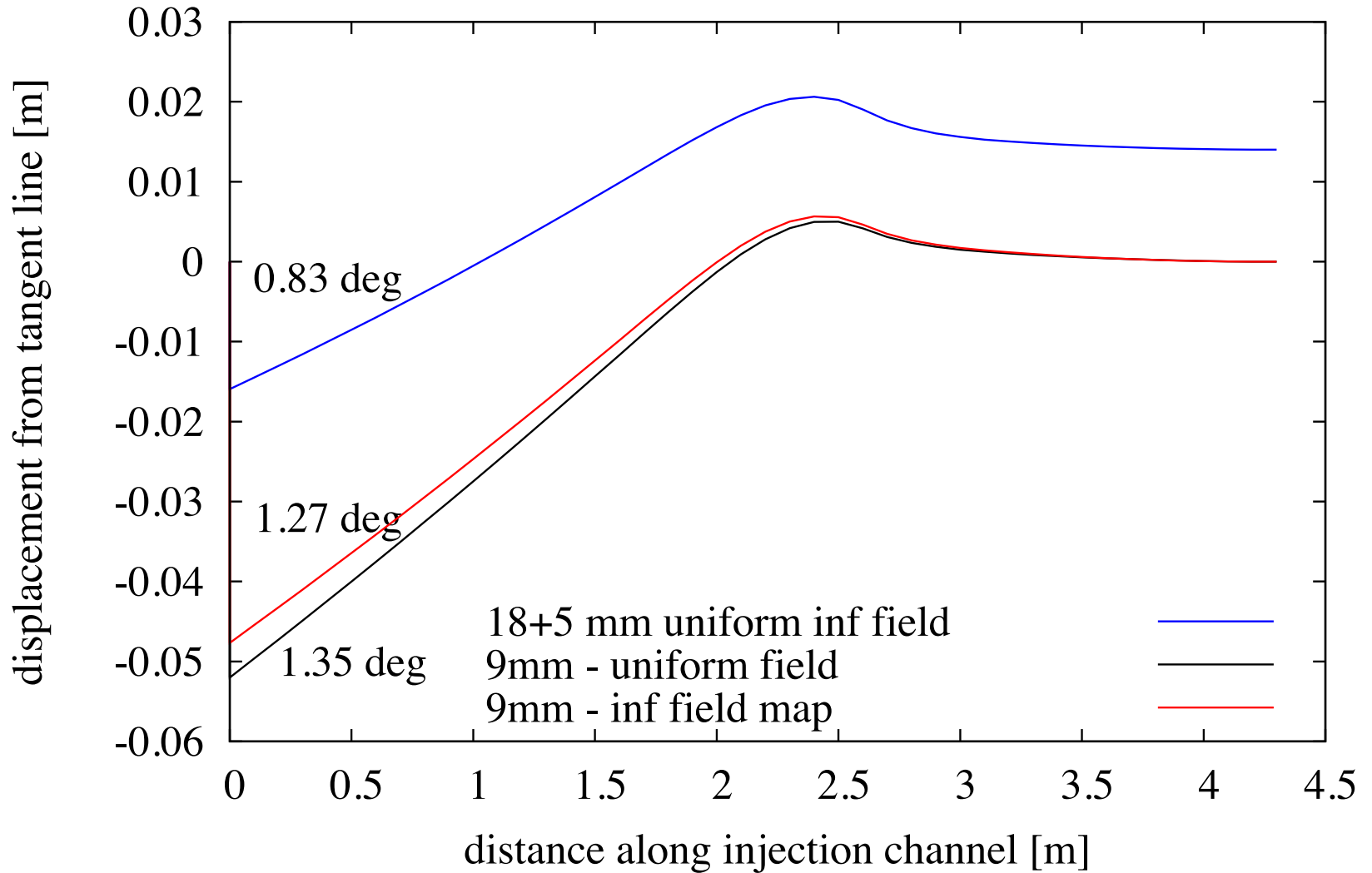
- At $t=0$, inject muons on to displaced closed orbit?
- Track 1000 turns
- Quad plate voltage ramps as in E821
- (Plot shows orbit at exit of Q4 on each turn)

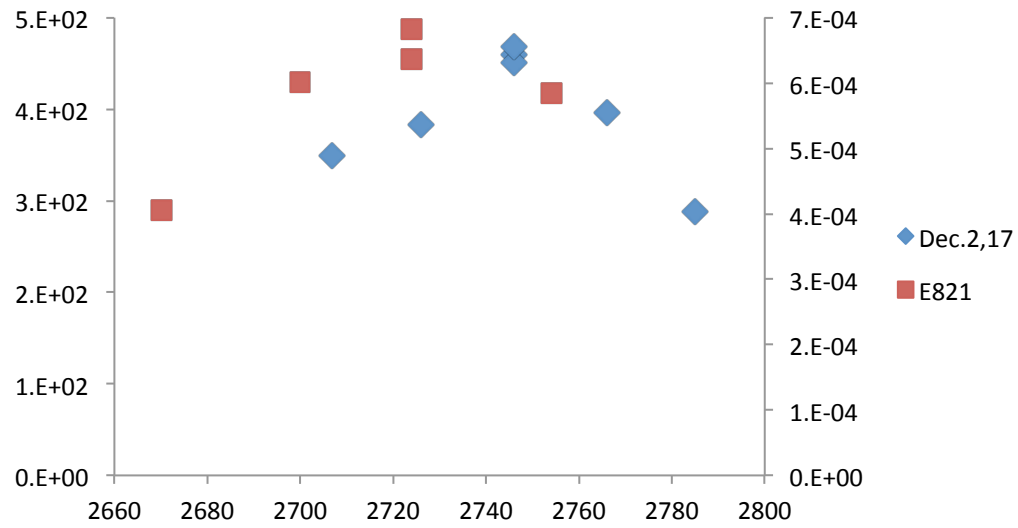
Orbit shifted to center of aperture as quads ramp to nominal (symmetric voltage)

A small horizontal betatron oscillation is introduced by the ramp



Trajectory through injection channel






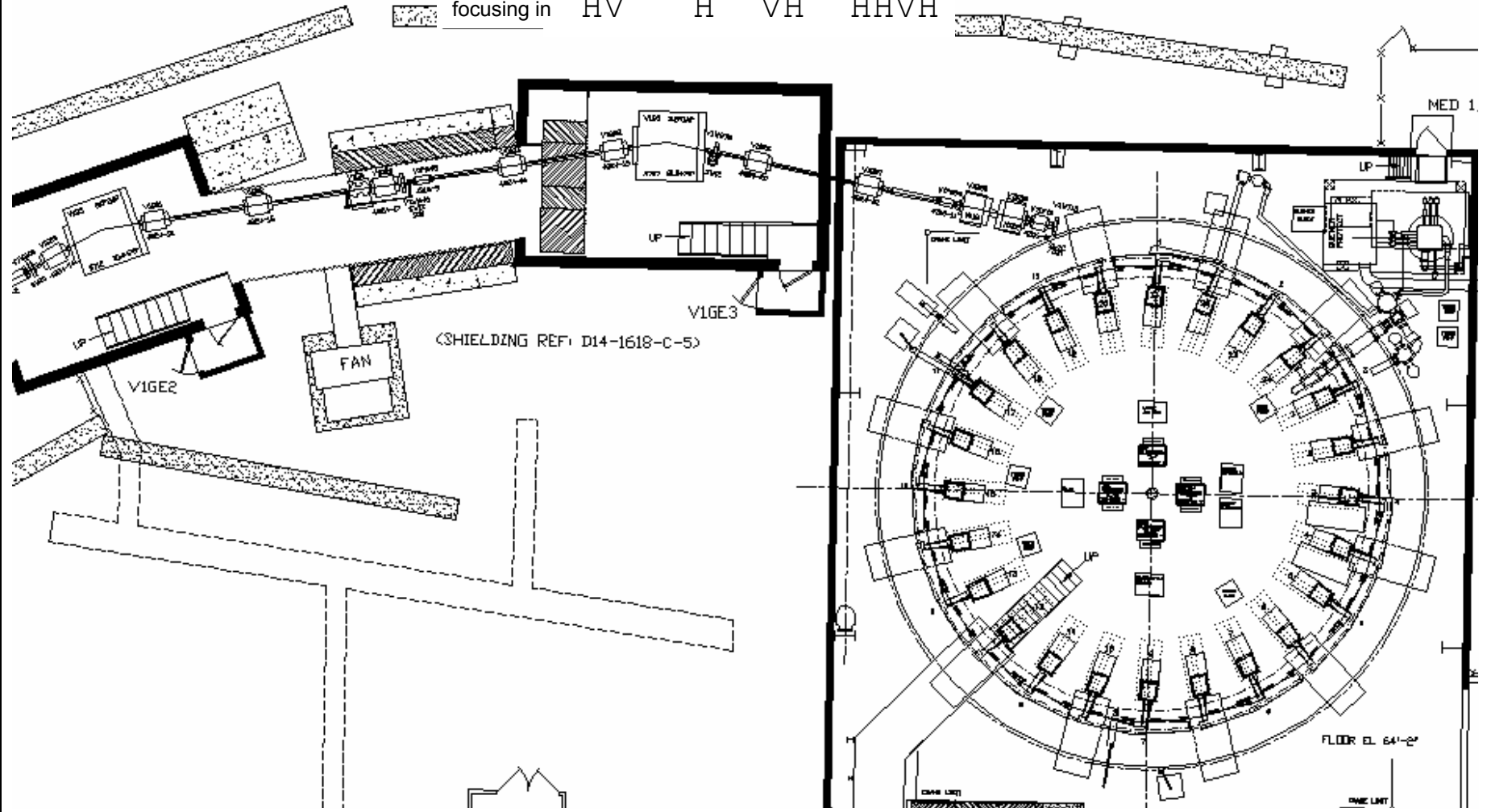
Higher inflector field compensates positive inflector exit angle

E 821

V line D5 to g-2 ring

DQQ	Q	QQD	QQQQ
HV	H	VH	HHVH

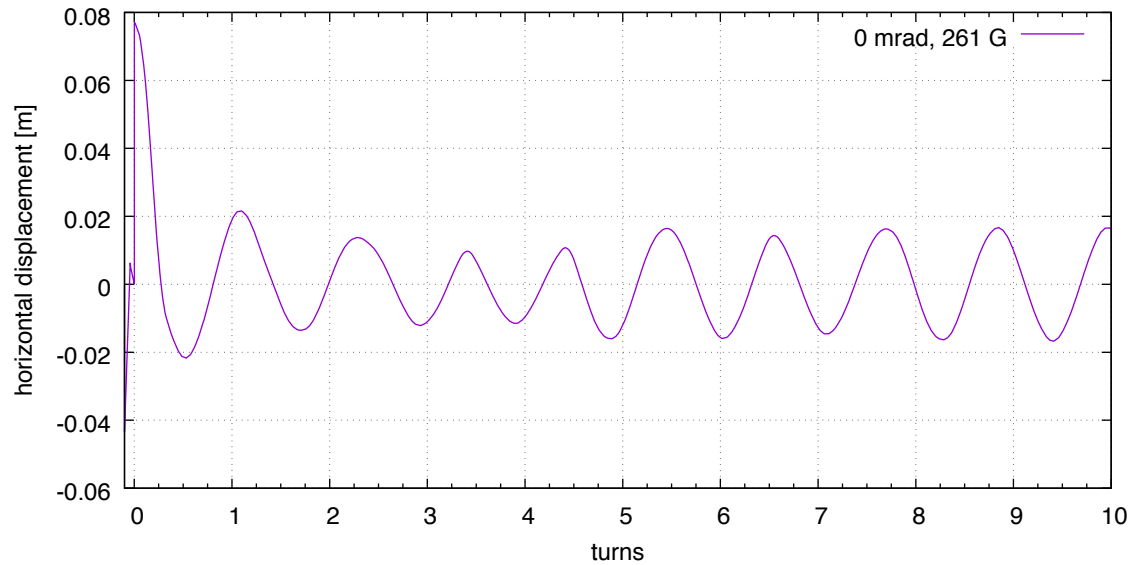
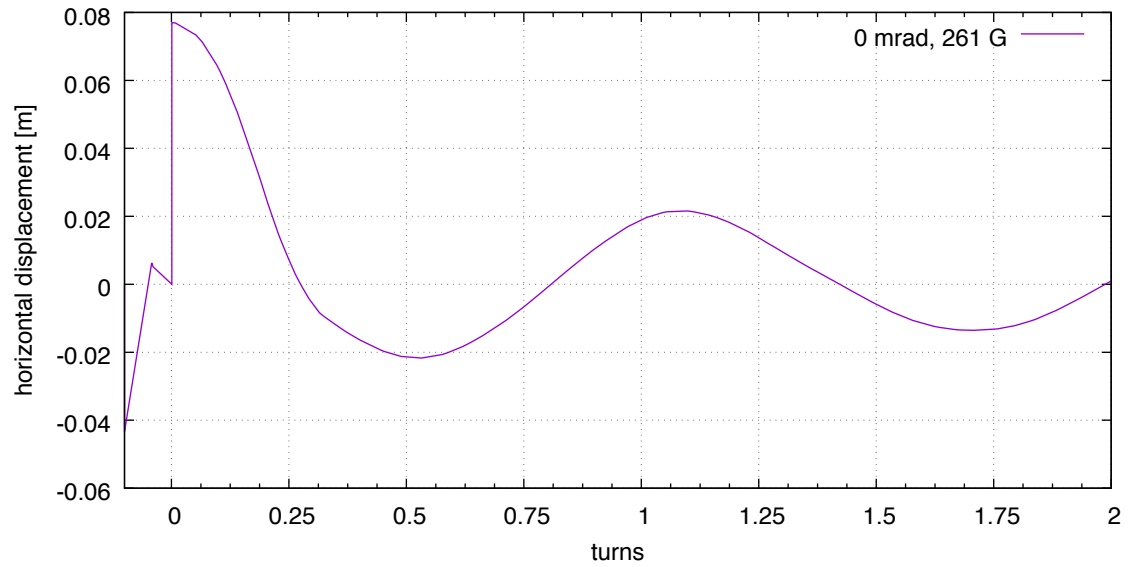
 focusing in



Optimum kick (minimize CBO) vs inflector exit angle – on momentum muon

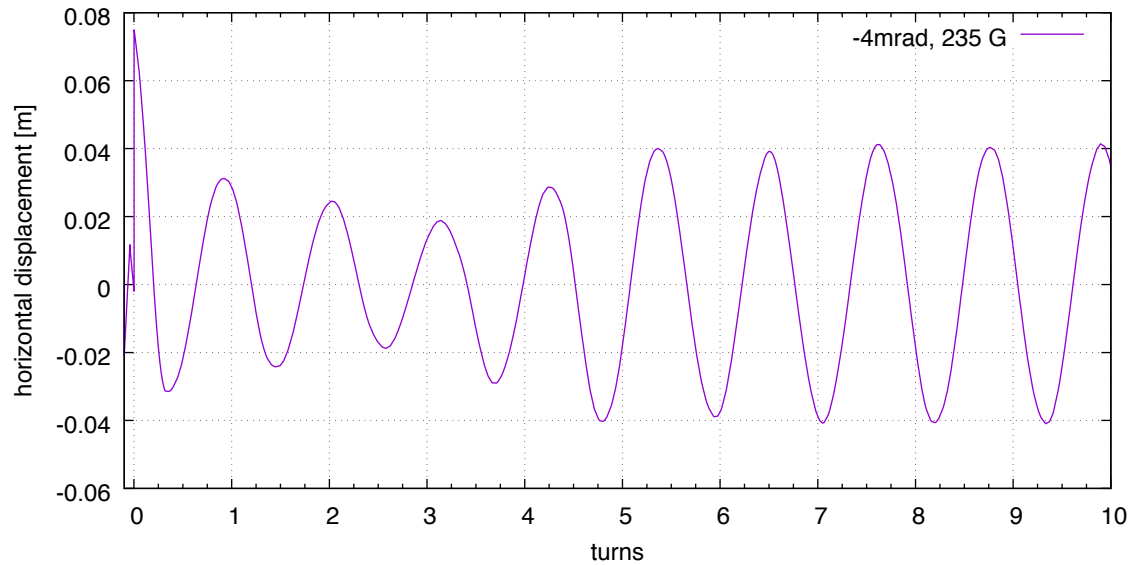
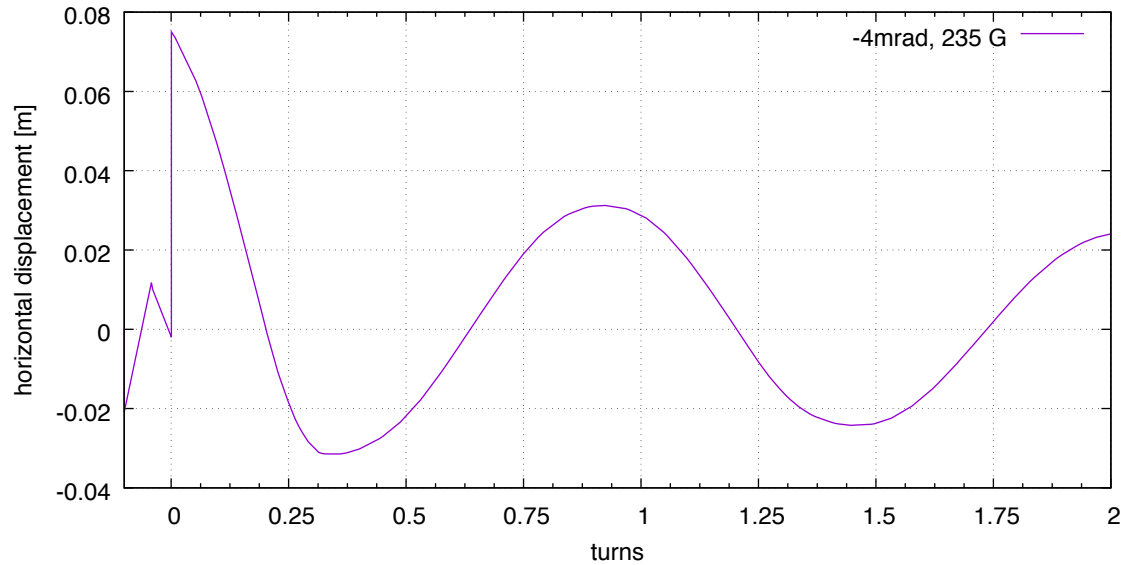
$\theta=0$

Muon trajectory is
along tangent line
(inflector center line)



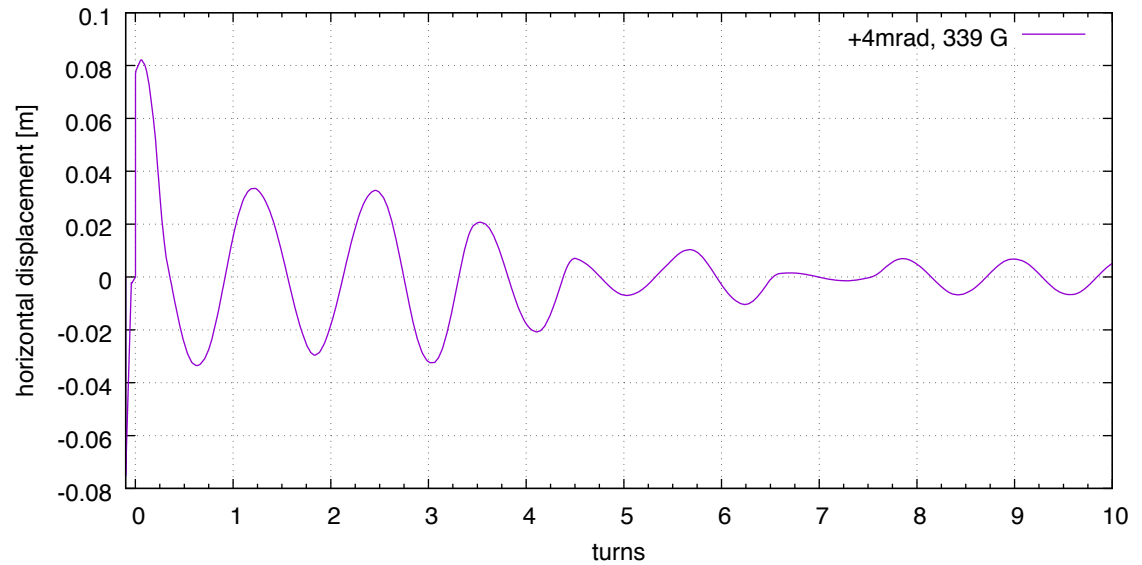
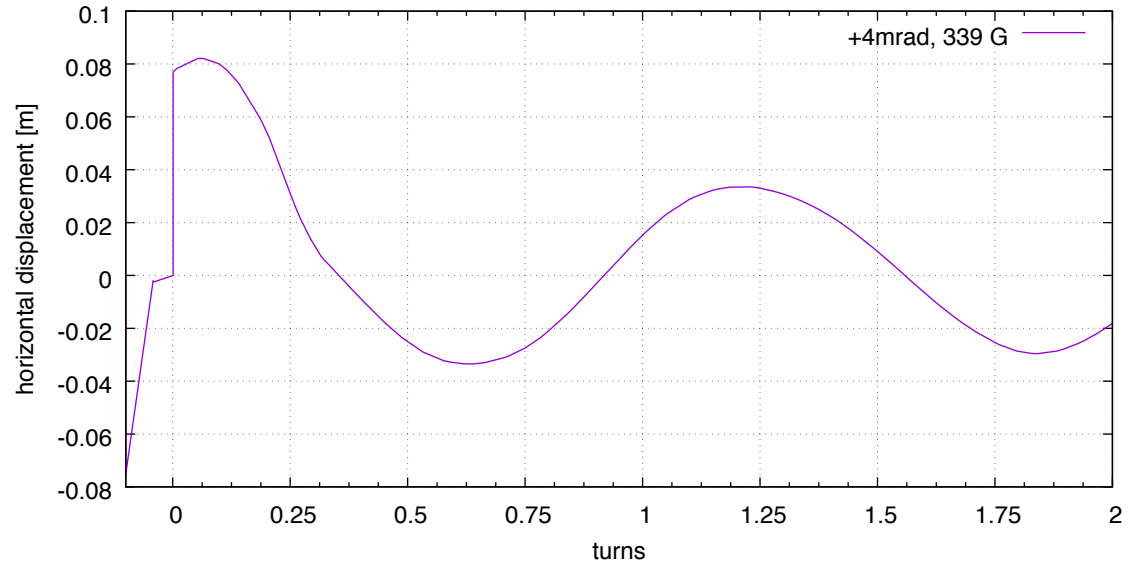
Optimum kick (minimize CBO) vs inflector exit angle – on momentum muon

Muon trajectory is tilted
-4mrad with respect to
tangent line (inflector
center line)

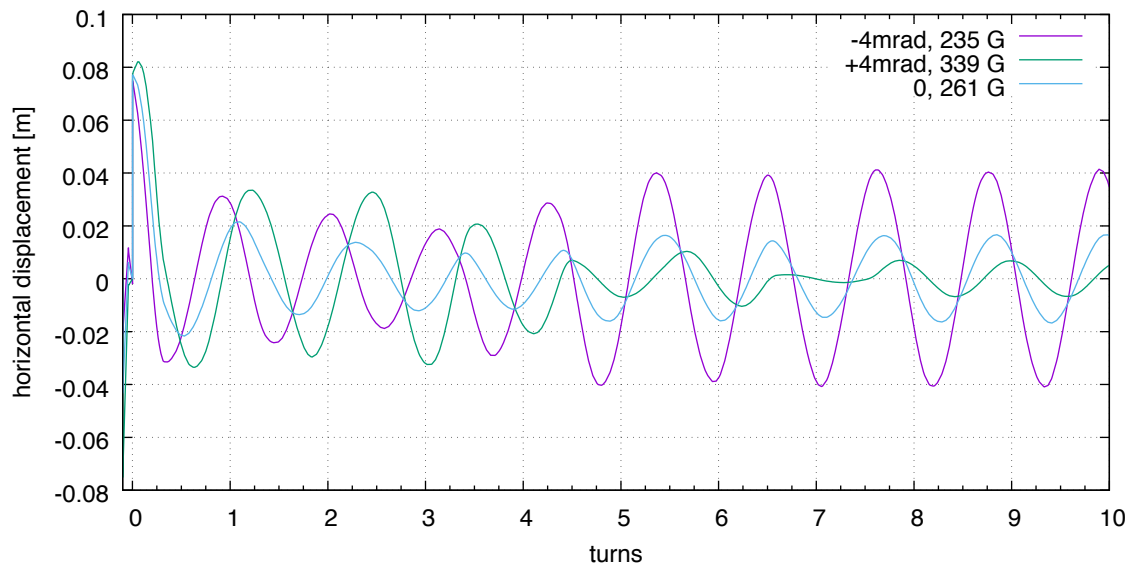
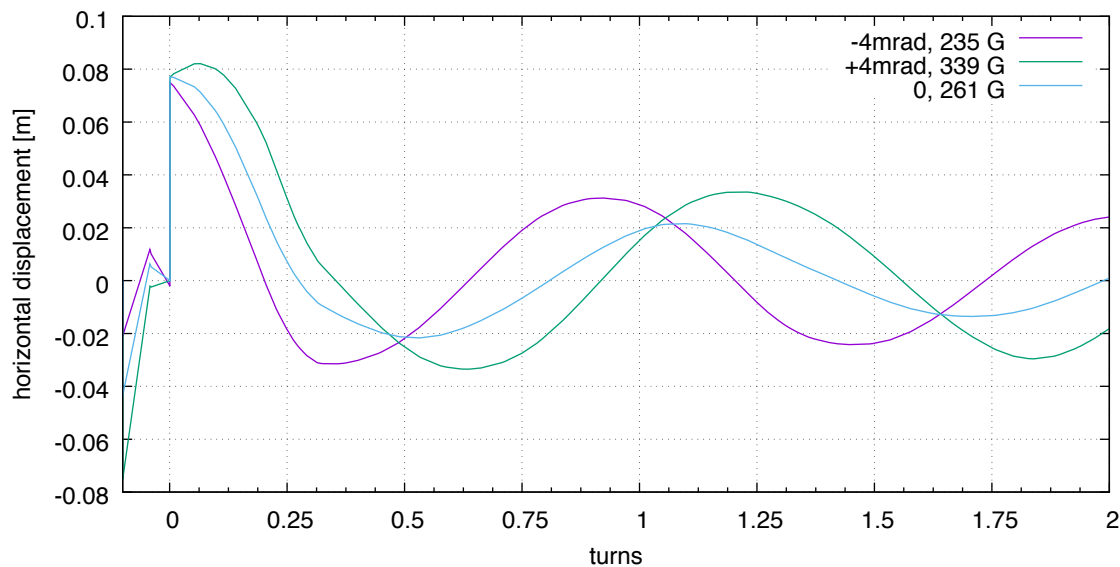


Optimum kick (minimize CBO) vs inflector exit angle – on momentum muon

Muon trajectory is t
+4mrad with respect to
tangent line (inflector
center line)



Optimum kick (minimize CBO) vs inflector exit angle – on momentum muon



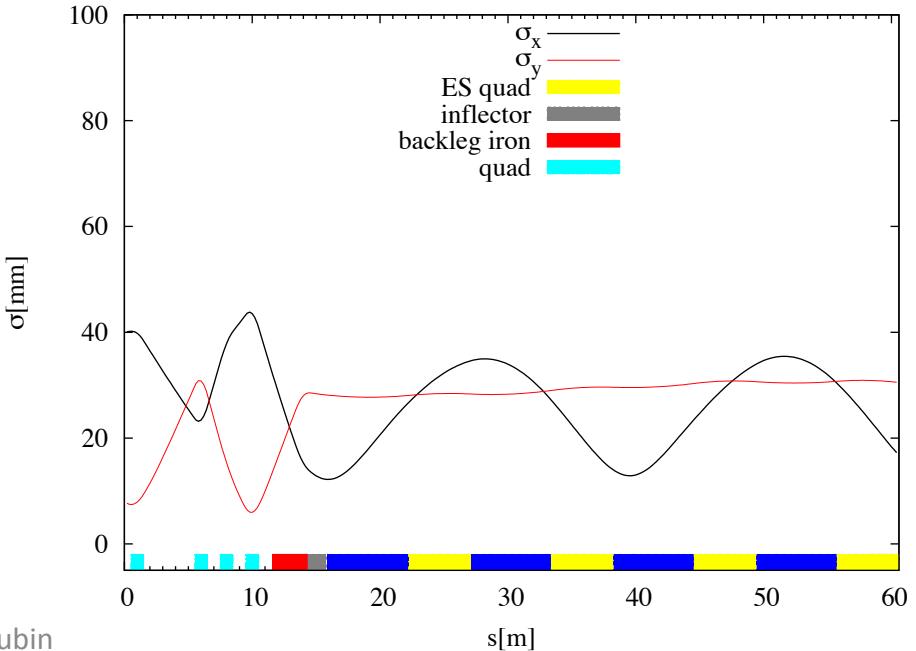
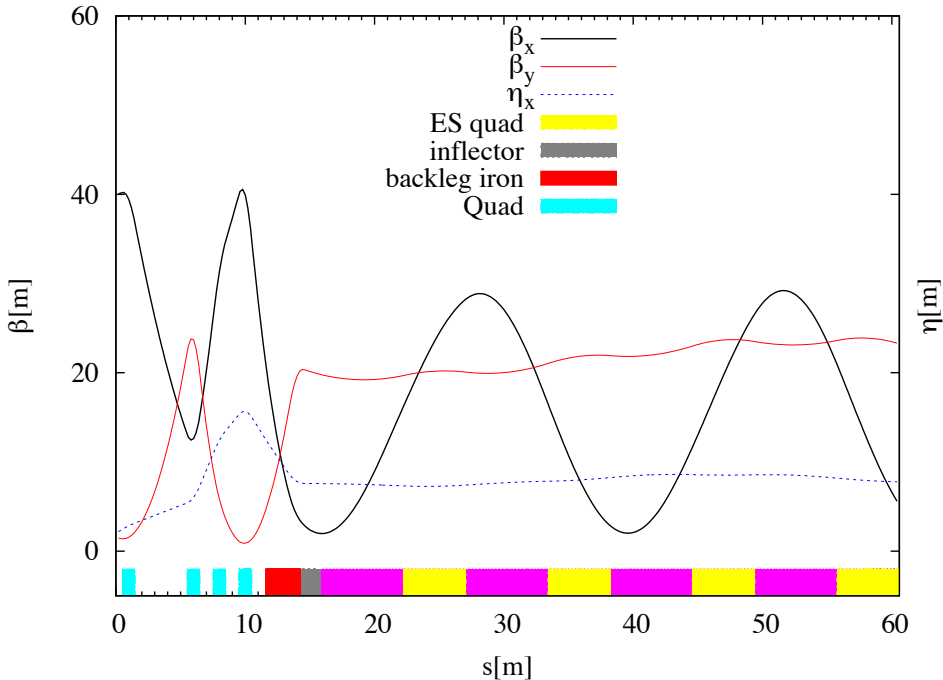
1. Dispersion at inflector exit matches dispersion into ring

- Ring aperture = $\pm 4\text{cm}$
- Inflector aperture = $\pm 9\text{mm}$
- Ring dispersion $\eta_0=8\text{m}$
- => Energy acceptance in ring $\Delta E/E = 0.5 \%$
- => Energy acceptance through inflector $\Delta E/E = 0.11 \%$

Real acceptance = 0.11 % and no losses in ring

$$\sigma_x = [\beta\varepsilon + (\eta\delta)^2]^{1/2}, \quad \delta=0.11\%$$

- Beam fits comfortably into ring aperture
- 0.11% energy acceptance
 - No energy modulation of beam size
 - “breathing” at twice the betatron tune



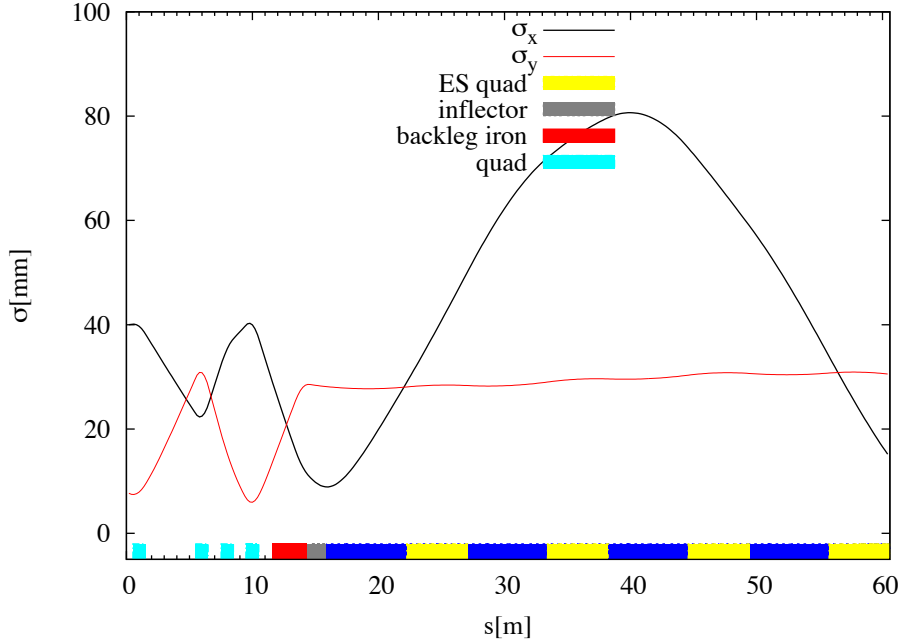
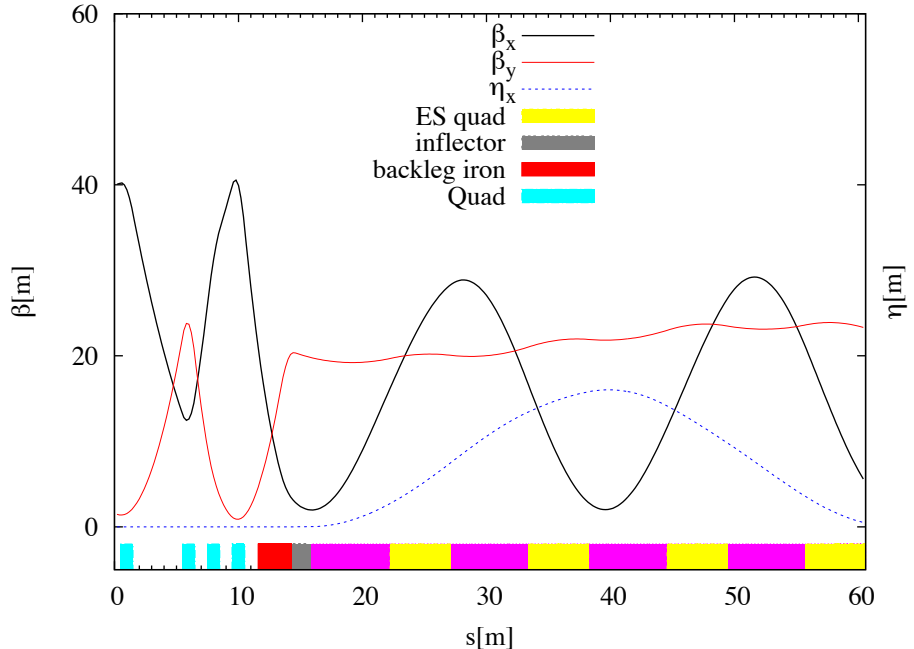
2. Dispersion at inflector exit is zero

- Ring aperture = $\pm 4\text{cm}$
- Inflector aperture = $\pm 9\text{mm}$
- Ring dispersion $\eta_{\min}=0 \Rightarrow \eta_{\max} = 2\eta_0 = 16\text{m}$
- \Rightarrow Energy acceptance in ring $\Delta E/E = 0.25 \%$
- \Rightarrow Energy acceptance through inflector $\Delta E/E = \infty$

Real acceptance = 0.25 %
and all losses are in ring

$$\sigma_x = [\beta\varepsilon + (\eta\delta)^2]^{1/2}, \quad \delta=0.05\%$$

- All energies fit through inflector.
- All muons with $\Delta E/E > 0.25 \%$ are scraped off in the ring
 - Deep modulation of beam size ($\eta_{\max}/\eta_{\min} = \infty$) at the betatron tune
 - “breathing” at twice the betatron tune



3. Dispersion at inflector exit is 1/2 match value

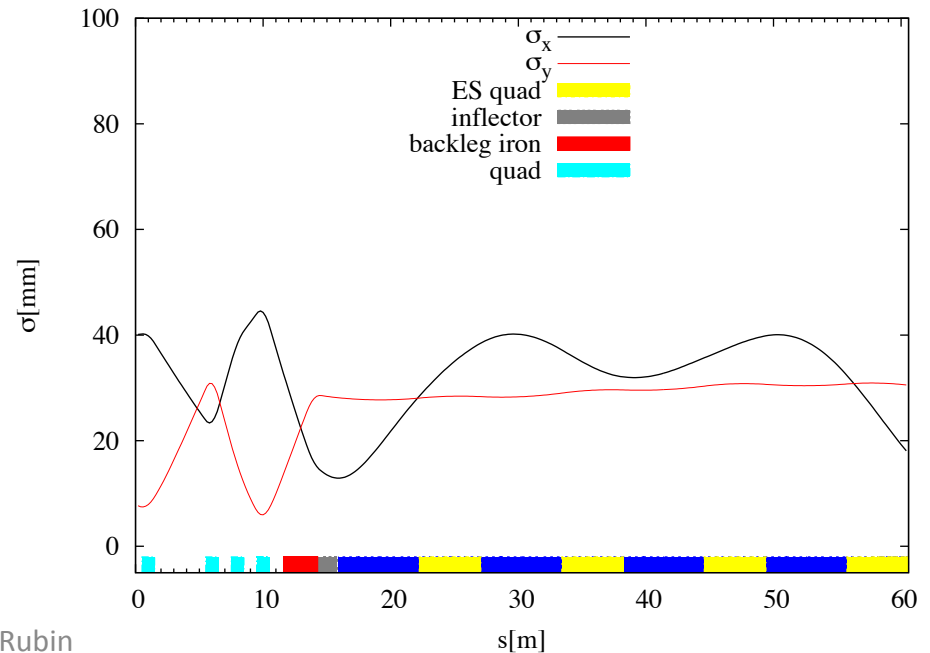
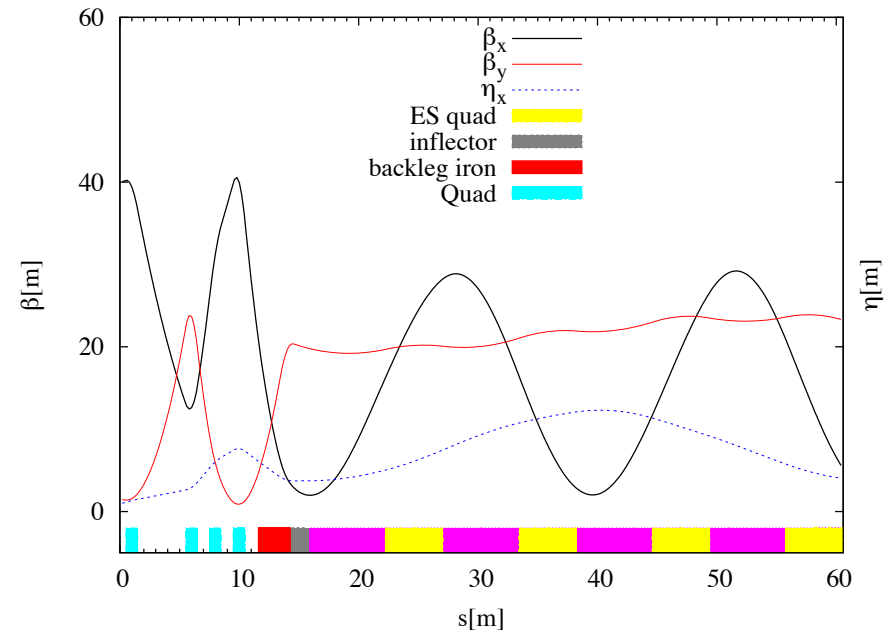
- Ring aperture = $\pm 4\text{cm}$
- Inflector aperture = $\pm 9\text{mm}$
- Ring dispersion $\eta_{\min} = \frac{1}{2}\eta_0 \Rightarrow \eta_{\max} = 1.5\eta_0 = 12\text{m}$
- \Rightarrow Energy acceptance in ring $\Delta E/E = 0.25\%$
- \Rightarrow Energy acceptance through inflector $\Delta E/E = 0.009/4 = 0.23\%$

Real acceptance = 0.23 %
with no losses in ring

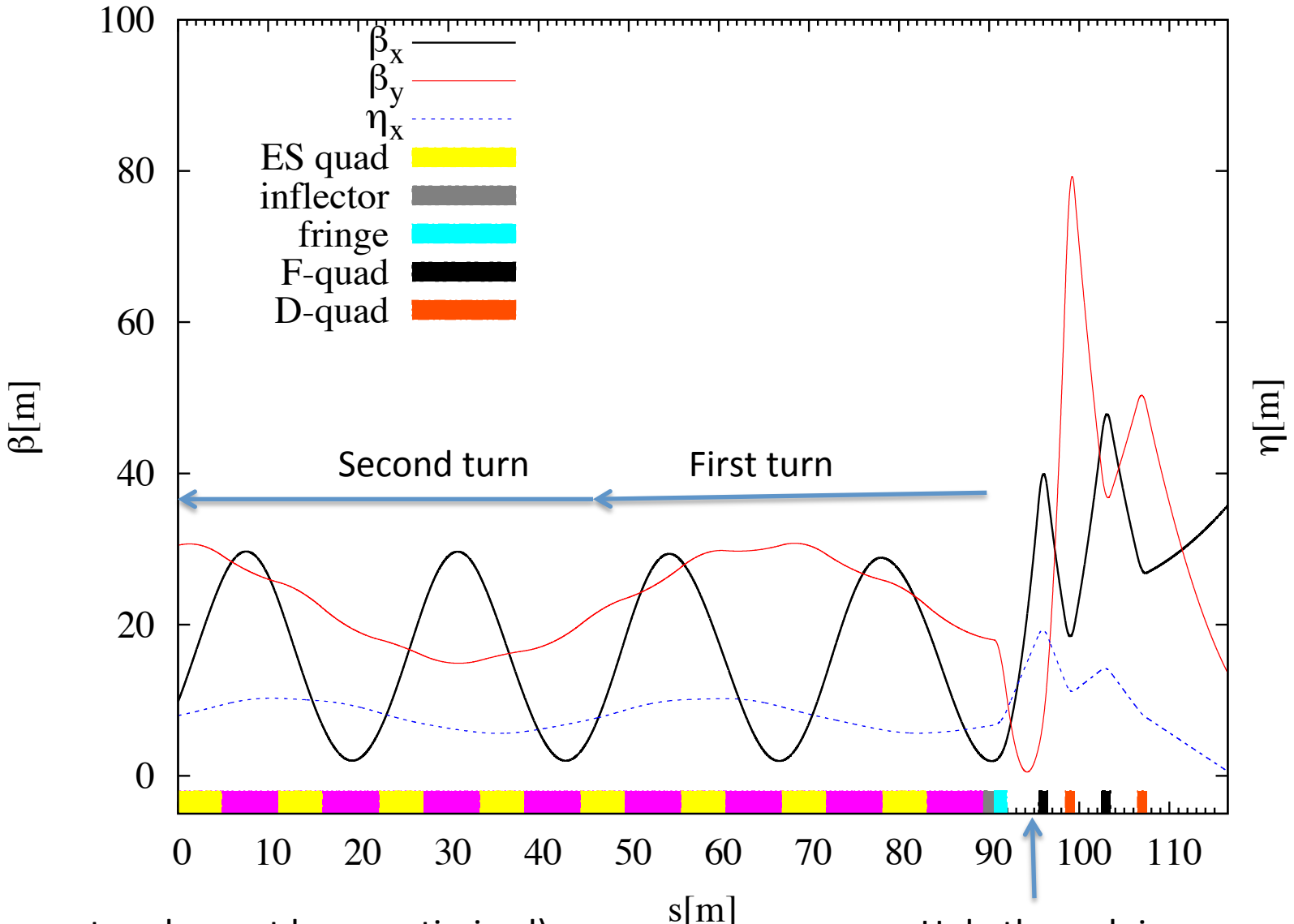
$$\sigma_x = [\beta\varepsilon + (\eta\delta)^2]^{1/2}, \quad \delta = 0.023\%$$

Acceptance of inflector \leq acceptance of ring.

- All muons with $\Delta E/E > 0.23\%$ are lost in inflector
- Modest modulation of beam size ($\eta_{\max}/\eta_{\min} = 3$) at the betatron tune
- “breathing” at twice the betatron tune



β, α at $s=0$ are chosen so that 40mm-mrad beam fits through inflector aperture.
 Quads upstream of iron (F-quads and D-quads) optimized to achieve reasonable values at 120m

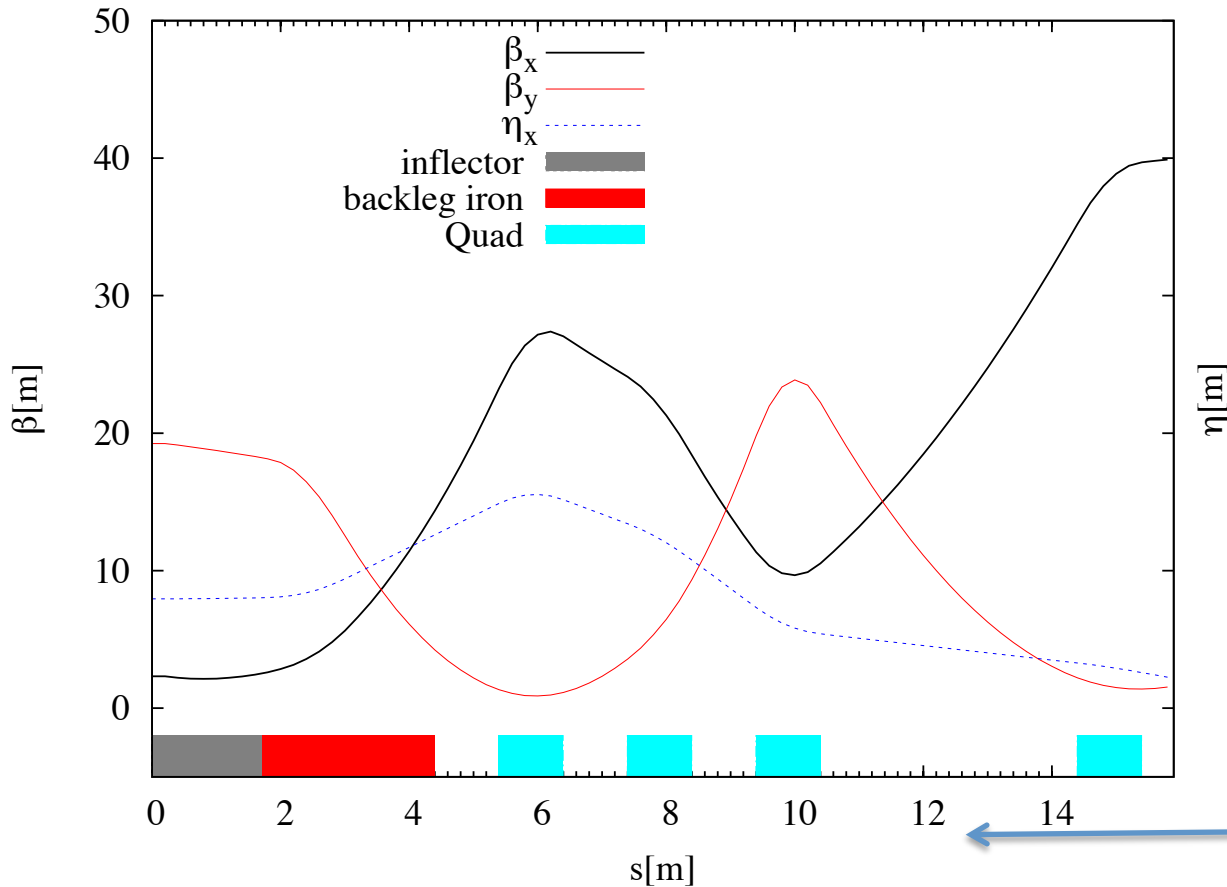


(Energy aperture has not been optimized)

Hole through iron

40 mm-mrad beam clears inflector if at exit
 $\beta_x=2.45$, $\alpha_x=-0.41$, $\beta_y=19.1$, $\alpha_y=0.045$

$\sigma_E/E = 0.15\%$ clears inflector if at exit
 $\eta=7.96$, $\eta'=0.057$

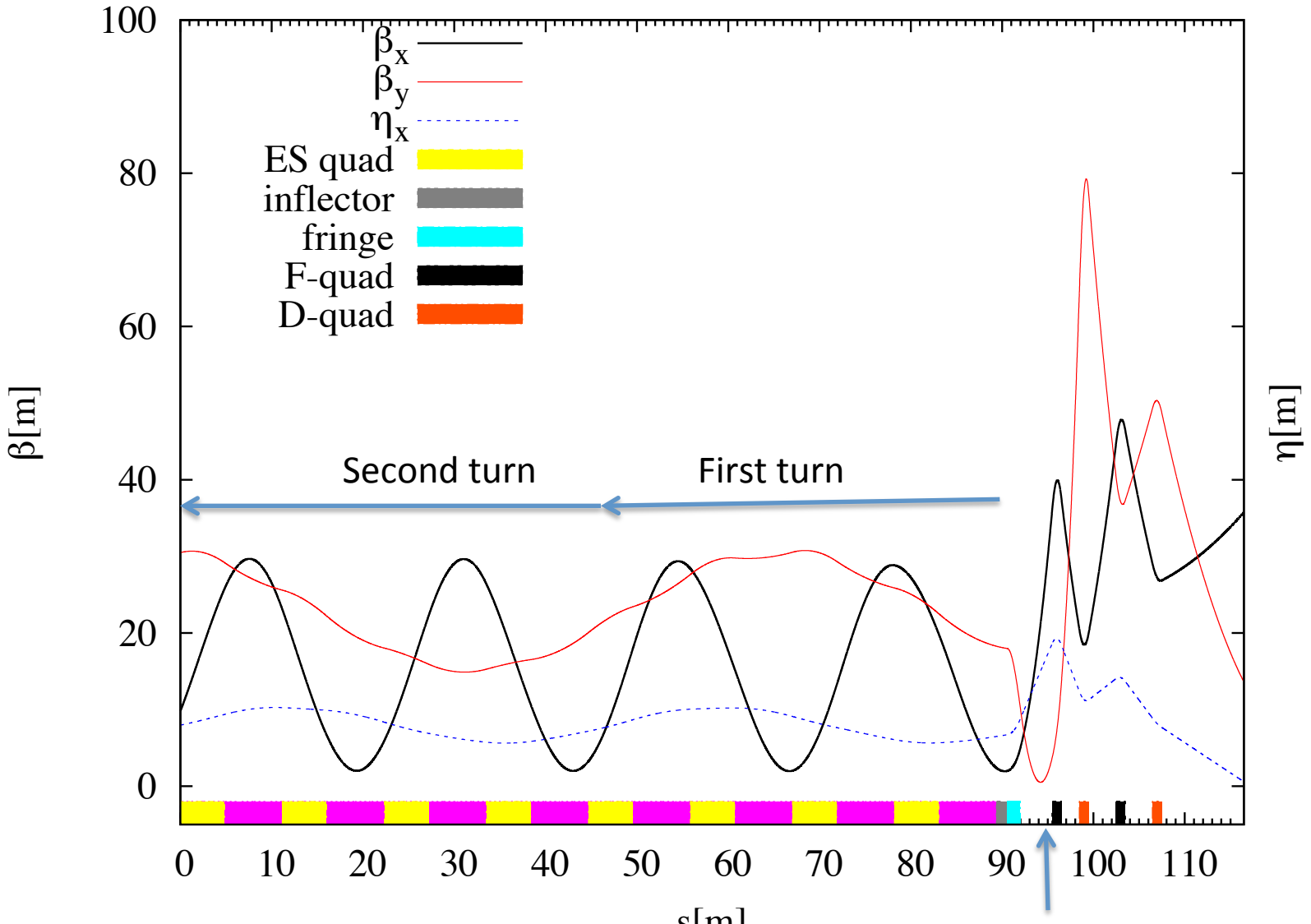


Quadrupoles
 optimized to minimize
 α, η, η' at entrance to
 beam line and with
 "reasonable" β

Beam travels right to left

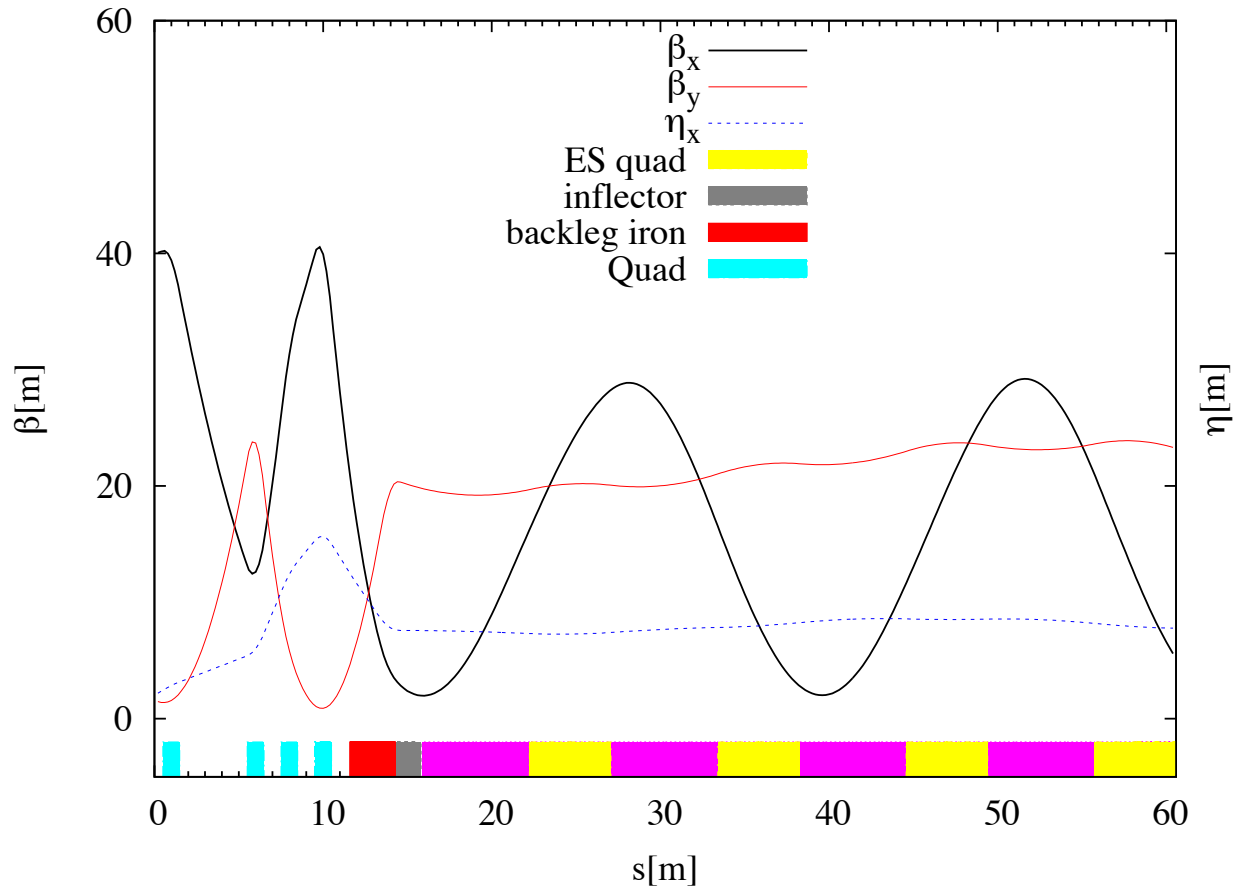
Propagate backwards from downstream end of inflector through backleg iron and through beam line quadrupoles

β, α at $s=0$ are chosen so that 40mm-mrad beam fits through inflector aperture.
 Quads upstream of iron (F-quads and D-quads) optimized to achieve reasonable values at 120m



(Energy aperture has not been optimized)

Hole through iron

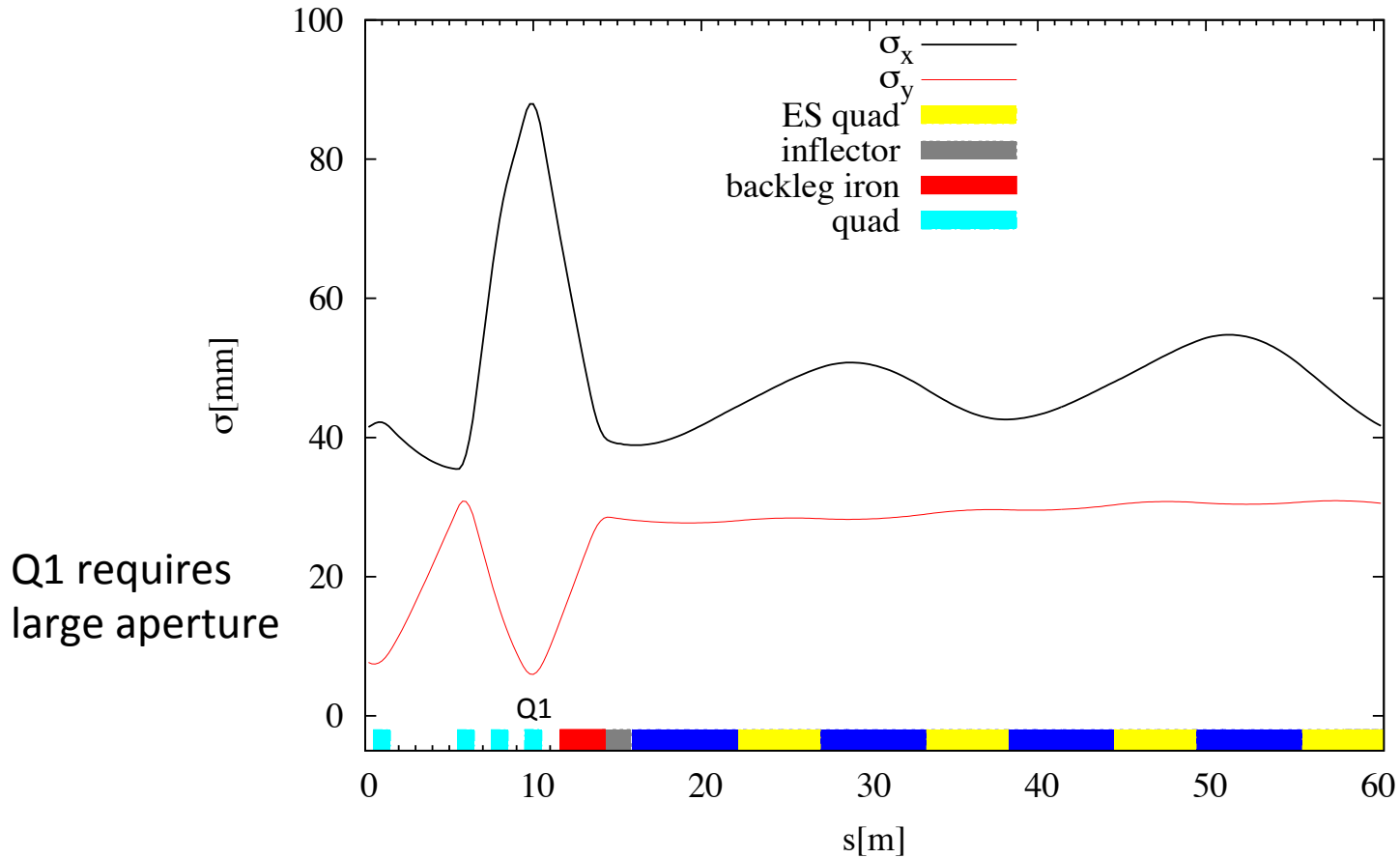


Beam travels left to right

Propagate forward through beam line quadrupoles, backleg iron, inflector and into ring

$$\sigma = [\beta\varepsilon + (\eta\delta)^2]^{1/2}$$

$$\varepsilon = 40 \text{ mm-mrad}, \quad \delta = 0.5\%$$



Propagate forward through beam line quadrupoles, backleg iron, inflector and into ring